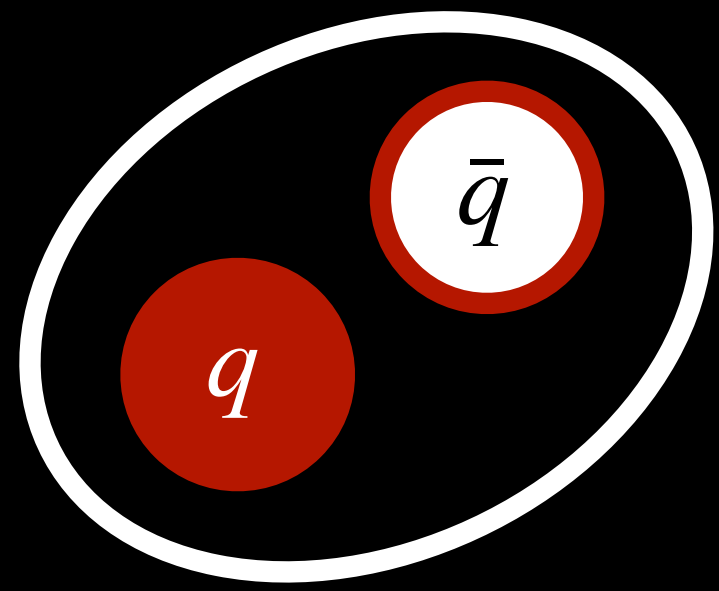


Exploring Photoproduced $\eta^{(\prime)}\pi^0$ Systems in the Search for Exotic Hadrons at GlueX

10th International Conference QNP – Barcelona, Spain
Hadron Spectroscopy

Zachary Baldwin, July 8, 2024

Constituent Quark Model



Total angular momentum | $J = 0, 1, 2, \dots$

Parity | $P = (-1)^{L+1}$

Charge Conjugation | $C = (-1)^{L+S}$

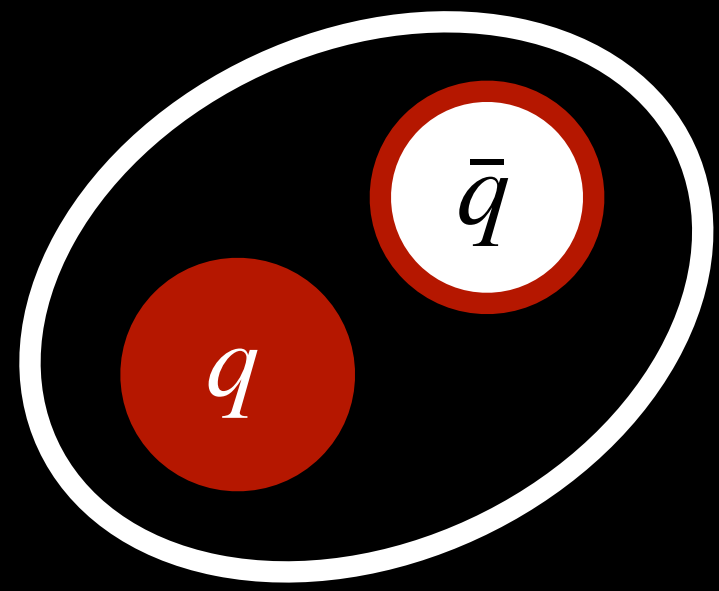
L is the relative orbital angular momentum of the q and \bar{q}

S is the total intrinsic spin of the $q\bar{q}$ pairs

Allowed J^{PC} quantum numbers

L	S	J^{PC}	L	S	J^{PC}	L	S	J^{PC}
0	0	0^{-+}	1	0	1^{+-}	2	0	2^{-+}
0	1	1^{--}	1	1	0^{++}	2	1	1^{--}
			1	1	1^{++}	2	1	2^{--}
			1	1	2^{++}	2	1	3^{--}

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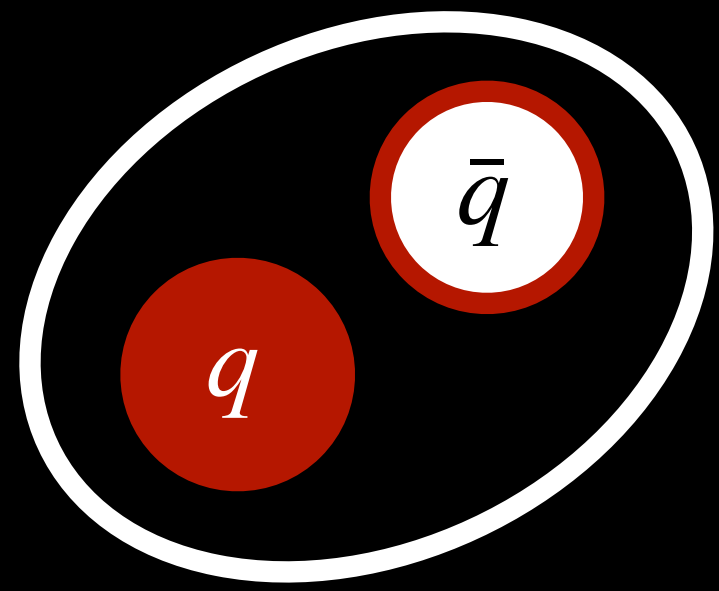
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Discovering forbidden quantum numbers would be immediate evidence of a non- $q\bar{q}$ state

Forbidden J^{PC} quantum numbers

$0^{--}, 0^{+-}, 1^{-+}, 2^{+-}$

Constituent Quark Model



Total angular momentum | $J = 0, 1, 2, \dots$

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Forbidden J^{PC} quantum numbers

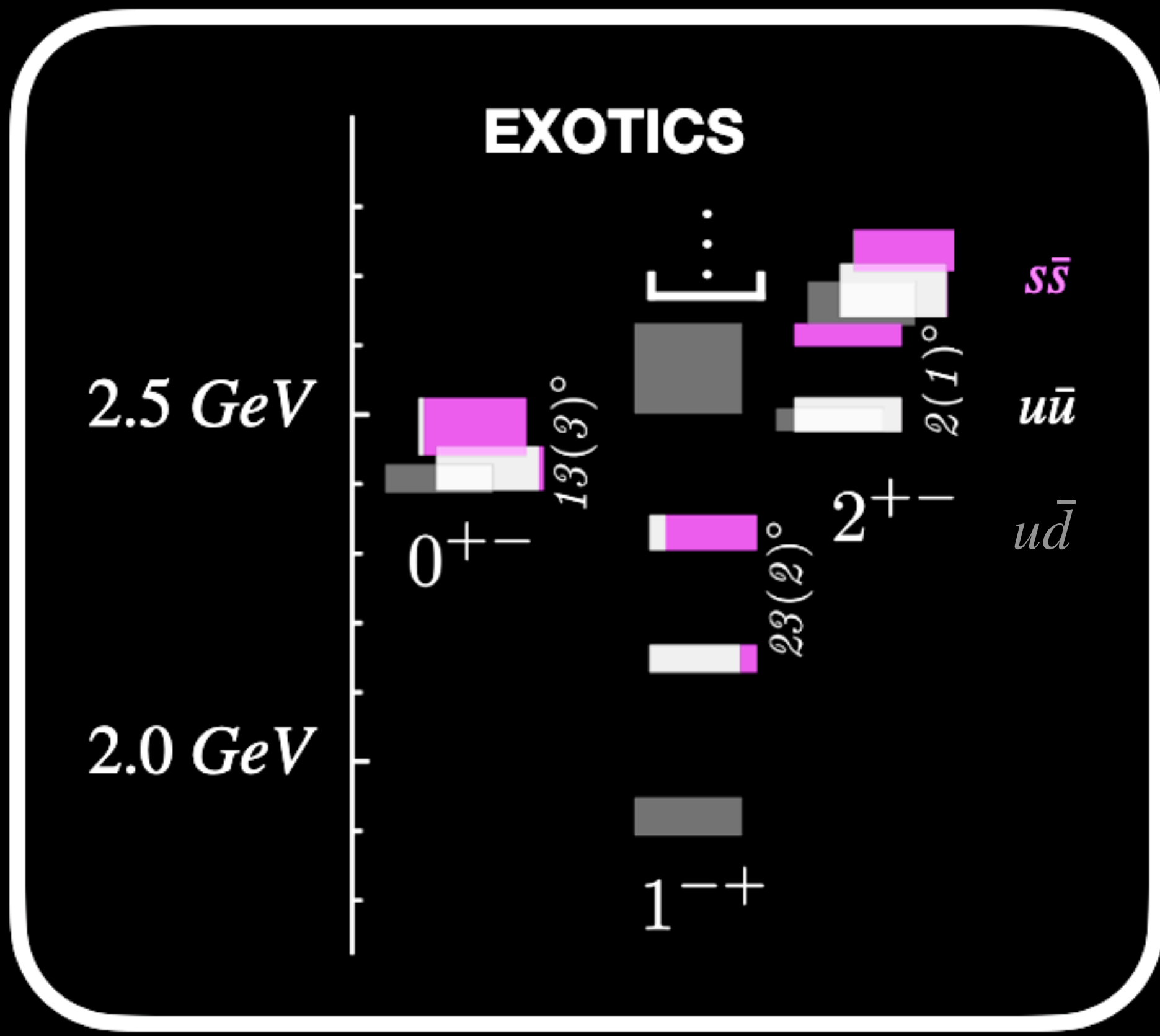
$0^{--}, 0^{+-}, 1^{-+}, 2^{+-}$

Do gluons play a larger role?

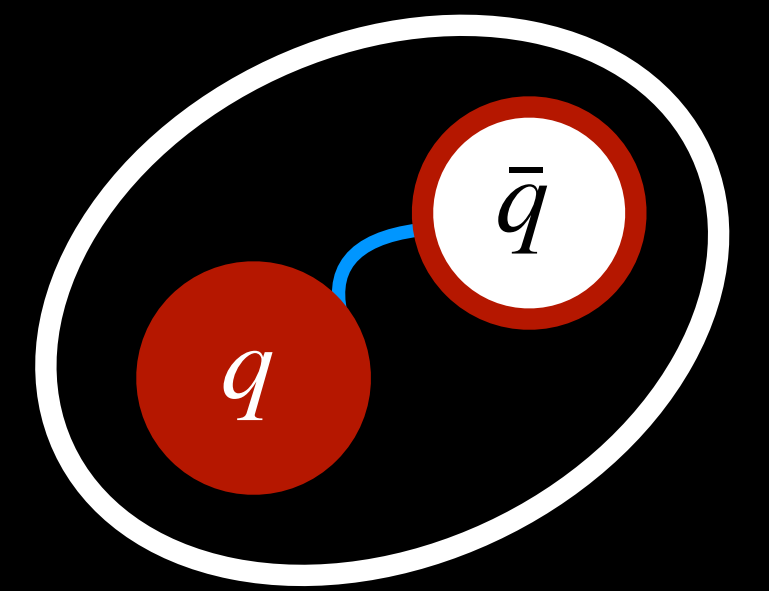
Lattice QCD predicts “gluonic excitations”, confirming mesons that are not in constituent quark model known as exotic mesons



J. Dudek et al. [Hadron Spectrum Collab], Phys. Rev. D 83, 111502 (2011)



Exotic Hybrid



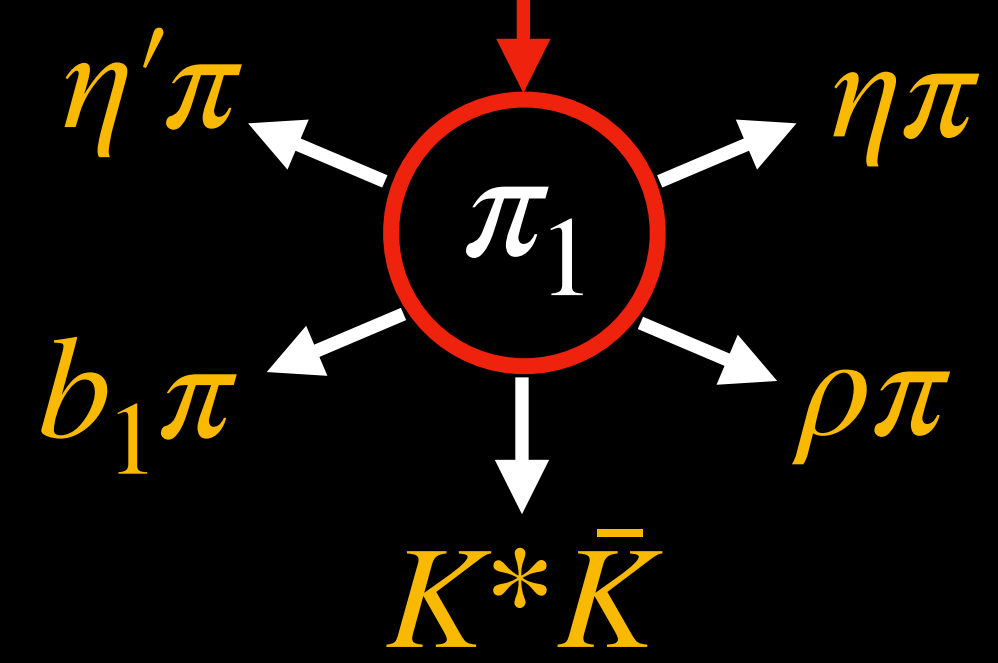
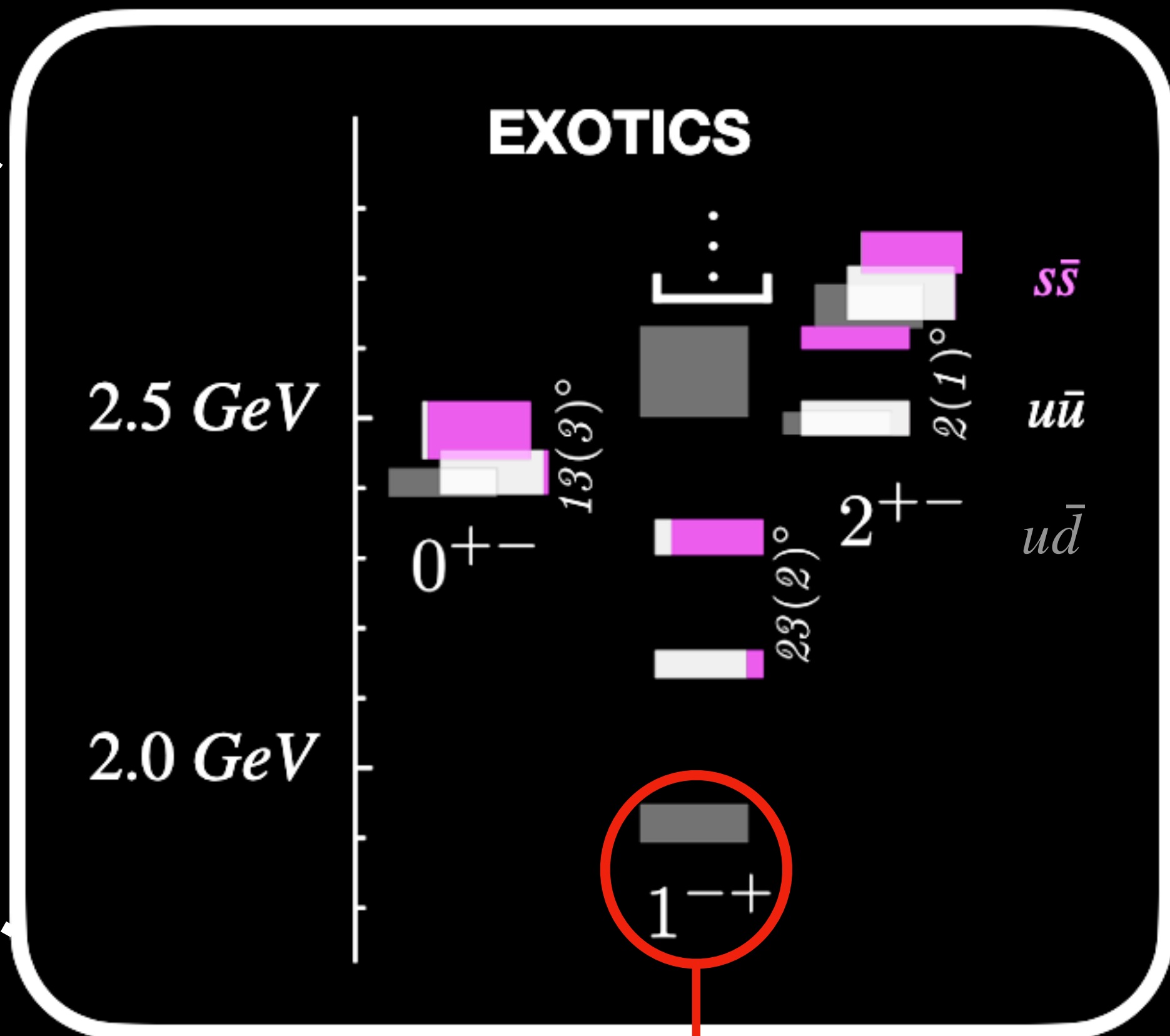
Glueball



