#### Axions as Hot Relics

The QCD Axion

Solution: Th QCD Axion

Axions via Gluons

Axion via Quarks

Axion via Leptons

Axions via Pions

# Cosmic Axion Background: the QCD axion as a hot relic

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Winter Meeting ICCUB 2022

<sup>1</sup>In collaboration with R.Z. Ferreira, F. Rompineve, F. D'Eramo, F. Arias-Aragon, J.L. Bernal, L.Merlo.

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Axions via Pions • In QCD lagrangian a term is allowed:

$$\mathcal{L}_{ heta} = rac{lpha_{m{s}}}{8\pi} heta m{G}_{\mu
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•  $G_{\mu\nu}\tilde{G}^{\mu\nu} = \partial_{\mu}K^{\mu}$ : total derivative  $\implies$  no classical effect

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  - Violates P and T (or equivalently, P and CP)
  - Periodicity:  $\theta = \theta + 2\pi$ .

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  - One effect: Neutron Electric Dipole Moment (nEDM)  $d_n = 5 \times 10^{-16} \theta$  e cm

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  - Measurement  $d_n < \mathcal{O}(10^{-26})$  e cm  $\implies |\theta| < 10^{-10}$

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- Why so small?

Axions as Hot Relics	Promote $\theta$ to a new scalar field, QCD Axion $(\theta \rightarrow \frac{a}{f})$ :
The QCD Axion	
Solution: The QCD Axion	
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Axions as Hot Relics

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Axions via Pions Promote  $\theta$  to a new scalar field, QCD Axion ( $\theta \rightarrow \frac{a}{f}$ ):

• Solves the "Strong CP problem"

$$\mathcal{L}_{a} = \frac{\alpha_{s}}{8\pi} \frac{a}{f} G_{\mu\nu} \tilde{G}^{\mu\nu}$$

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Axions via Pions Promote  $\theta$  to a new scalar field, QCD Axion  $(\theta \rightarrow \frac{a}{7})$ :

Solves the "Strong CP problem"

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• Integrating by parts:  $\mathcal{L}_{a} = \frac{\alpha_{s}}{8\pi} \frac{\partial_{\mu} a}{f} K^{\mu}$ ,  $\implies$  continuous shift symmetry  $a \rightarrow a + c$ (No potential)

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• Boundary term sensitive to QCD Instantons,

1 Induces a potential  $V(a) \propto -\cos(a/f)$ ;

- 2  $a \rightarrow 0 \implies$  Drives  $\mathcal{GP}$  to zero
- S  $\implies$  Axion mass  $m_a \approx \sqrt{V''}|_{a=0} = 0.57 \left(\frac{10^7 GeV}{f}\right) eV$

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● f (Axion "decay constant") ⇔ m<sub>a</sub>



The QCD Axion

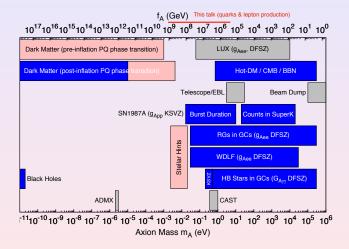
Solution: The QCD Axion

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Axion

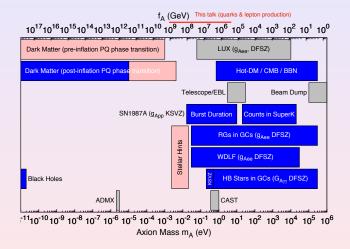
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 Byproduct: Axion can be Dark Matter (classical oscillating field) NOT our goal here, see tomorrow



