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The diffusive shock acceleration of a population of relativistic electrons on internal shocks is one of the main scenarios to account for the multi-wavelength (MWL) flux variability observed in relativistic jets of active galactic nuclei. In addition to observations of flux variability, constraints are also provided by very-long-baseline interferometry (VLBI), which shows a large variety of moving and standing emission zones with distinct behavior.

We will present a model combining relativistic magneto-hydrodynamic jet simulations (MPI-AMRVAC code) with radiative transfer (RIPTIDE code). We simulate the evolution of standing and moving emission zones in the jet and study their MWL signatures from the radio to the X-ray band by taking into account relativistic effects (Doppler beaming and light crossing) effect (LCE)).

We focus our attention on strong interactions between a fast moving shock and stationary recollimation shocks, to study how such events lead to a significant perturbation of the stationary jet structure.

Sufficiently strong shock - shock interactions lead to the appearance of trailing components, which appear in the wake of the leading moving shock. We characterize such relaxation shocks by two observational markers, one the *fork pattern* is visible in the radio band in the time-distance plot of bright VLBI components and one at higher frequencies under the form of *flare* echoes. Our results provide a coherent interpretation of radio VLBI observations in several radio galaxies.

Jet setup

Our relativistic jet is,

- \rightarrow cylindrical and supersonic;
- \rightarrow in over-pressure (8) compared to a static and uniform ambient medium (AM);
- \rightarrow subject to a perturbation (ejecta);
- \rightarrow set with $L_{\rm kin} = 10^{46} \, {\rm erg. s^{-1}}$ the total kinetic luminosity (9) and $R_{\text{jet}} = 0.1$ pc the jet radius (10).

Component	$ ho ~(\mathrm{cm}^{-3})$	$p \; (\mathrm{dyn} \cdot \mathrm{cm}^{-2})$	γ
AM	10^{3}	1	1
jet	0.1	1.5	3
ejecta	0.1	1.5	24

Table 1: Parameters used in simulation.



Figure 1: Electron number density $n_{\rm e}$ and $\gamma_{\rm e,max}$ in the jet. The main perturbation is located around 120 $R_{\rm jet}$ following by a relaxation shock.



Figure 3: Core separation of components vs. time plot adapted from (11) observed with the VLBA 2 cm Survey and MOJAVE monitoring programs.

Application to 3C 111

CHARACTERISTIC MULTI-WAVELENGTH EMISSION SIGNATURES FROM STRONG SHOCK-SHOCK INTERACTIONS IN PERTURBED RELATIVISTIC JETS

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> MPI-AMRVAC (Message-Passing Interface Adaptive Mesh Refinement Advection Code) (1) :

Input : ambient medium, jet and ejecta initial conditions.

- radiative cooling is taken into account;
- energy and evolved in time in the fluid.

Output : 2D save states of the jet fluid and electron fluid.



 \rightarrow 3C 111 is a FR II type AGN where trailing components have been observed by (11) correlated with standing radio nodes (C1 & **C2**) (respectively at 0.1 and 3 mas from the radio core) (12, 13)

Our scenario

 \Rightarrow Fork patterns are visible on two different periods with similar observed behaviors in simulations positioned on standing shocks;

 \Rightarrow First event: drift of a perturbed standing shock $(\mathbf{C1} \rightarrow \mathbf{F})$ after interaction with a leading one (\mathbf{E}) (Fig. 2, left);

Second event: propagation of several relaxation shocks (E2, E3 & **E4**) after interaction between **C2** and **E1** ($\mathbf{E} \rightarrow \mathbf{E1}$).

- poral resolution);



 \blacksquare Shock detection method visible in (2) where the electrons fluid is injected;

Evolution of the electrons fluid upper cutoff Lorentz factor $\gamma_{e,\max}(3)$ where synchrotron

 \blacksquare The turbulent magnetic field strength B_{turb} is set as a fraction of 1% of the thermal

Input : 2D save states from MPI-AMRVAC, observation angle θ_{obs} and frequency ν .

- lowing approximations proposed by (6);
- effect (LCE) following method proposed by (7);
- and computation of the observed synchrotron flux.

Output : 2D synchrotron maps as observed at the frequency ν .



Conclusion and prospects

 \checkmark Simulations of strong shock - shock interactions lead to a great variety of emission regions;

 \checkmark Two observational markers of relaxation shocks : fork events and flare echoes are victims of observational limitations (spatial and tem-

 \checkmark Comparison with 3C111 seems promising, our scenario may be applied to other sources;

 \checkmark Characterization of relaxation shocks can help us to constrain the jet physics and to test the likelihood of the "shock - shock" scenario.

- Addition of the synchrotron polarization;
- Simulations of more complex jets (dedicated simulations on $M \, 87$ jet) and the reproduction of the MWL (radio to VHE γ \rightarrow SSC implemented!) broadband emission and variability associated with a standing shock (as HST-1 in M 87 in 2005).



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RIPTIDE code (*Radiation and Integration Processes with Time Dependence*) (4, 5):

Computation of the synchrotron emissivity and absorption parameters in each cell fol-

■ Transformation imposed by relativistic effects such as Doppler beaming and light crossing

Resolution of the radiative transfer equation along a given line of sight for an observer

 \Rightarrow allows to probe the standing shock structure through the small flaring scale.