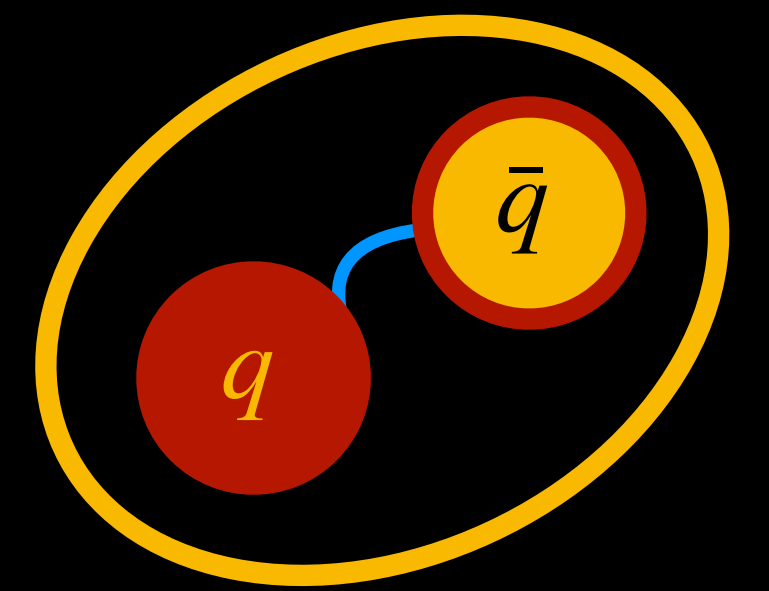
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Exotic Hybrid



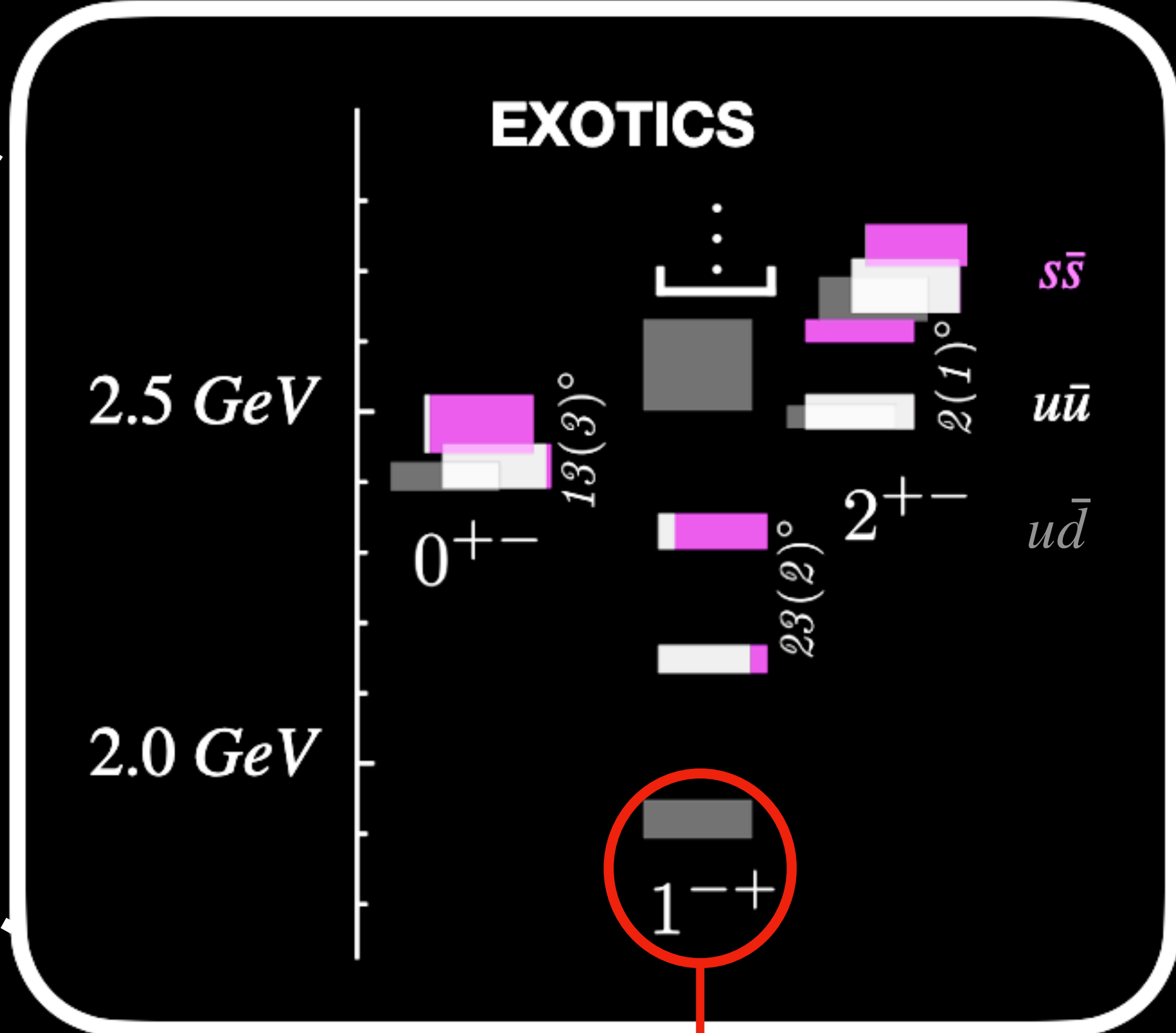
Glueball



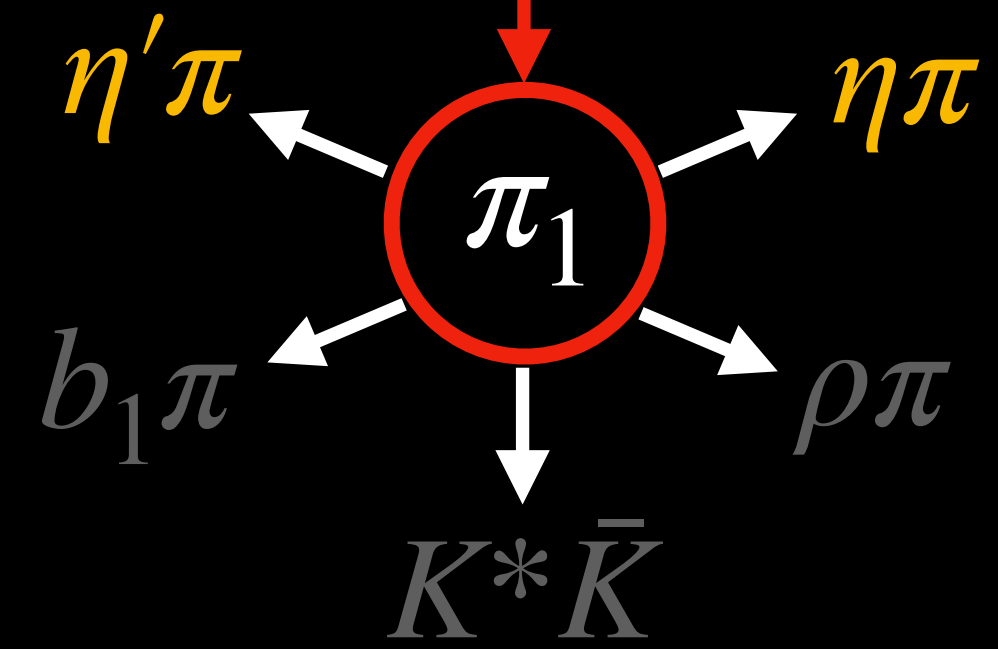
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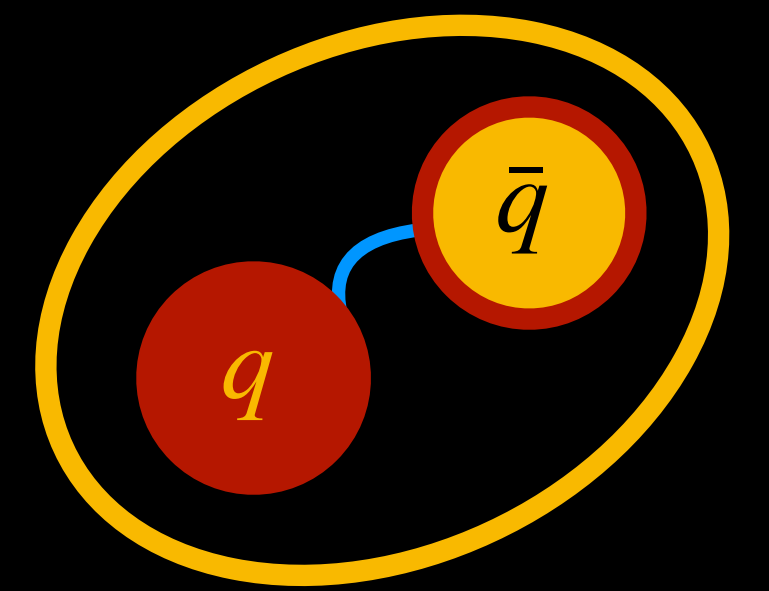
J. Dudek et al. [Hadron Spectrum Collab], Phys. Rev. D 83, 111502 (2011)



FOCUS



Exotic Hybrid



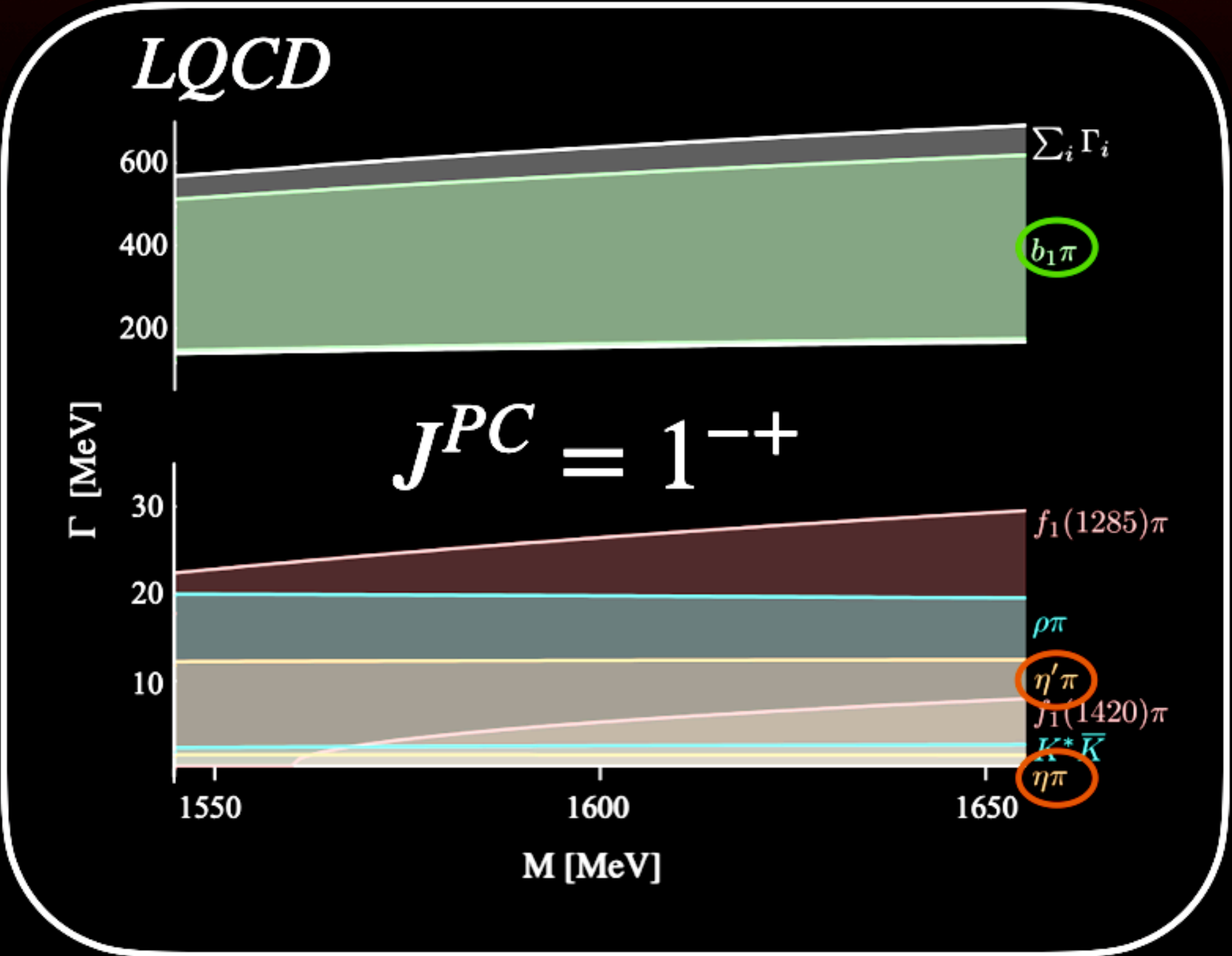
Glueball



L	S	P	D	F	...
J^{PC}	0^{++}	1^{-+}	2^{++}	3^{-+}	...

LQCD calculations provide insight into 1^{-+} state decay modes

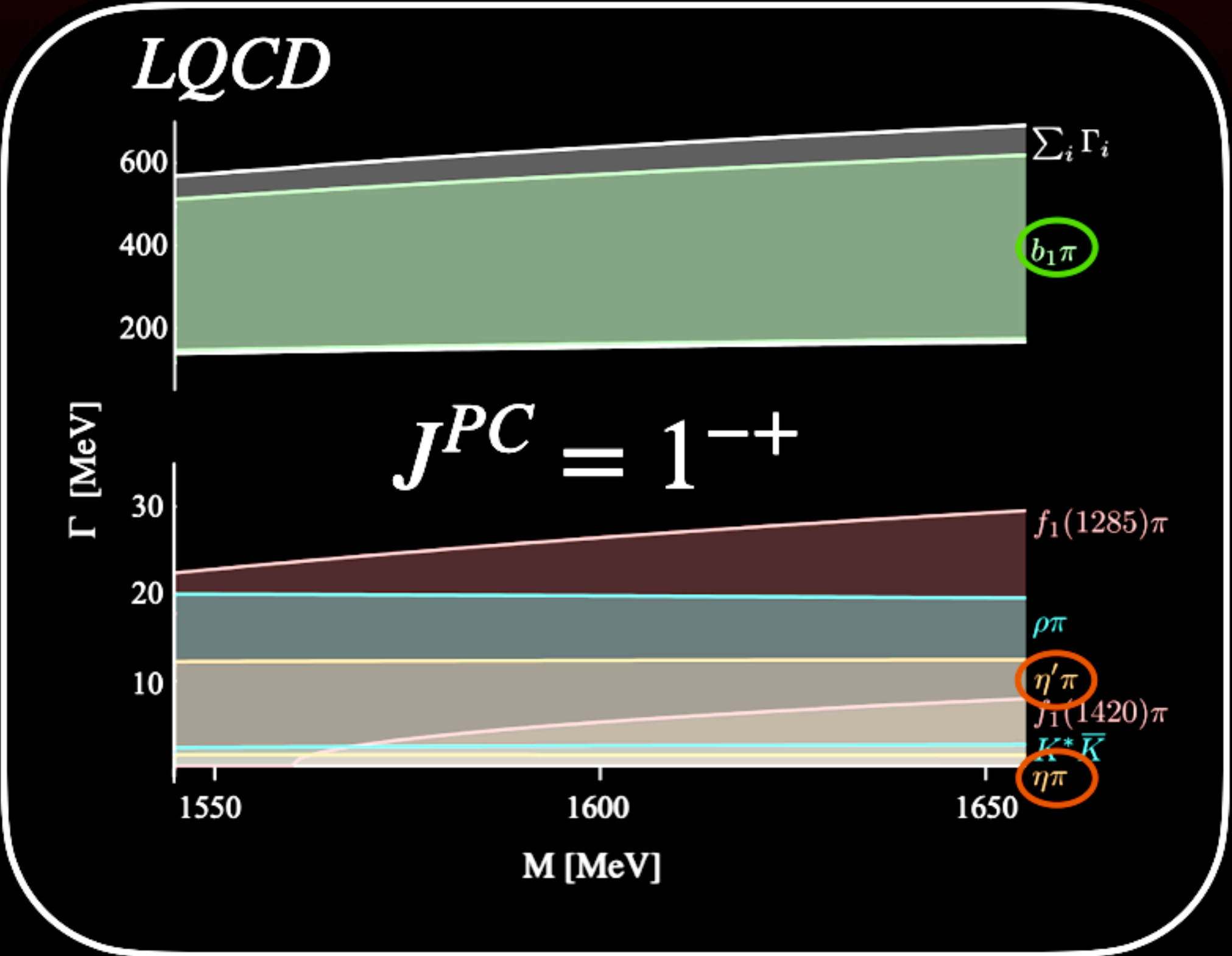
	Γ_i / MeV
$\eta\pi$	$0 \rightarrow 1$
$\eta'\pi$	$0 \rightarrow 12$
$b_1\pi$	$139 \rightarrow 529$



What percentage of π_1 is expected in $\eta^{(\prime)}\pi^0$?