Axion

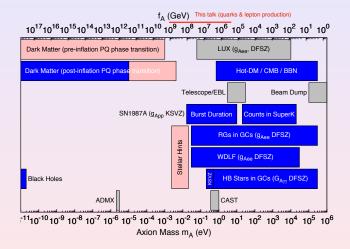
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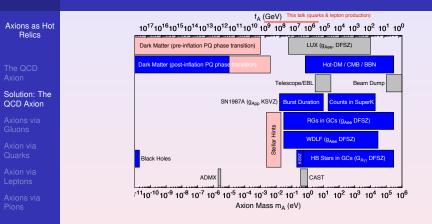
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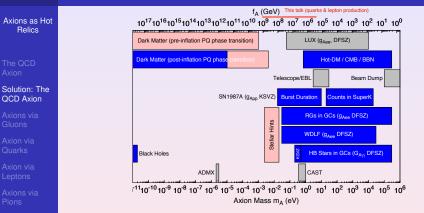
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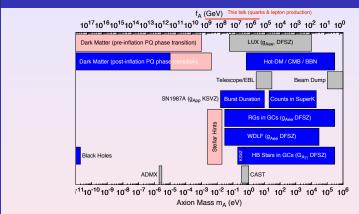
 Axion quanta produced by scatterings → Dark Radiation, visible in CMB (*Cosmic Axion Background*, similar to Cosmic Microwave Background, and Cosmic Neutrino Background)

Axions as Hot

Relics

Solution: The

**QCD** Axion



- Axion quanta produced by scatterings → Dark Radiation, visible in CMB (*Cosmic Axion Background*, similar to Cosmic Microwave Background, and Cosmic Neutrino Background)
- Caveat: Constraints based on individual couplings with *e*, *γ*, nucleons... Expected *O*(1/*f*), but model dependent.

# **QCD** Axion

#### Axions as Hot Relics

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#### Solution: The QCD Axion

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- Axion effective lagrangian:
  - May couple with continuous shift symmetry with all SM

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Only breaking: Instanton-induced (tiny) mass

Axions as Hot Relics

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Axions via Pions • Due to  $\frac{\alpha_s}{8\pi} \frac{a}{f} G_{\mu\nu} \tilde{G}^{\mu\nu}$  QCD Axions can be produced by gluon scatterings in the Early Universe

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- Due to  $\frac{\alpha_s}{8\pi} \frac{a}{f} G_{\mu\nu} \tilde{G}^{\mu\nu}$  QCD Axions can be produced by gluon scatterings in the Early Universe
- Can be produced at high *T* and decouples at  $T \lesssim T_{DEC}$  $\rightarrow$  hot relic (dark radiation)

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(M.Turner, 1987; Masso, F. Rota, and G. Zsembinszki, 2003, Salvio, Strumia, Xue, 2014)

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Scattering rate (via gluons) vs. Hubble



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Figure: (Massò et al. Phys.Rev. D66 (2002).).

 $\Gamma_{s} \equiv \langle \sigma v \rangle \cdot n_{g}^{EQ} = \left( \frac{\alpha_{s}}{2\pi f} 
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• At 
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 thermal equilibrium

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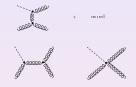


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  ight)^2 g_s^2 T^3$  vs.  $H pprox rac{T^2}{M_{Pl}}$ .
- At  $T > T_{DEC} \equiv$  thermal equilibrium
- Example: •  $f = 10^8 GeV \implies T_{DEC} \approx TeV$ •  $f = 10^9 GeV \implies T_{DEC} \approx 100 TeV$

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#### Axions via Gluons

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Axions via Pions

- If a particle:
  - Was in equilibrium at  $T > T_{DEC}$

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- 2 Decouples at some  $T \lesssim T_{DEC}$
- Has negligible mass

Axions as Hot Relics

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  - Has negligible mass
- After decoupling is dark radiation, (if *m* ≪ O(0.1 ~ 1*eV*)) (like neutrinos)
- $\implies$  Observable by CMB (and BBN)

(mostly affects expansion rate, Matter-Radiation equality...)