A. Woss et al. [Hadron Spectrum Collab], Phys. Rev. D 103, 054502 (2021)

LQCD calculations provide insight into 1^{-+} state decay modes



	Γ_i / MeV
$\eta\pi$	$0 \rightarrow 1$
$\eta'\pi$	$0 \rightarrow 12$
$b_1\pi$	$139 \rightarrow 529$



$$\frac{B(\pi_1 \rightarrow b_1\pi)}{B(\pi_1 \rightarrow \eta\pi)} > \frac{139}{1} > 100$$

$$\frac{B(\pi_1 \rightarrow b_1\pi)}{B(\pi_1 \rightarrow \eta'\pi)} > \frac{139}{12} > 10$$

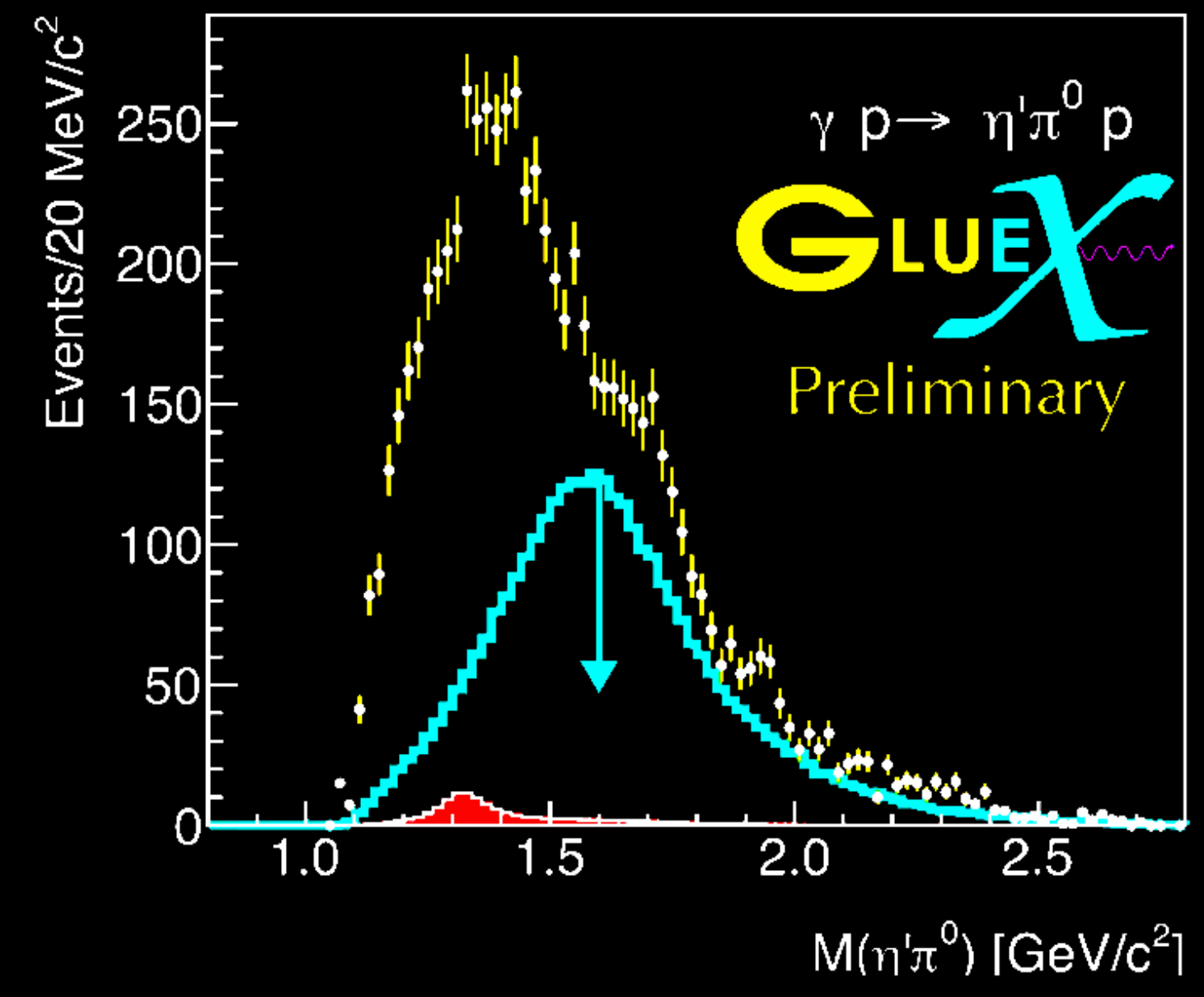
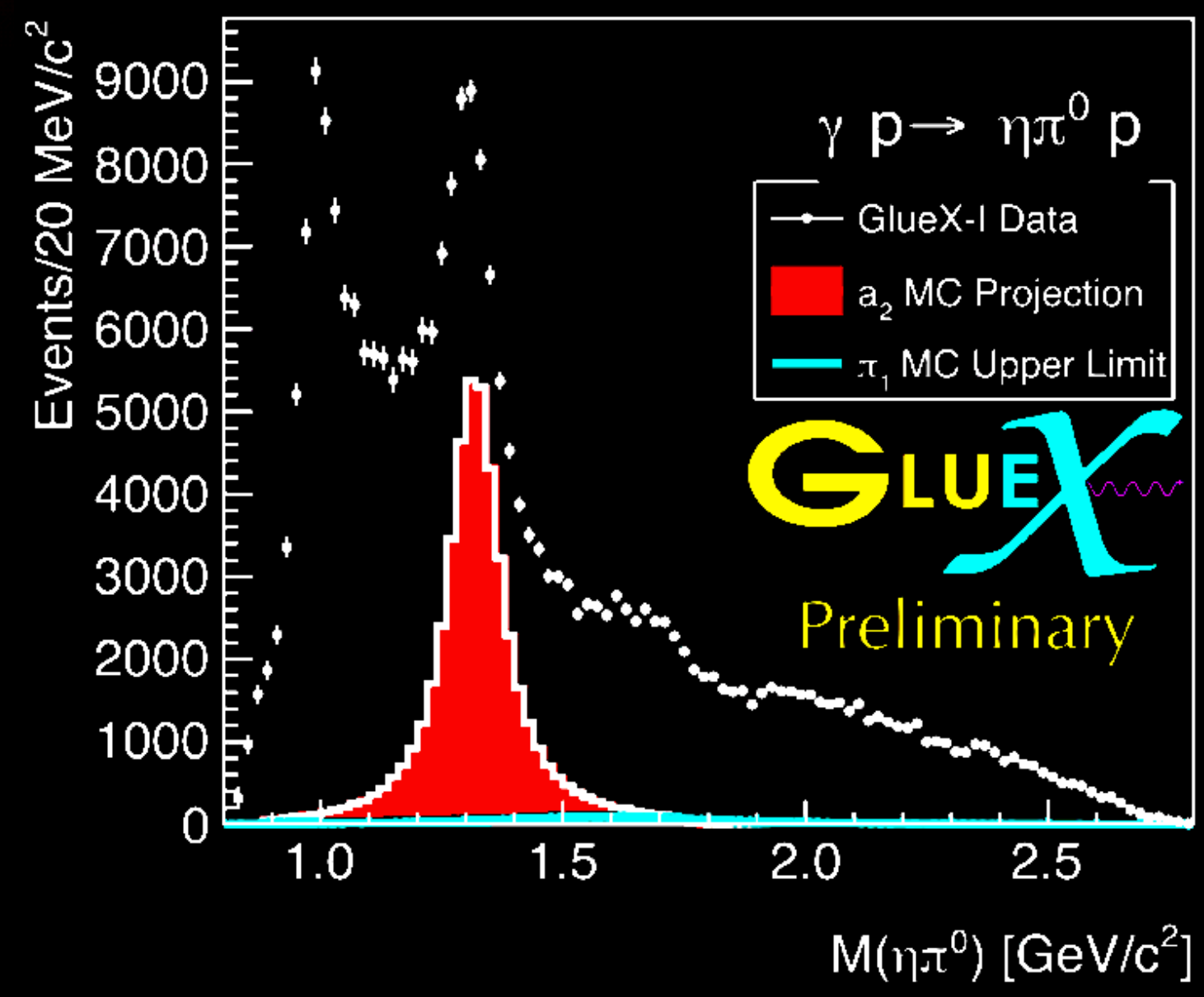
$$B(\pi_1 \rightarrow \eta\pi) \lesssim 1\%$$

$$B(\pi_1 \rightarrow \eta'\pi) \lesssim 10\%$$

A. Woss et al. [Hadron Spectrum Collab], Phys. Rev. D 103, 054502 (2021)

HOT OFF THE PRESS

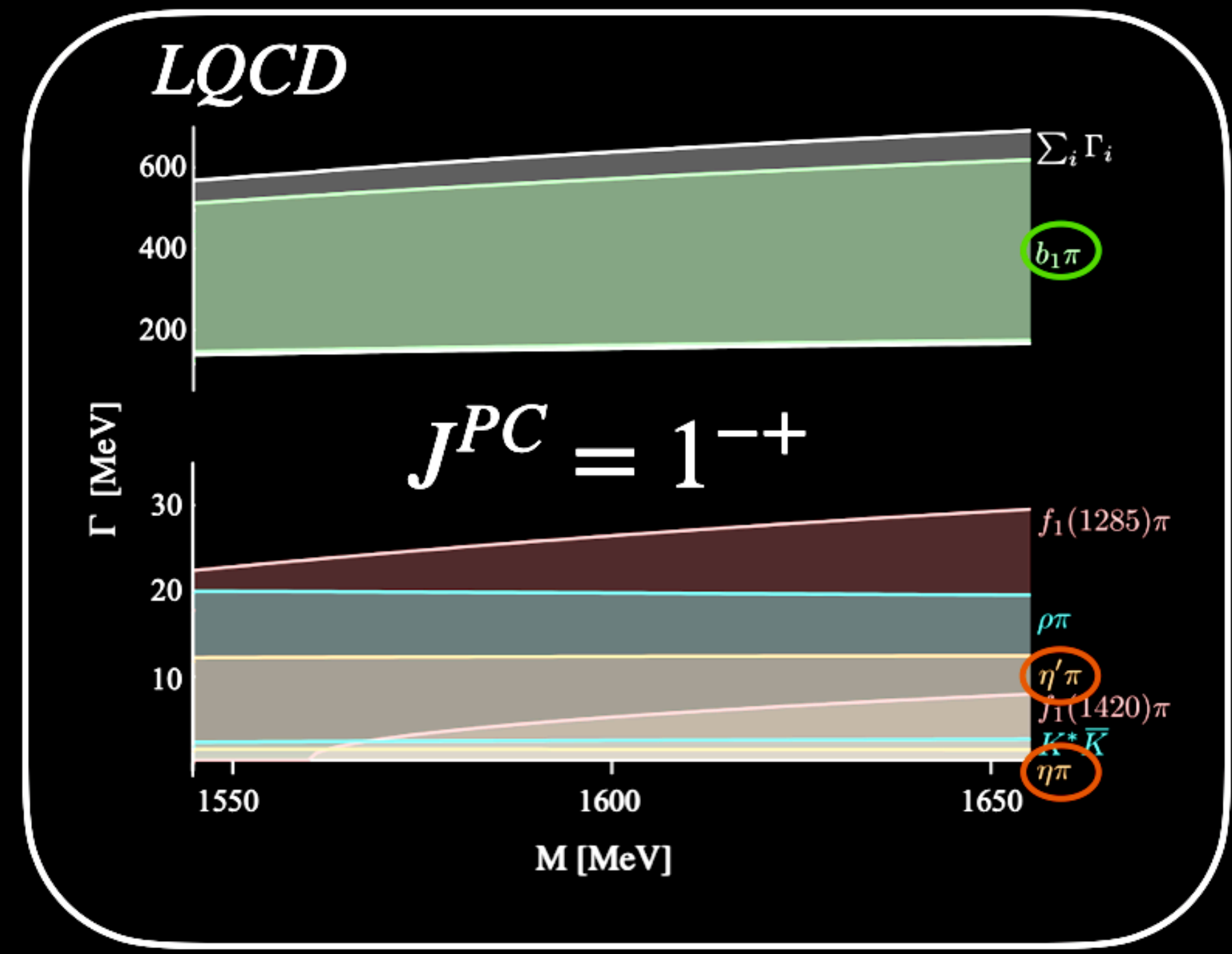
A.Fzal et al. An upper limit on the photoproduction cross section of the spin-exotic $\pi_1(1600)$, arXiv (2024)



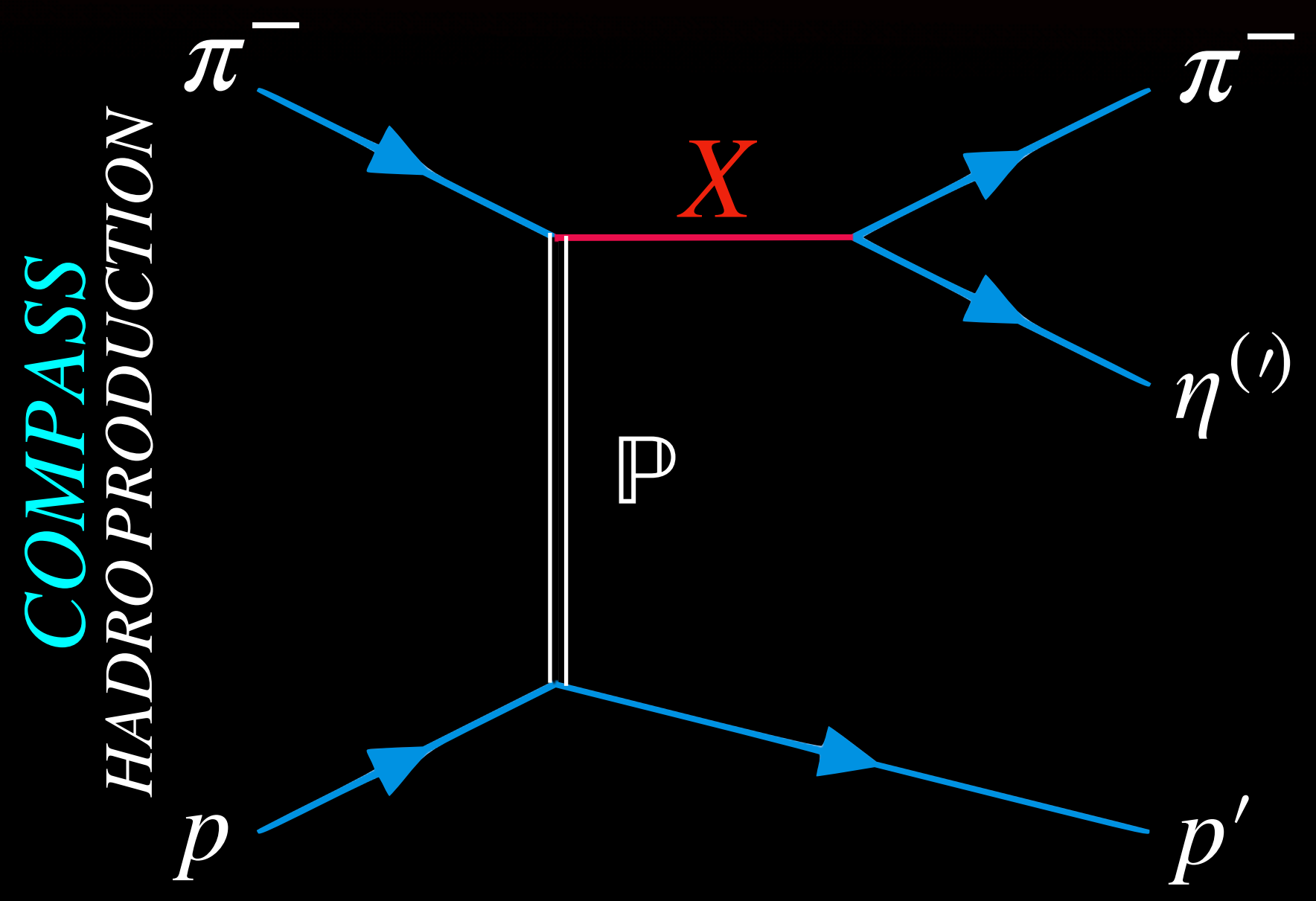
- **Based on upper limit π_1 cross section $\omega\pi\pi$**
 - first photo-produced upper limit cross sections of spin exotic candidate

- **Projection of $\pi_1 \rightarrow \eta^{(\prime)} \pi^0$ upper limit**
 - uses strong a_2 signal measurements as a reference

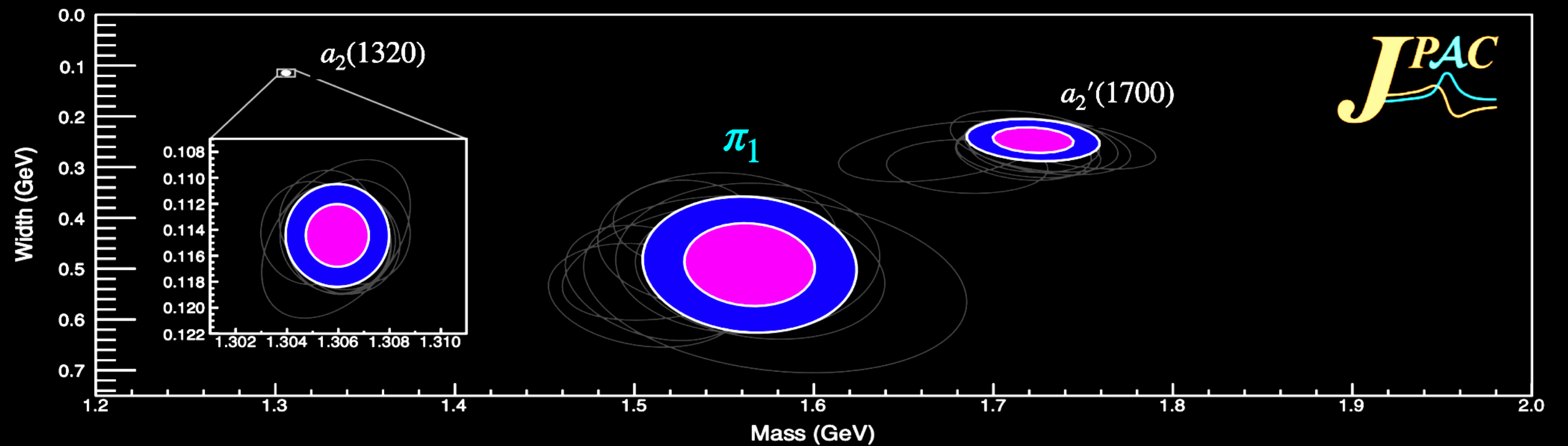
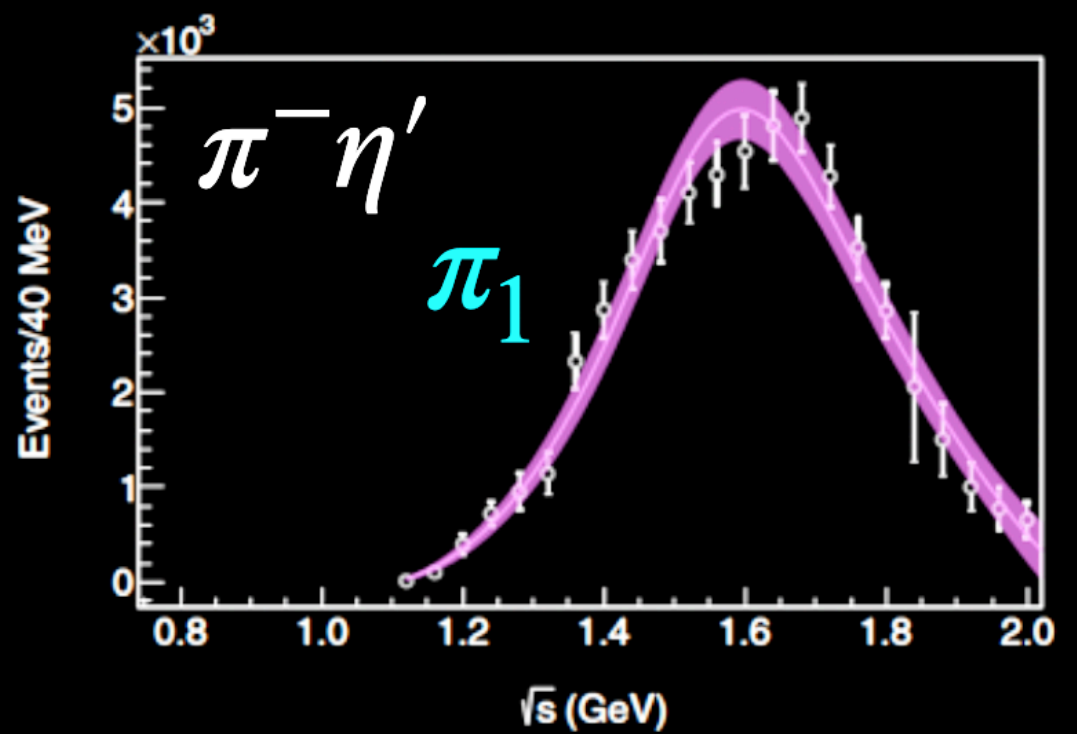
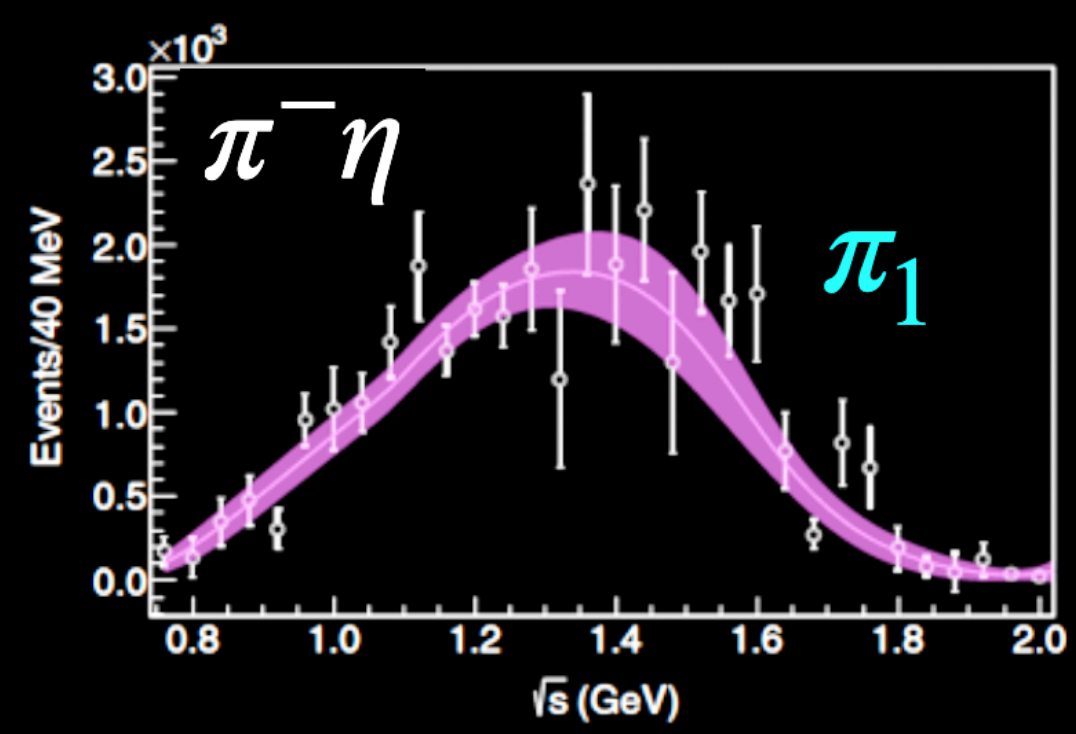
Further overview: A. Austregesilo – Tues. 9:40 am



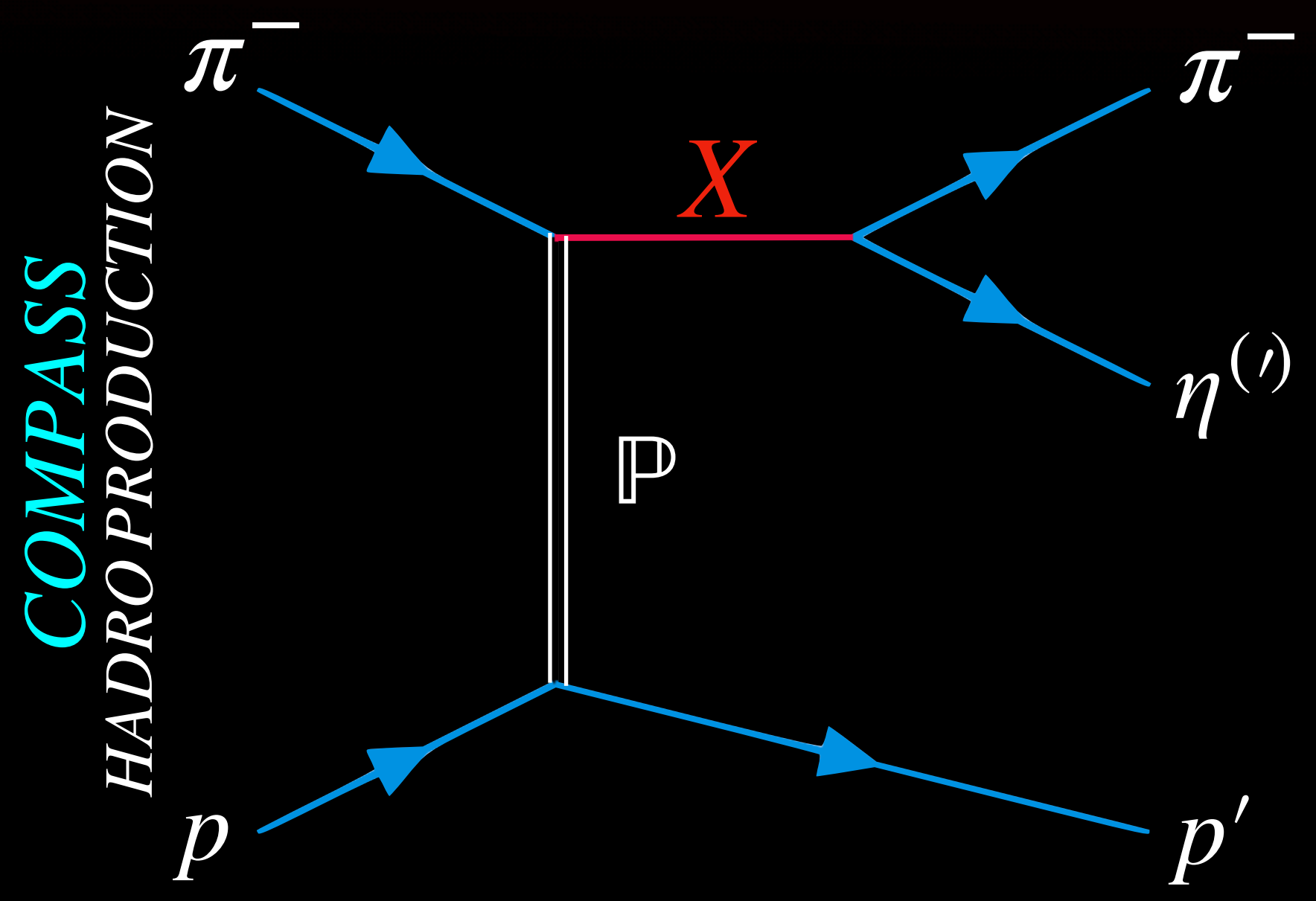
A. Woss et al. [Hadron Spectrum Collab], Phys. Rev. D 103, 054502 (2021)



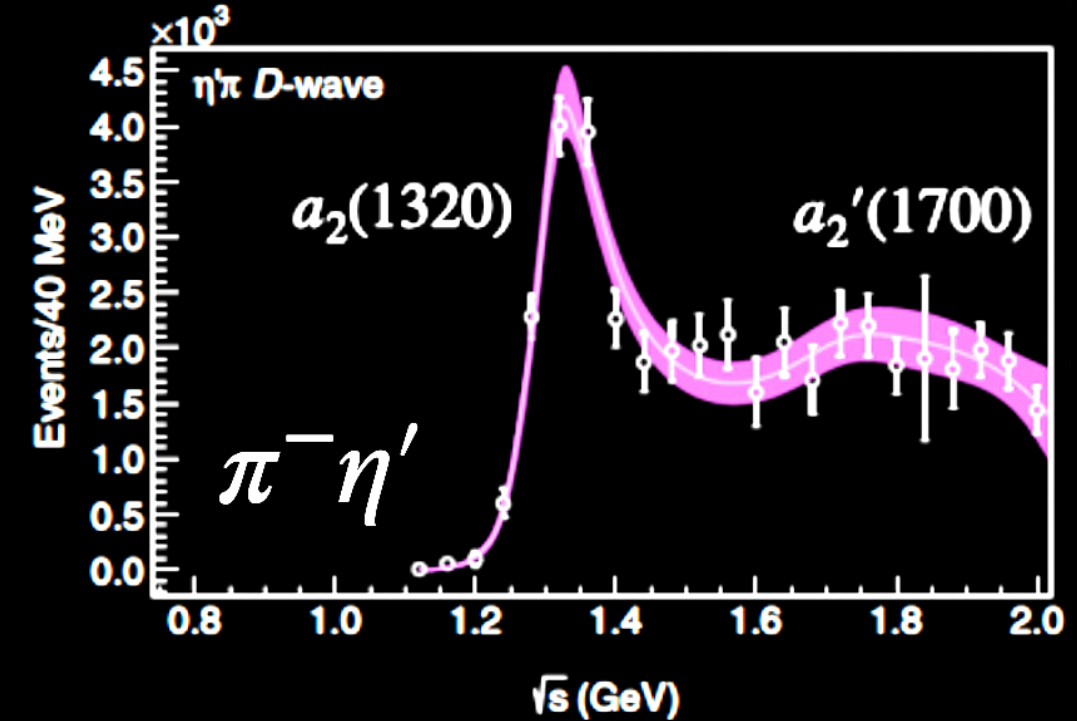
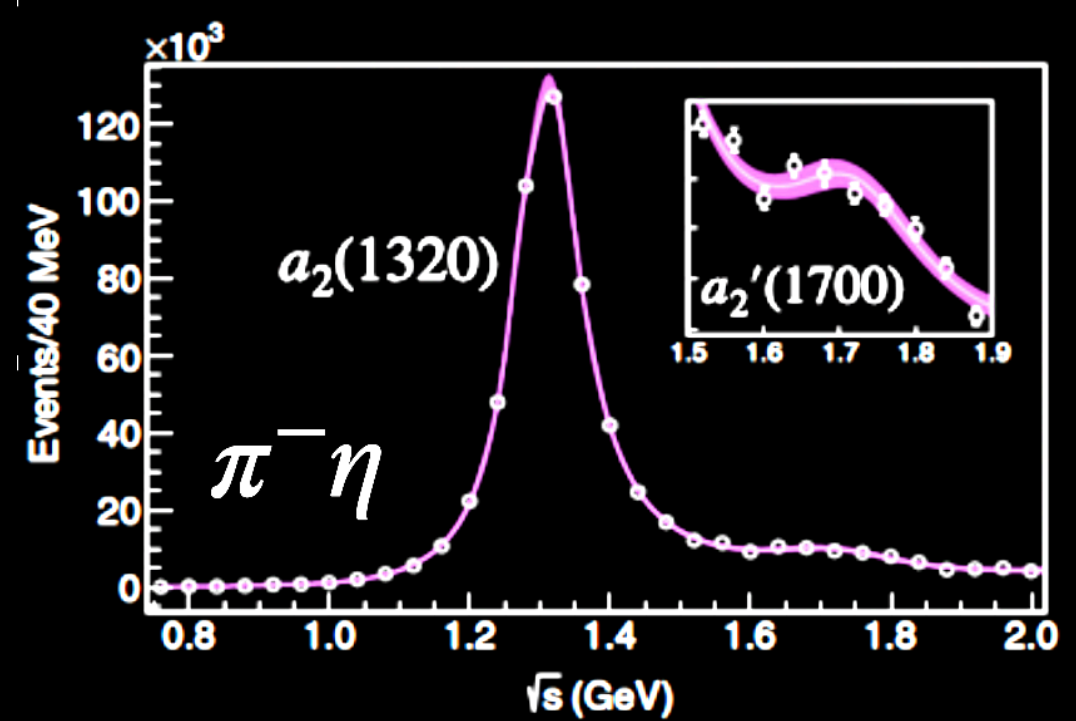
- **JPAC analysis utilizing COMPASS data**
 - extended with $\bar{p}p$ data
B.Kopf et.al. Eur. Phys. J. C 81, 1056 (2021)
- **Coupled channel fit to both $\eta^{(1)}\pi$ systems**
 - also described the dominate a_2 meson



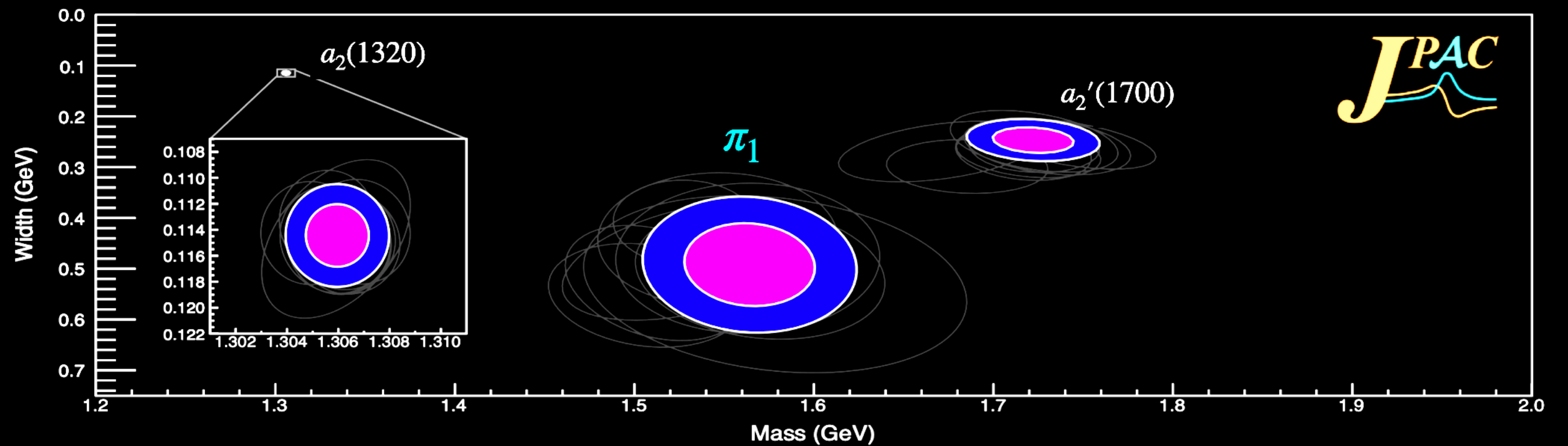
A. Rodas et al. [Joint Physics Analysis Center], PRL 122, 042002 (2019)