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- Traditionally parameterized by effective neutrino number

•  $N_{\rm eff} = 3.046 + \Delta N_{eff}$ 

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- Traditionally parameterized by effective neutrino number
- $N_{\rm eff} = 3.046 + \Delta N_{eff}$

• 
$$\Delta N_{eff} pprox rac{13.6}{g_{*,DEC}^{4/3}}$$

# $\Delta N_{\rm eff}$ diluted by $g_{*,DEC}$

Axions as Hot Relics

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Solution: The QCD Axion

#### Axions via Gluons

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Axion via Leptons

Axions via Pions  Abundance ΔN<sub>eff</sub> diluted if total number of relativistic species in the plasma g<sub>\*,DEC</sub> is large

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# $\Delta N_{\rm eff}$ diluted by $g_{*,DEC}$

Axions as Hot Relics

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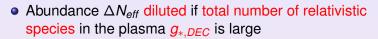
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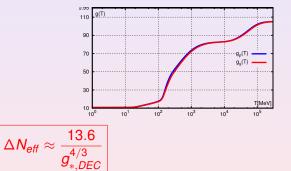
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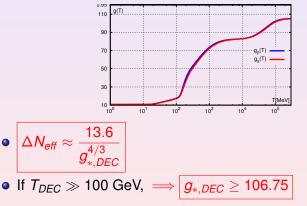
# $\Delta N_{\rm eff}$ diluted by $g_{*,DEC}$

Axions as Hot Relics

#### Axions via Gluons

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• Abundance  $\Delta N_{eff}$  diluted if total number of relativistic species in the plasma  $g_{*,DEC}$  is large



•  $\Rightarrow \Delta N_{eff} \leq 0.027$  (only upper bound!) (marginally detectable,  $1\sigma$ , by CMB-Stage 4 experiments)

Axions as Hot Relics

- The QCD Axion
- Solution: The QCD Axion

Axions via Gluons

#### Axion via Quarks

Axion via Leptons

Axions via Pions  If f ≤ 10<sup>9</sup>-10<sup>10</sup> GeV dominant channels can be via quarks & leptons <sup>2</sup> with T<sub>DEC</sub> ≤ Electroweak scale

<sup>&</sup>lt;sup>2</sup>A.N. & R.Z.Ferreira, PRL 2018; D'Eramo, Ferreira, A.N., Bernal JCAP 2018, F. Arias-Aragón et al. JCAP 2021.

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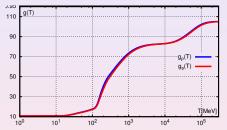
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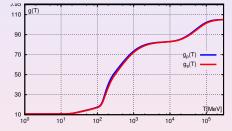
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### ADVANTAGES:

 $\bigcirc g_*^{SM} \text{ is smaller } \implies \text{ larger } N_{eff}$ 

2 Here we are confident on  $g_*^{SM} \implies \text{Precise predictions}$ 

3 Lower  $f \implies$  more accessible by direct searches (CAST, IAXO)

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Axions via Pions • If a is directly coupled to SM heavy quarks (c, b, t):

$$\mathcal{L}_{a-q} = \partial_{\mu}a \sum_{i} rac{c_{i}}{2f} \bar{q}_{i} \gamma^{\mu} \gamma^{5} q_{i} \,,$$

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• Scattering rate (via quarks, *e.g.*  $qg \leftrightarrow qa$ ) vs. Hubble



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• If  $m_q = 0 \implies$  the vertex vanishes

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Scattering rate (via quarks, e.g. qg ↔ qa) vs. Hubble



• If  $m_q = 0 \implies$  the vertex vanishes

Indeed:

- This coupling can be rotated away  $q 
  ightarrow e^{irac{c_{l}a}{2t}\gamma^{5}}q$
- But it reappears in the mass term  $m_q \bar{q} e^{i \frac{c_l^a}{f} \gamma^5} q$

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$$\mathcal{L}_{a-q} = \partial_{\mu} a \sum_{i} \frac{c_{i}}{2f} \bar{q}_{i} \gamma^{\mu} \gamma^{5} q_{i},$$

Scattering rate (via quarks, e.g. qg ↔ qa) vs. Hubble



• If  $m_q = 0 \implies$  the vertex vanishes

Indeed:

- This coupling can be rotated away  $q 
  ightarrow e^{irac{c_{l}a}{2t}\gamma^{5}}q$
- But it reappears in the mass term  $m_q \bar{q} e^{i \frac{c_l^a}{T} \gamma^5} q$

 $\Gamma_s = \left(\frac{c_i}{f}\right)^2 g_s^2 m_q^2 T$ 

Axions as Hot Relics

The QCD Axion

Solution: Th QCD Axion

Axions via Gluons

#### Axion via Quarks

Axion via Leptons

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$$\Gamma_{s} = \left(\frac{c_{i}}{f}\right)^{2} g_{s}^{2} m_{q}^{2} T \cdot e^{-\frac{m_{q}}{T}} \text{vs. } H \approx \frac{T^{2}}{M_{Pl}}$$

Axions as Hot Relics

The QCD Axion

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Axions via Gluons

Axion via Quarks

Axion via Leptons

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<sup>&</sup>lt;sup>3</sup> R.Ferreira & A.N., PRL 2018. See also Turner PRL 1987, Brust et al. JHEP 2013, Baumann et al. PRL 2016.

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• Ratio peaks at  $T \approx m_q$ 

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• Ratio peaks at  $T \approx m_q$ 

• Axions produced dominantly via quarks

 $1~{\rm GeV} \lesssim T \lesssim 100 {\rm GeV}$ 

• Range  $10^9 \text{GeV} \gtrsim f/c_i \gtrsim 10^7 \text{GeV}^{-3}$ (partly in tension with SN bounds, if all  $c_i = 1$ )

<sup>&</sup>lt;sup>3</sup>R.Ferreira & A.N., PRL 2018. See also Turner PRL 1987, Brust et al. JHEP 2013, Baumann et al. PRL 2016.

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- Interesting for direct detection (e.g. IAXO),  $m_a \approx 10^{-1} \sim 10^{-3} eV$ ,

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Axions as Hot Relics

- The QCD Axion
- Solution: The QCD Axion

Axions via Gluons

#### Axion via Quarks

Axion via Leptons

Axions via Pions

- $g_{*,DEC}$  is smaller at 1 GeV  $\lesssim T \lesssim 100$ GeV
- Prediction: larger  $N_{\rm eff} \lesssim 0.045$  (\*Not just upper bound!\*)

Axions as Hot Relics

The QCD Axion

Solution: The QCD Axion

Axions via Gluons

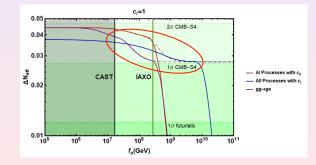
#### Axion via Quarks

Axion via Leptons

Axions via Pions

- $g_{*,DEC}$  is smaller at  $1 \text{ GeV} \lesssim T \lesssim 100 \text{GeV}$
- Prediction: larger N<sub>eff</sub> \$\le 0.045\$ (\*Not just upper bound!\*)
- Solving Boltzmann equations for n<sub>a</sub>:

(R.Ferreira & A.N., PRL 2018; F.Arias-Aragon et al. JCAP, 2021)



 $10^9 \text{GeV} \gtrsim f/c_i \gtrsim 10^7 \text{GeV}$ ,  $5 \times 10^{-3} \text{eV} \lesssim m_a \lesssim 0.5 \text{eV}$ ( $c_i = 1$ , for QCD Axion) as the set of the

Axions as Hot Relics

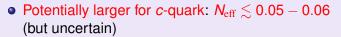
- The QCD Axion
- Solution: The QCD Axion

Axions via Gluons

#### Axion via Quarks

Axion via Leptons

Axions via Pions



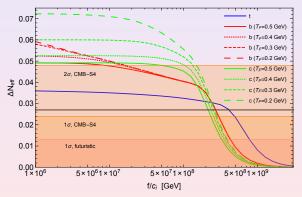


Figure: R.Ferreira & A.N., PRL 2018.