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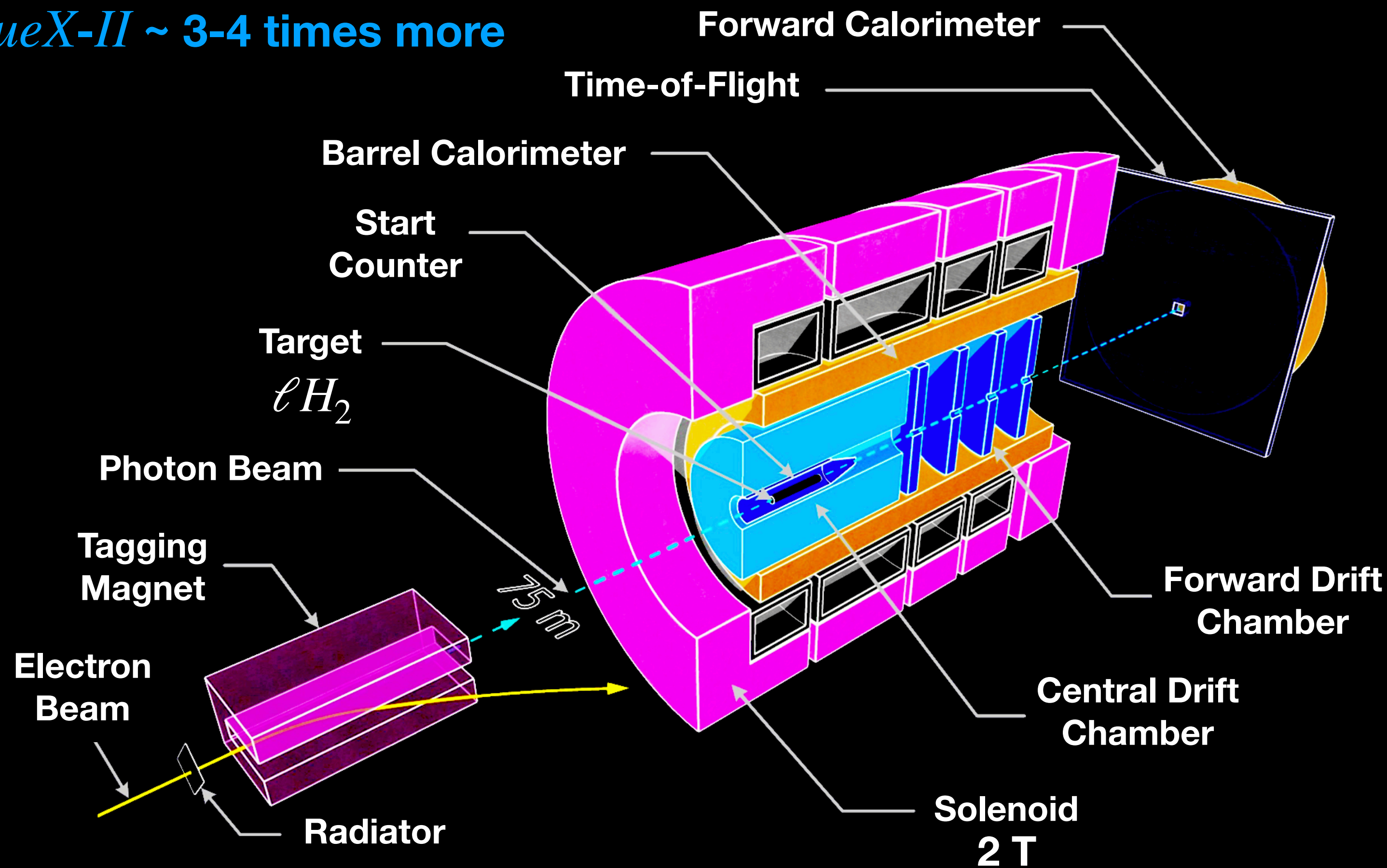
Use dominant $a_2(1320)$ as reference signal for the π_1 exotic hybrid search



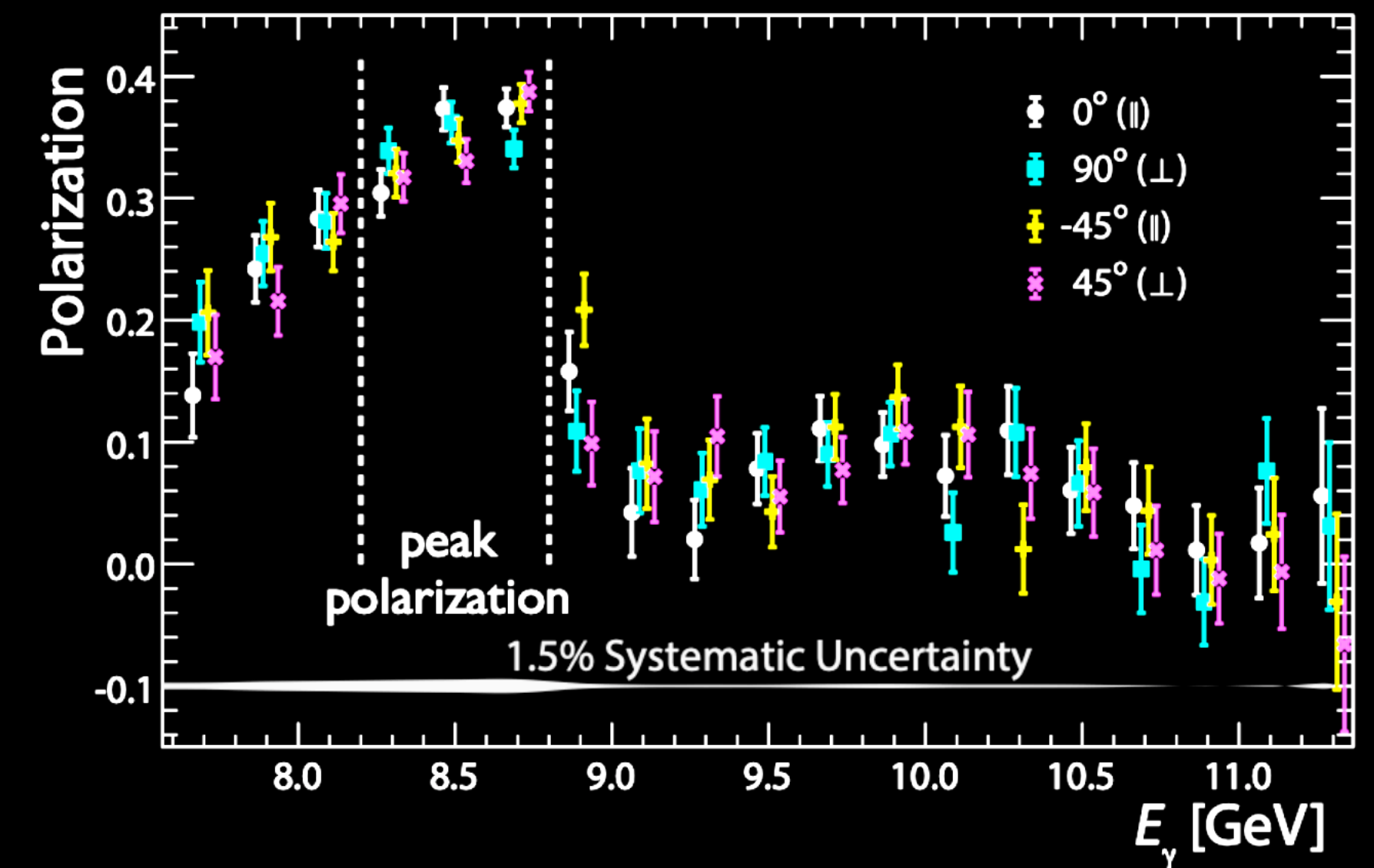
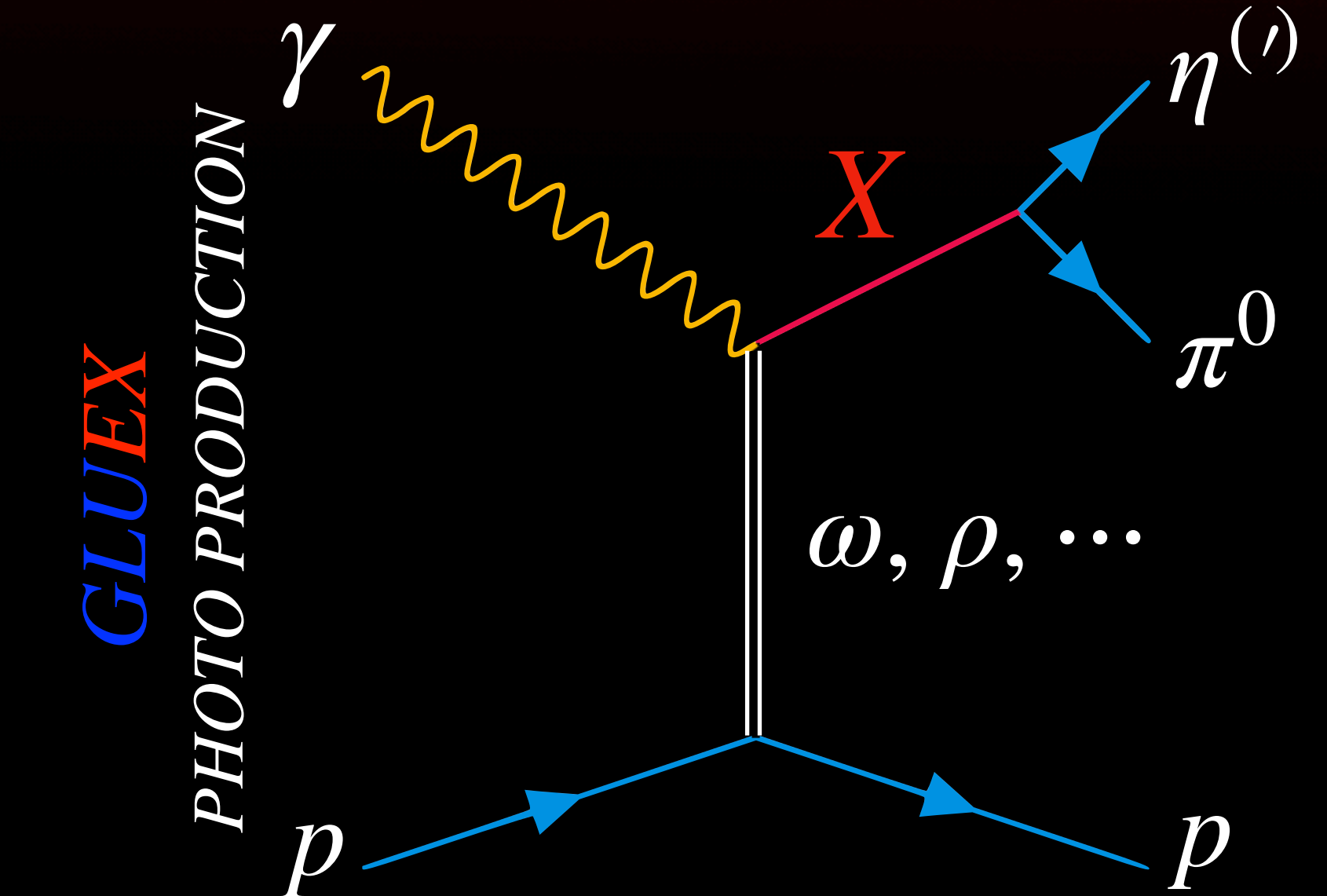
A. Rodas et al. [Joint Physics Analysis Center], PRL 122, 042002 (2019)

Experiment Details

- designed to reconstruct final state particles from $\gamma p \rightarrow pM$
- 4 polarization orientations
- *GlueX-I* collected $\int L = 125 \text{ pb}^{-1}$ in coherent peak
- *GlueX-II* ~ 3-4 times more



S. Adhikari et al., NIM A 987, 164807 (2021)

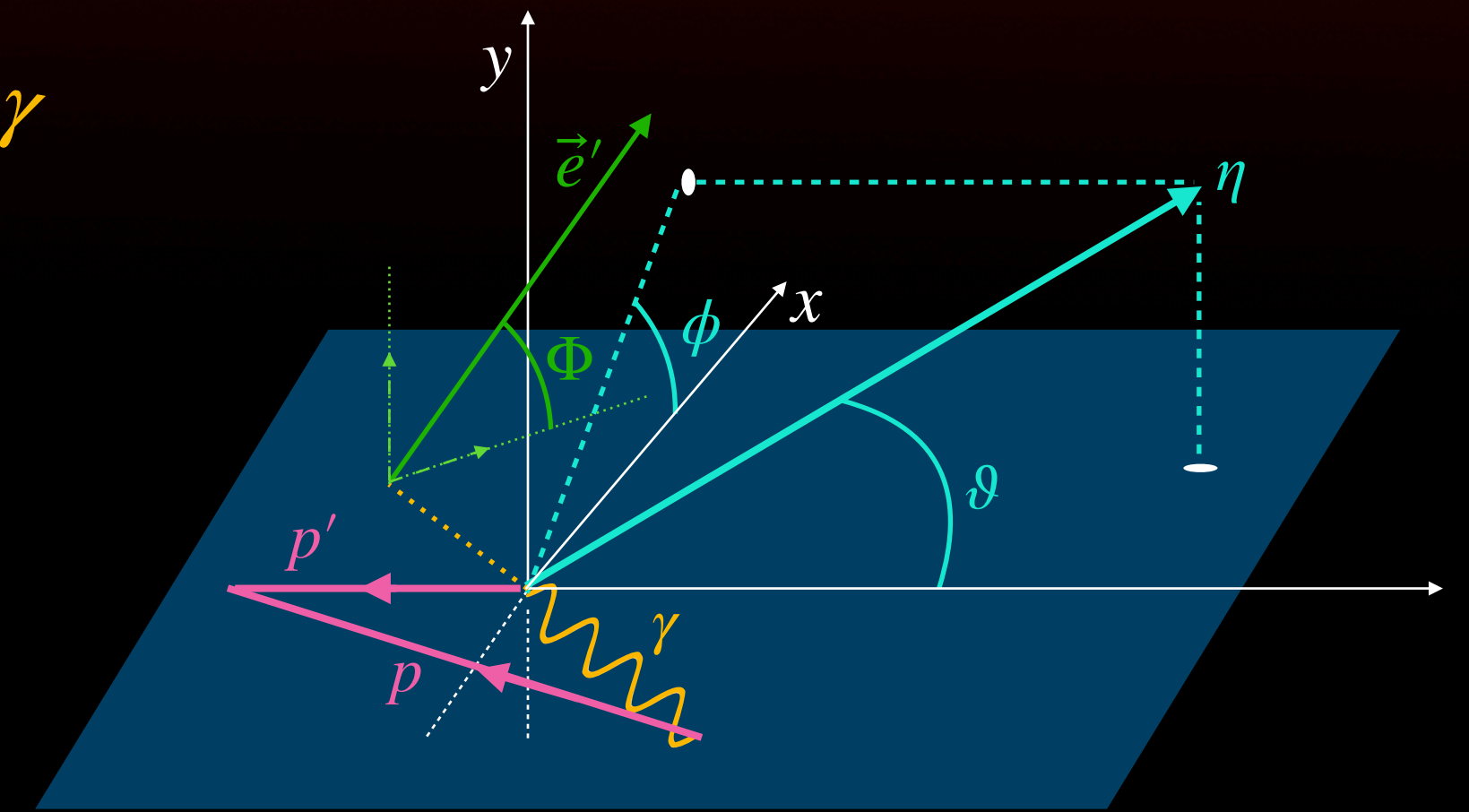


$$\mathcal{I}(\Omega, \Phi) = \sum_{\alpha=0}^3 \mathcal{I}_{\alpha}(\Omega) P_{\gamma}^{\alpha}(\Phi)$$

Linearly polarized γ

$$P_{\gamma}^{\alpha}(\Phi) = (1, -P_{\gamma} \cos(2\Phi), -P_{\gamma} \sin(2\Phi), 0)$$

$$\mathcal{I}(\Omega, \Phi) = \mathcal{I}_0(\Omega) - \mathcal{I}_1(\Omega) P_{\gamma} \cos(2\Phi) - \mathcal{I}_2(\Omega) P_{\gamma} \sin(2\Phi)$$

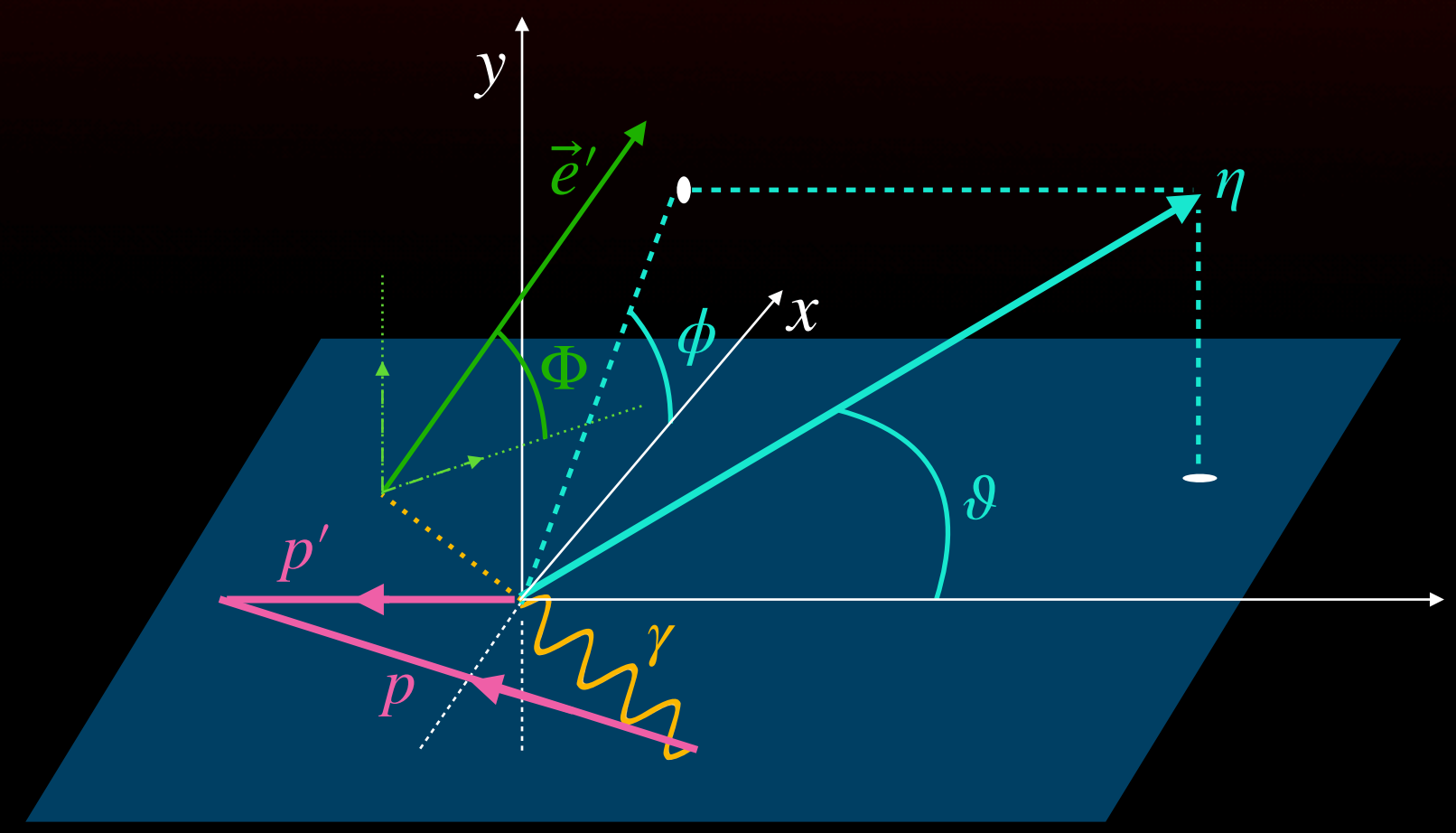


V. Mathieu et al. [JPAC], PRD 100, 054017 (2019)

- Describes all two-pseudoscalar systems (i.e. all $\eta^{(\prime)} \pi$)