### Hot Axions via Quark Decays

Axions as Hot Relics

The QCD Axion

Solution: Th QCD Axion

Axions via Gluons

#### Axion via Quarks

Axion via Leptons

Axions via Pions

### • *a* – *q* interaction can be flavor non-diagonal

$$\mathcal{L}_{a-q} = \partial_{\mu}a \sum_{q \neq q'} \bar{q'} \gamma^{\mu} \left( \mathcal{V}_{q'q} + \mathcal{A}_{q'q} \gamma^5 \right) q + \mathrm{h.c.} \; ,$$

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### Hot Axions via Quark Decays

Axions as Hot Relics

The QCD Axion

Solution: Th QCD Axion

Axions via Gluons

#### Axion via Quarks

Axion via Leptons

Axions via Pions • a - q interaction can be flavor non-diagonal

$$\mathcal{L}_{a-q} = \partial_{\mu} a \sum_{q \neq q'} \bar{q'} \gamma^{\mu} \left( \mathcal{V}_{q'q} + \mathcal{A}_{q'q} \gamma^5 \right) q + \mathrm{h.c.} \; ,$$

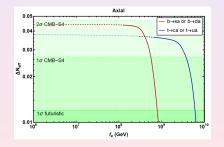


Figure: F.Arias-Aragon, F.D'Eramo, R.Z.Ferreira, A. N , L.Merlo, JCAP 2021.

• More efficient than scatterings (larger  $f/c \lesssim 10^{10} \text{ GeV}$ )

Axions as Hot Relics

The QCD Axion

Solution: The QCD Axion

Axions via Gluons

Axion via Quarks

Axion via Leptons

Axions via Pions

- The same can be done with leptons ( $\mu$  and  $\tau$ ) <sup>4</sup>
- a-electron uninteresting (strongly constrained)

Axions as Hot Relics

The QCD Axion

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Axion via Leptons

Axions via Pions

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$$\mathcal{L}_{a-\ell} = \partial_{\mu}a \sum_{i} \frac{c_{i}}{2f} \bar{\ell}_{i} \gamma^{\mu} \gamma^{5} \ell_{i},$$



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Axions as Hot Relics

The QCD Axion

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- Slightly smaller  $f/c_{\ell}$
- Ratio peaks at  $T \approx m_{\ell} \implies$  Larger  $N_{eff}$

<sup>4</sup>F.D'Eramo, A.N.,R.Z.Ferreira, J.L.Bernal, JCAP 2018 → ( ) → ( )

Axions as Hot Relics

The QCD Axion

Solution: The QCD Axion

Axions via Gluons

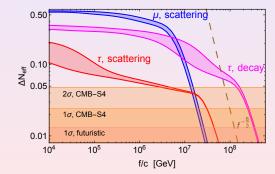
Axion via Quarks

Axion via Leptons

Axions via Pions



• Ratio peaks at  $T \approx m_{\ell} \implies \text{Larger } N_{eff}$ 



• Caveat:  $\mu$  scattering constrained by SN cooling at  $f/c_{\mu}\gtrsim 10^8 GeV$  (Bolling et al. PRL 2020, Croon et al. JHEP 2021)

## Axion-Pion coupling

Axions as Hot Relics

DFSZ model: a couples to u-type and d-type quarks,
KSVZ model: no coupling to SM fermions

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The QCD Axion

Solution: The QCD Axion

Axions via Gluons

Axion via Quarks

Axion via Leptons

Axions via Pions

## Axion-Pion coupling

Axions as Hot Relics

C

- Axions via Pions

• DFSZ model: a couples to u-type and d-type quarks, KSVZ model: no coupling to SM fermions

• Coupling to pions:

$$\mathcal{L}_{a\pi} = \frac{c_{a\pi}}{f_{\pi}} \frac{\partial_{\mu} a}{f} \left[ 2 \partial^{\mu} \pi^{0} \pi^{+} \pi^{-} - \pi_{0} \left( \partial^{\mu} \pi^{+} \pi^{-} - \pi^{+} \partial^{\mu} \pi^{-} \right) \right],$$

where

$$c_{a\pi} = -rac{1}{3}c_u^0 - c_d^0 - rac{1-z}{1+z}$$
.  $z \equiv rac{m_u}{m_d} \simeq 0.47^{+0.06}_{-0.07},$ 

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## Axion-Pion coupling

Axions as Hot Relics

C

- Axions via Pions

 DFSZ model: a couples to u-type and d-type quarks, KSVZ model: no coupling to SM fermions

**DFSZ**: 
$$c_u^0 = \frac{1}{3}\cos^2(\beta)$$
,  $c_d^0 = \frac{1}{3}\sin^2(\beta)$ ,  
**KSVZ**:  $c_u^0 = c_d^0 = 0$ ,