- New basis $\rightarrow Z_l^m(\Omega, \Phi) = Y_l^m(\Omega) e^{-i\Phi}$

- Described by 3 angles: $\left[\begin{array}{l} \cos \vartheta_{\eta^{(\prime)}} \\ \phi_{\eta^{(\prime)}} \\ \Phi \end{array} \right]$ in the resonance frame
 \rightarrow btw the polarization and production plane

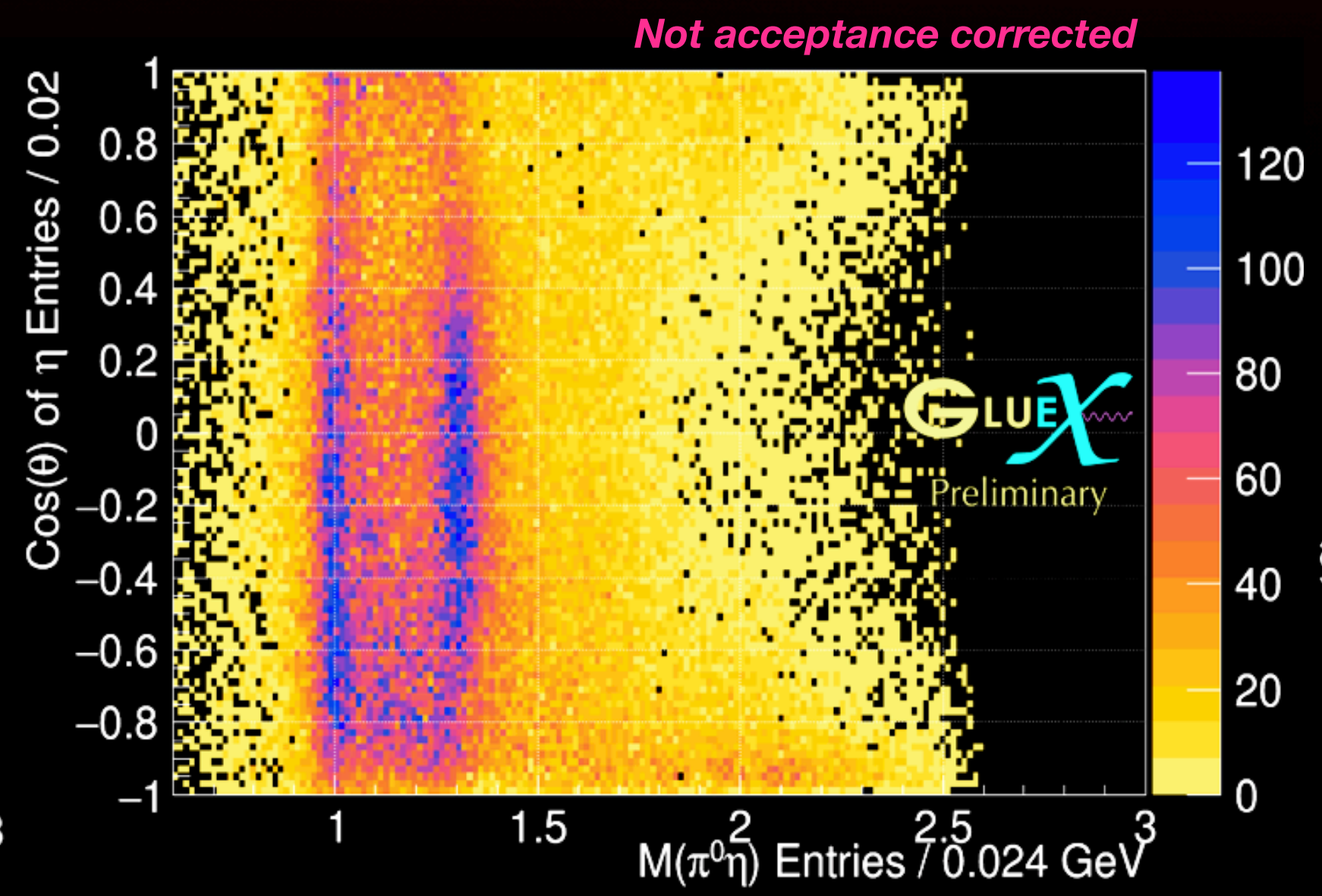
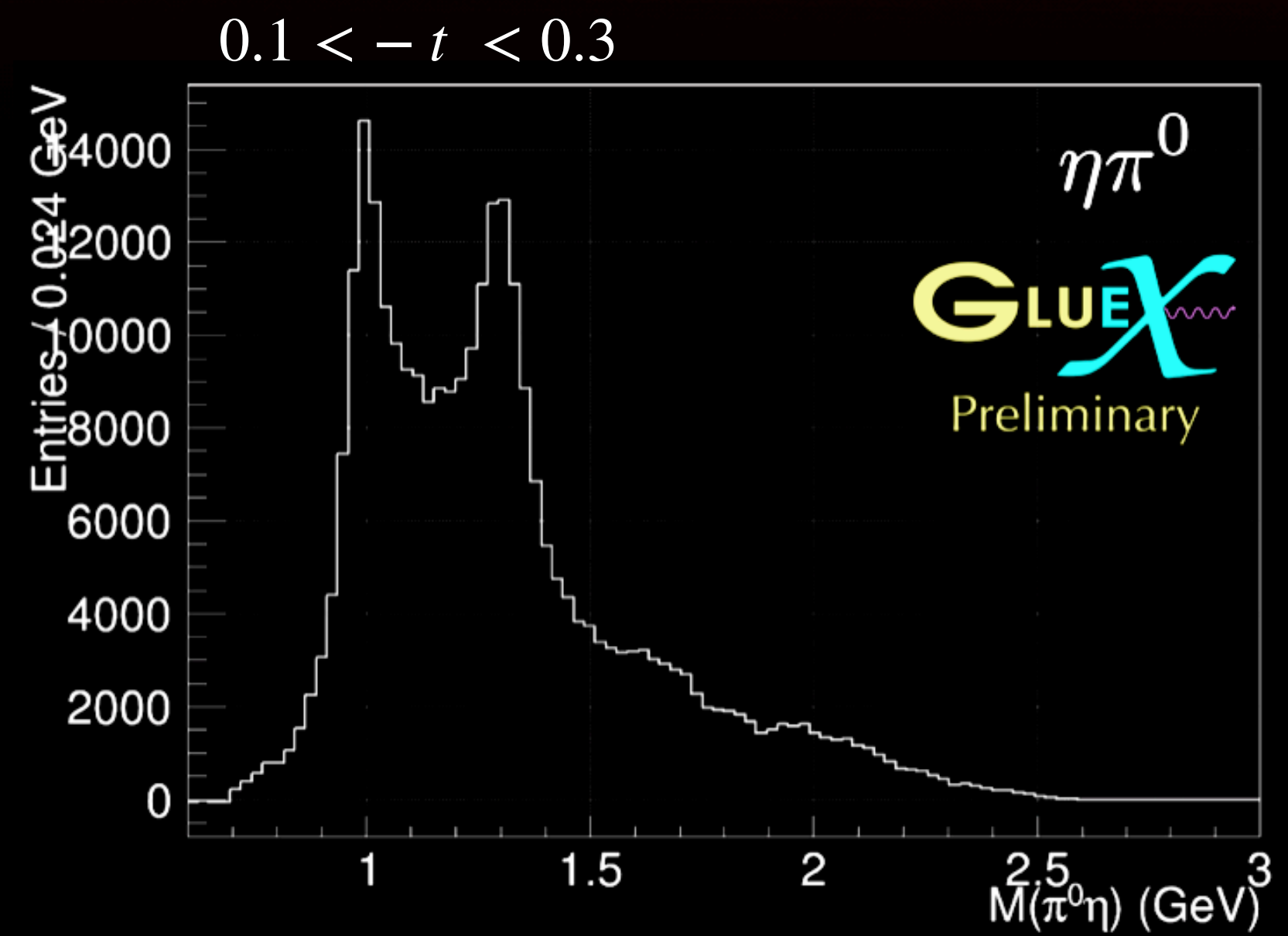


V. Mathieu et al. [JPAC], PRD 100, 054017 (2019)

- Reflectivity corresponds to exchange being natural (+1) and unnatural (-1) parity

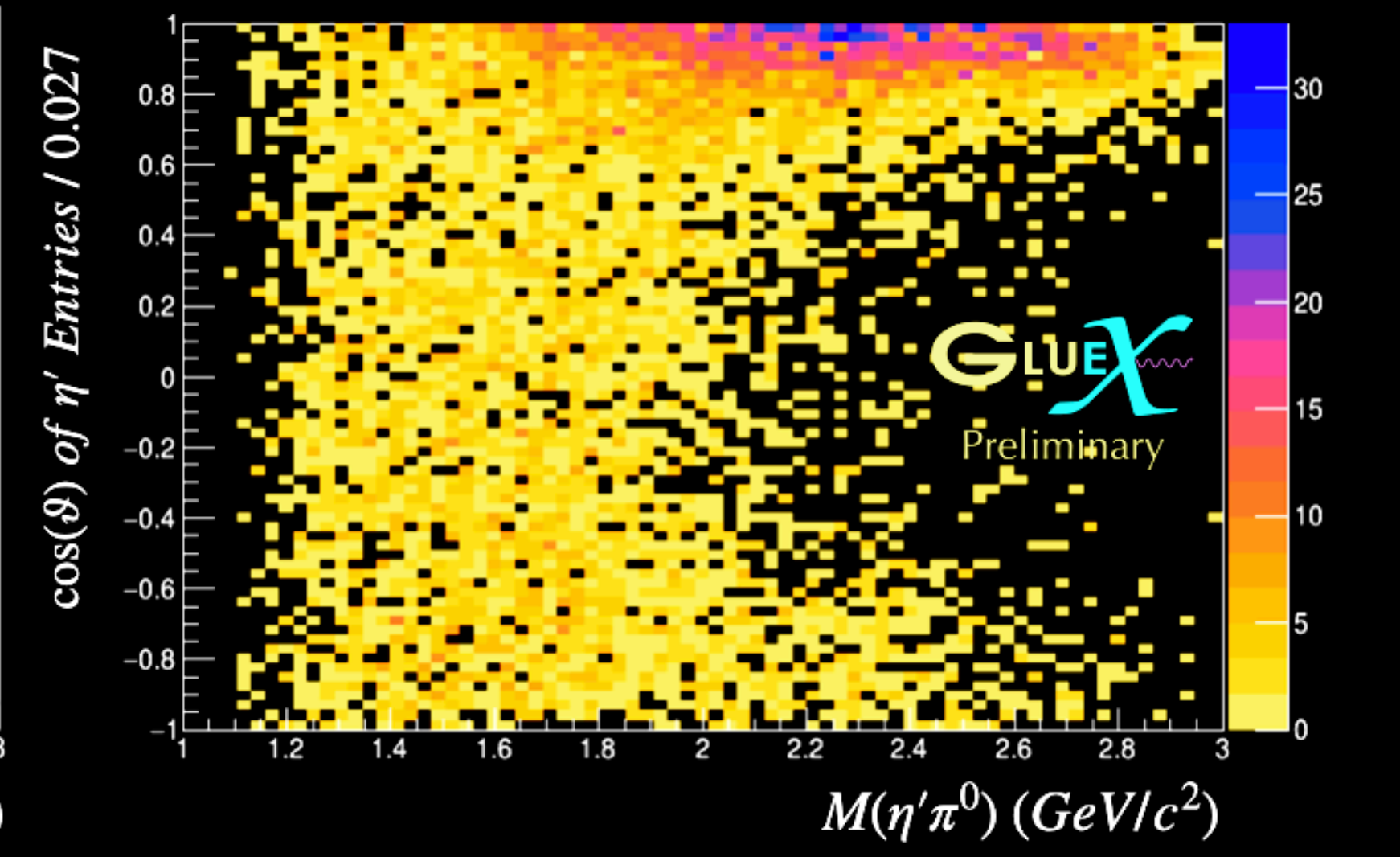
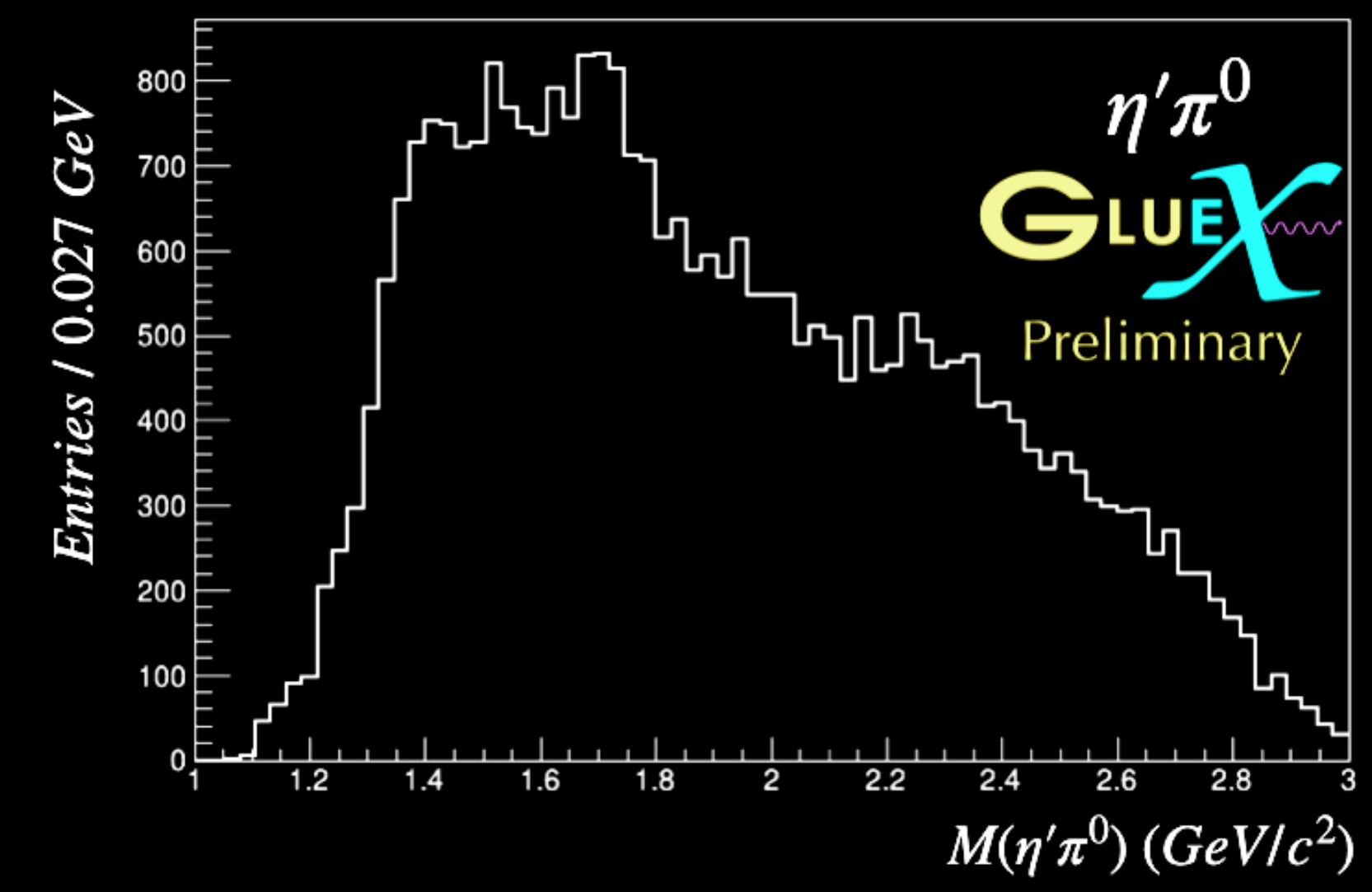
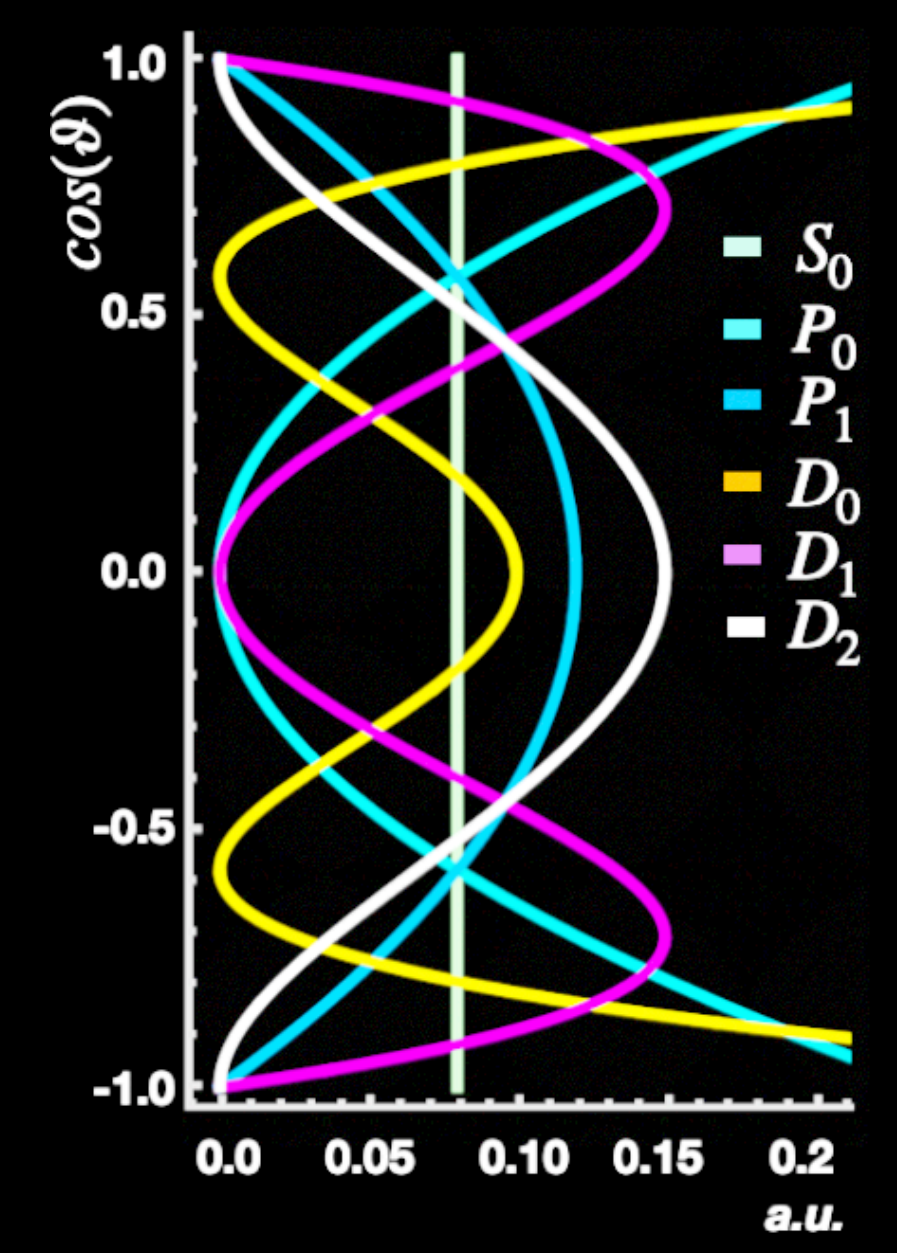
- 4x more amplitudes than hadro-production

$$\Rightarrow \mathcal{I}(\Omega, \Phi) = 2\kappa \sum_k \left\{ (1 - P_\gamma) \left| \sum_{l,m} [l]_m^{(-)} \mathcal{R}e[Z_l^m(\Omega, \Phi)] \right|^2 + (1 - P_\gamma) \left| \sum_{l,m} [l]_m^{(+)} \mathcal{I}m[Z_l^m(\Omega, \Phi)] \right|^2 + \right. \\ \left. (1 + P_\gamma) \left| \sum_{l,m} [l]_m^{(+)} \mathcal{R}e[Z_l^m(\Omega, \Phi)] \right|^2 + (1 + P_\gamma) \left| \sum_{l,m} [l]_m^{(-)} \mathcal{I}m[Z_l^m(\Omega, \Phi)] \right|^2 \right\}$$

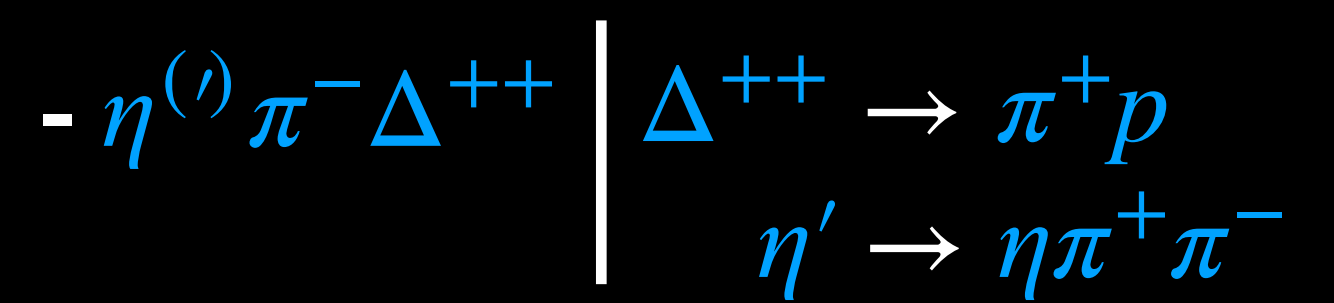


• Neutral decay modes

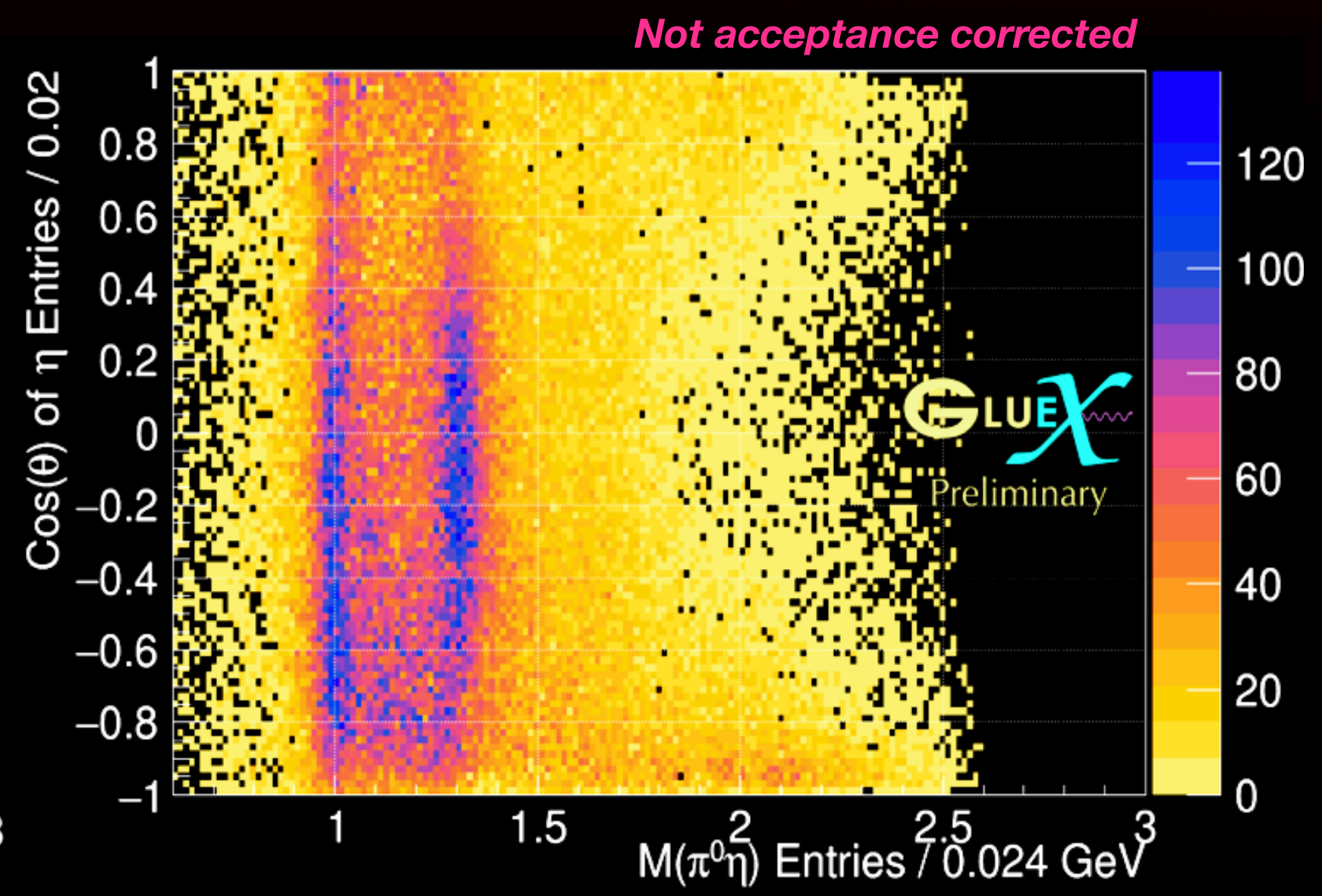
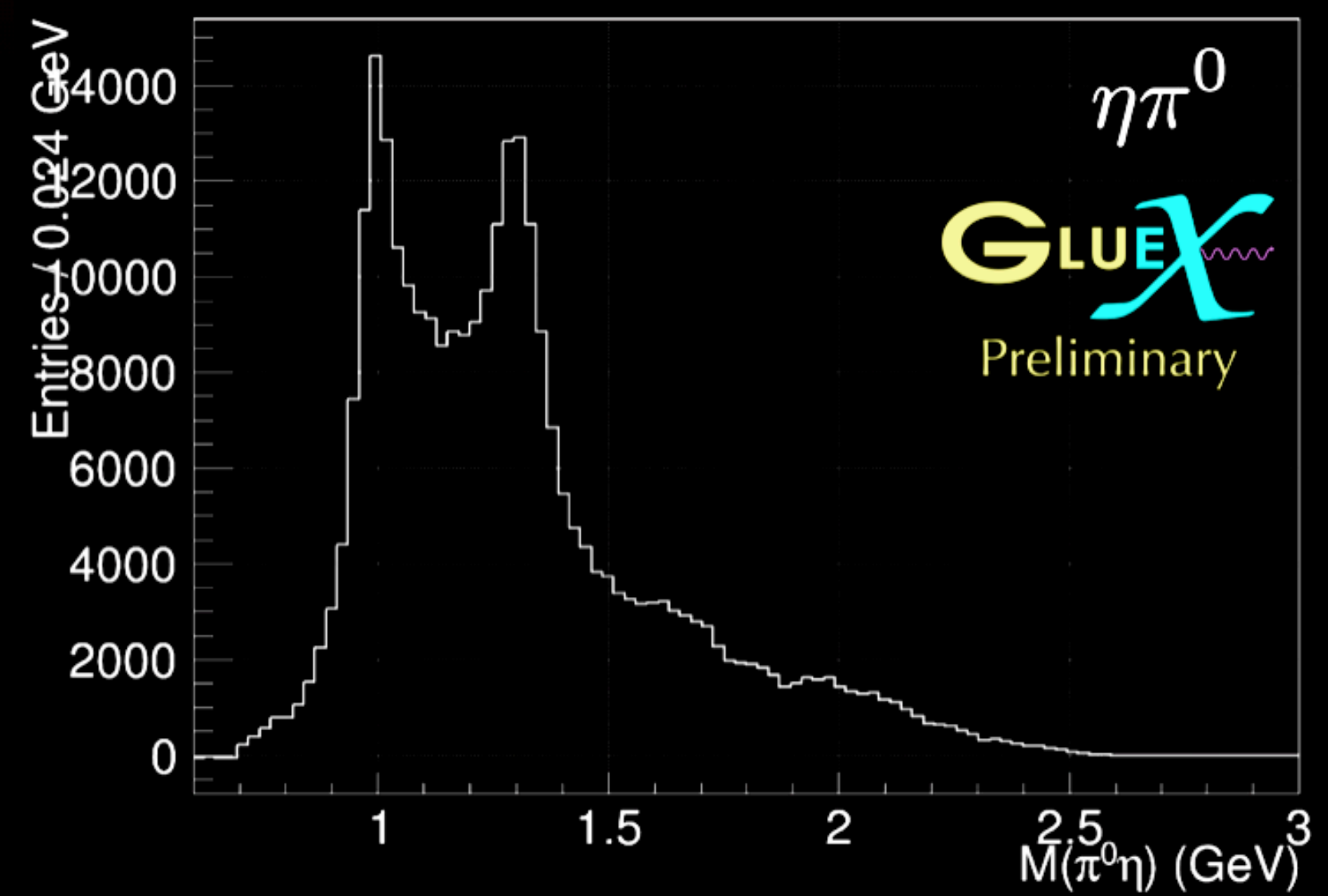
- $\eta\pi^0 \rightarrow 4\gamma$
- $\eta'\pi^0 \rightarrow 4\gamma\pi^+\pi^-$



• Charged decay modes also being analyzed ...