• Coupling to pions:

$$\mathcal{L}_{a\pi} = \frac{\mathcal{C}_{a\pi}}{f_{\pi}} \frac{\partial_{\mu} a}{f} \left[ 2 \partial^{\mu} \pi^{0} \pi^{+} \pi^{-} - \pi_{0} \left( \partial^{\mu} \pi^{+} \pi^{-} - \pi^{+} \partial^{\mu} \pi^{-} \right) \right],$$

where

$$c_{a\pi} = -rac{1}{3}c_u^0 - c_d^0 - rac{1-z}{1+z}$$
.  $z \equiv rac{m_u}{m_d} \simeq 0.47^{+0.06}_{-0.07},$ 

 KSVZ :
  $c_{a\pi} \simeq 0.12^{+0.023}_{-0.018}$ ,

 DFSZ :
  $c_{a\pi} \simeq 0.12^{+0.023}_{-0.018} - \frac{1}{9}\cos(2\beta)$ .

## CMB Bounds on DFSZ



The QCD Axion

Solution: The QCD Axion

Axions via Gluons

Axion via Quarks

Axion via Leptons

Axions via Pions

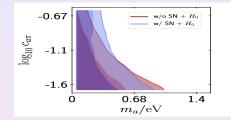


Figure: Constraints due to pion production Planck 18 + BAO (+ Pantheon + SH0ES H<sub>0</sub>)

### CMB Bounds on DFSZ

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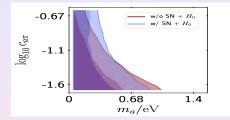


Figure: Constraints due to pion production Planck 18 + BAO (+ Pantheon + SH0ES  $H_0$ ) For DFSZ-II: muon production is also relevant for  $c_{a\pi} \leq O(0.1)$ :

DFSZ-I	Planck 18+BAO (+SN+H <sub>0</sub> )
$C_{a\pi} = 0.225$	$m_a \le 0.20 \; (0.29) \; { m eV}$
$c_{a\pi} = 0.0225$	$m_a \leq$ 0.84 (0.82) eV
DFSZ-II	Planck 18+BAO (+SN+H <sub>0</sub> )
$c_{a\pi} = 0.225$	$m_a \leq 0.20 \; (0.29) \; { m eV}$
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## CMB Bounds on DFSZ

Axions as Hot Relics

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Axions via Gluons

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Axion via Leptons

Axions via Pions

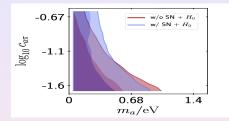


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Caveat! Perturbative pion cross-section calculation in Chiral PT breaks down at  $T \gtrsim 60 \text{ MeV} (\text{Di Luzio et al. 2021, arXiv 2101.10330.})$ 

Axions as Hot Relics

- The QCD Axion
- Solution: The QCD Axion
- Axions via Gluons
- Axion via Quarks
- Axion via Leptons
- Axions via Pions

• If  $f \leq \mathcal{O}(10^9)$  GeV, coupling with quarks and leptons (with  $c_i = \mathcal{O}(1)$ ) dominates over  $\frac{\alpha_s}{8\pi} \frac{a}{f} \tilde{GG}$ 

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2 Efficiency peaks at  $T \approx m_f$ 

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- Sor quarks (t, b): N<sub>eff</sub> ≤ 0.05 (measurable at 2σ by CMB S4) (\*maybe higher for c-quark?)

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- Solution Non-diagonal couplings  $\implies$  production via Decays more efficient ( $f \leq \mathcal{O}(10^{10})$  GeV)

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- **?** Pion production bound on DFSZ axion:  $m_a \lesssim 0.2 \text{ eV}$  (at large  $c_{a\pi}$ ), but relaxed  $m_a \lesssim 0.6 0.8 \text{ eV}$  for small  $c_{a\pi}$

(\*Caveat: Pion cross-section calculation should break down at  $T\gtrsim$  60 MeV (Di Luzio et al. 2021))

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Future CMB experiments will tell in a few years, plus direct detection (e.g. IAXO)