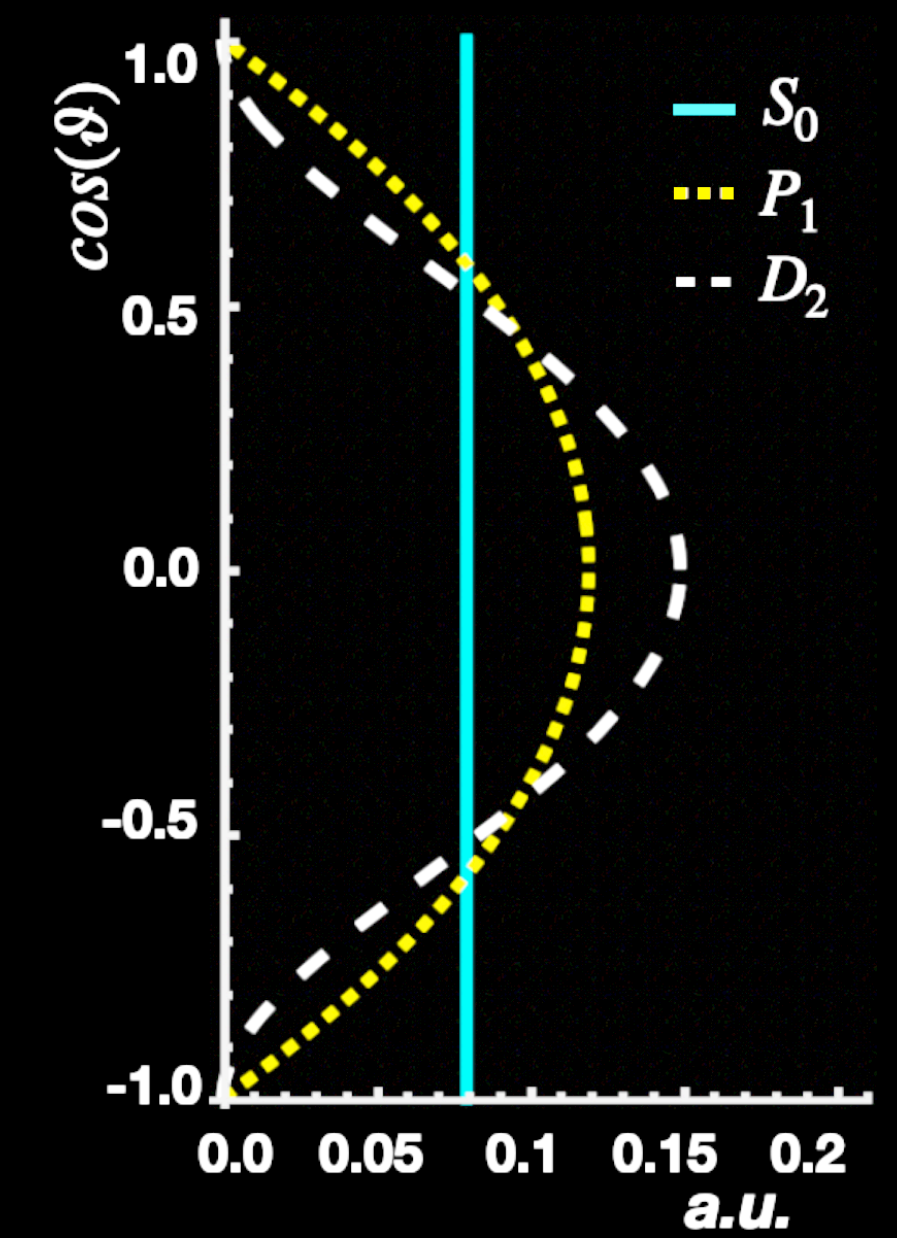


$0.1 < -t < 0.3$

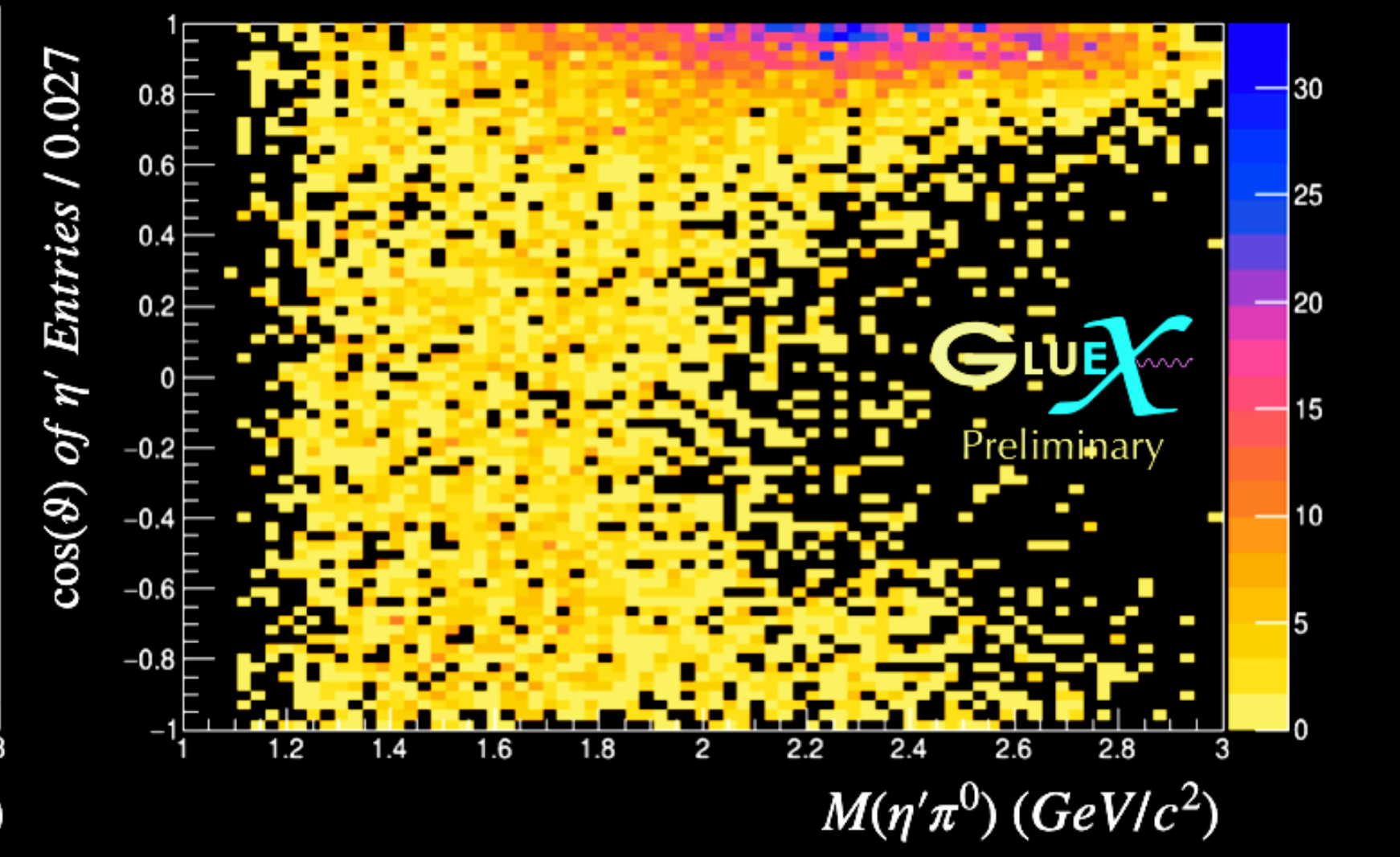
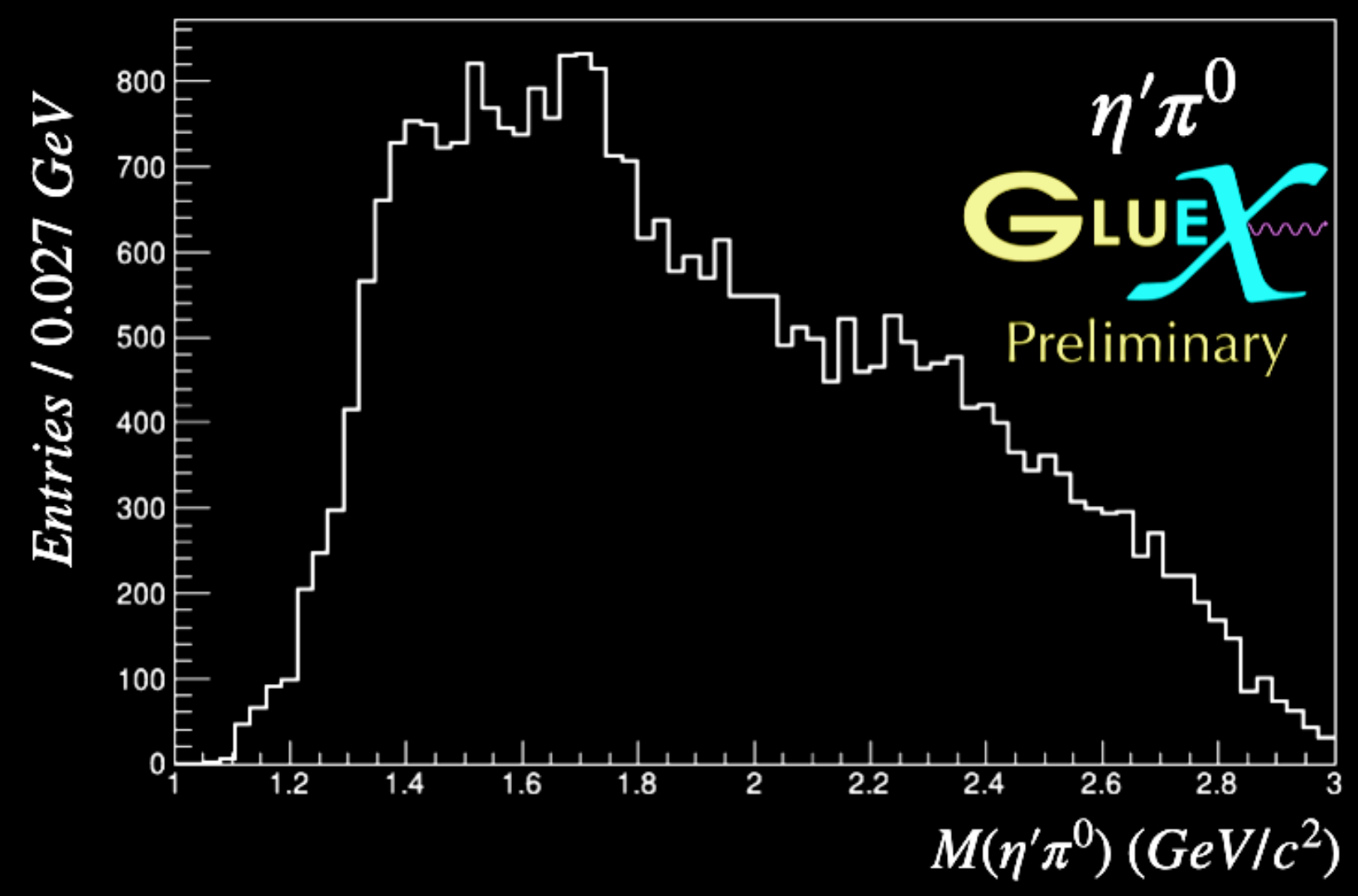


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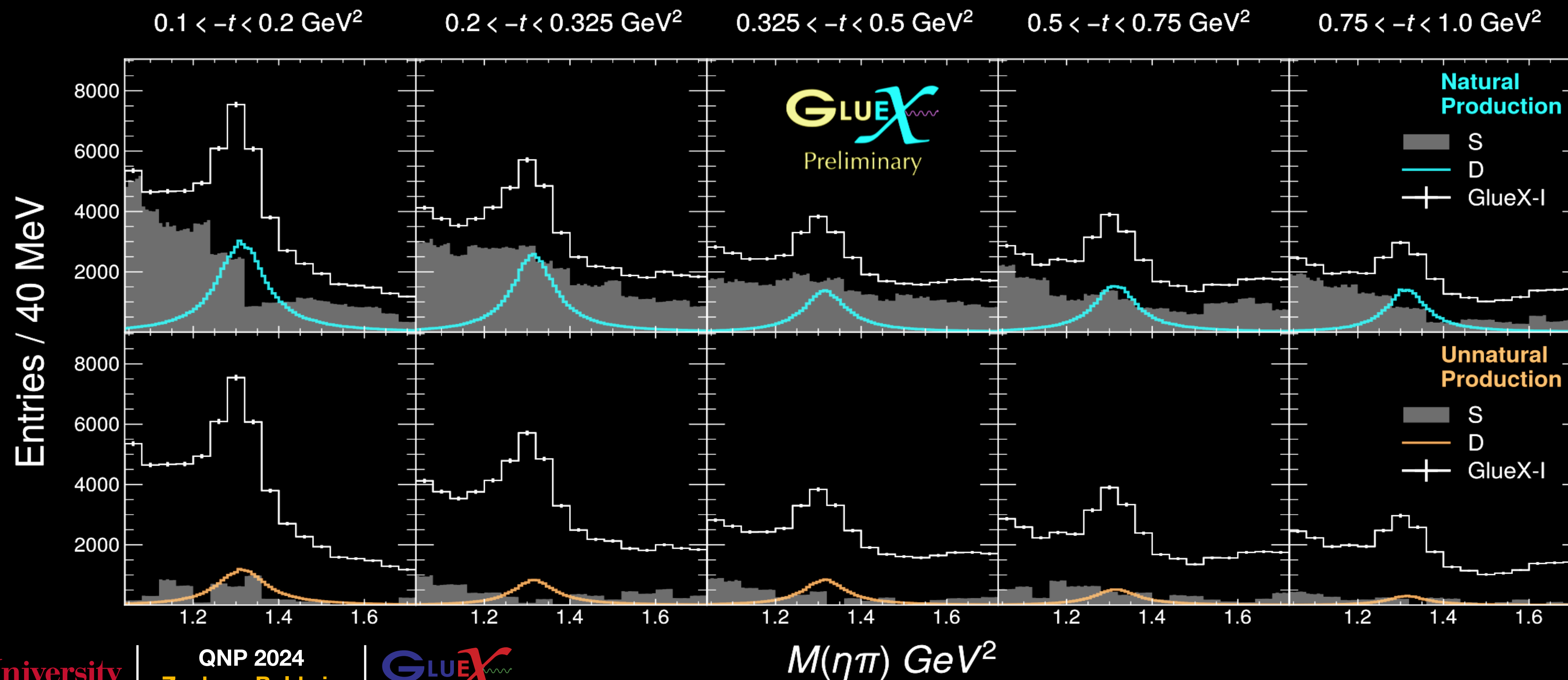
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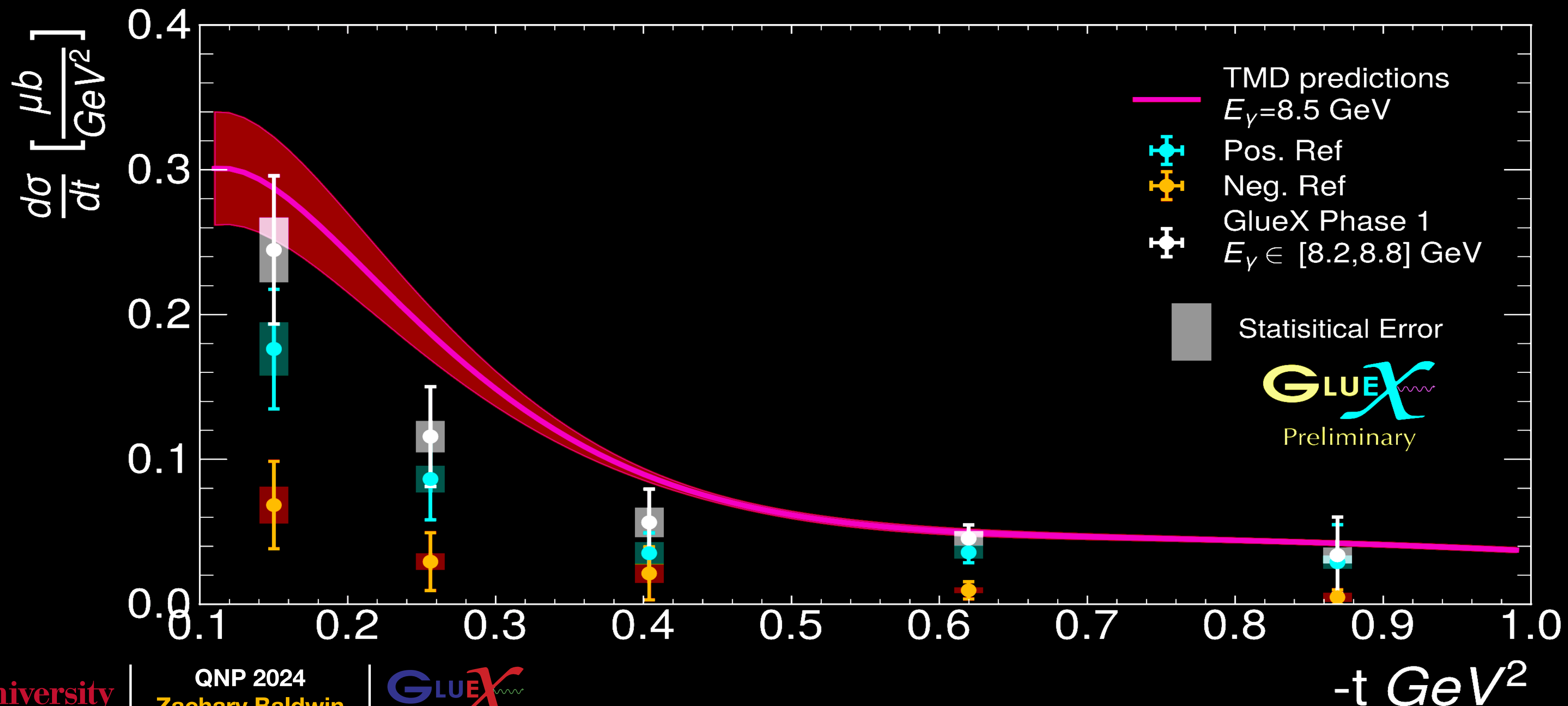
• Dominant S_0 and D_2 contributions



- Assume $a_2(1320)$ and $a_2(1700)$ are text book Breit-Wigner resonances
 - share only 1 common phase parameter for each in the D_{waves}
- S_{wave} contributions more complicated
 - define mass independent piecewise parameterization
- Individual fit results across $-t$
 - coherent sums of (+) and (-) reflectivities



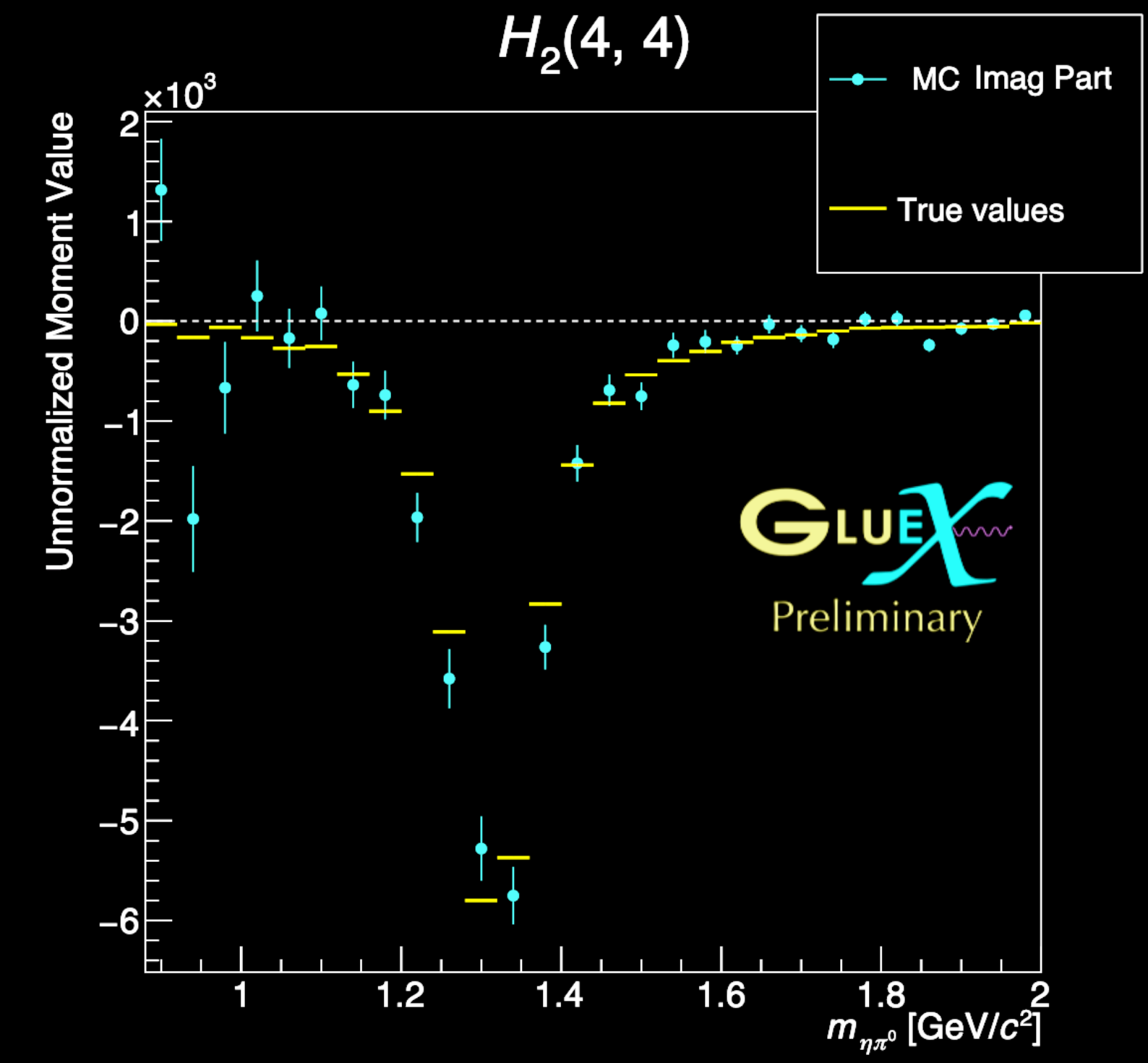
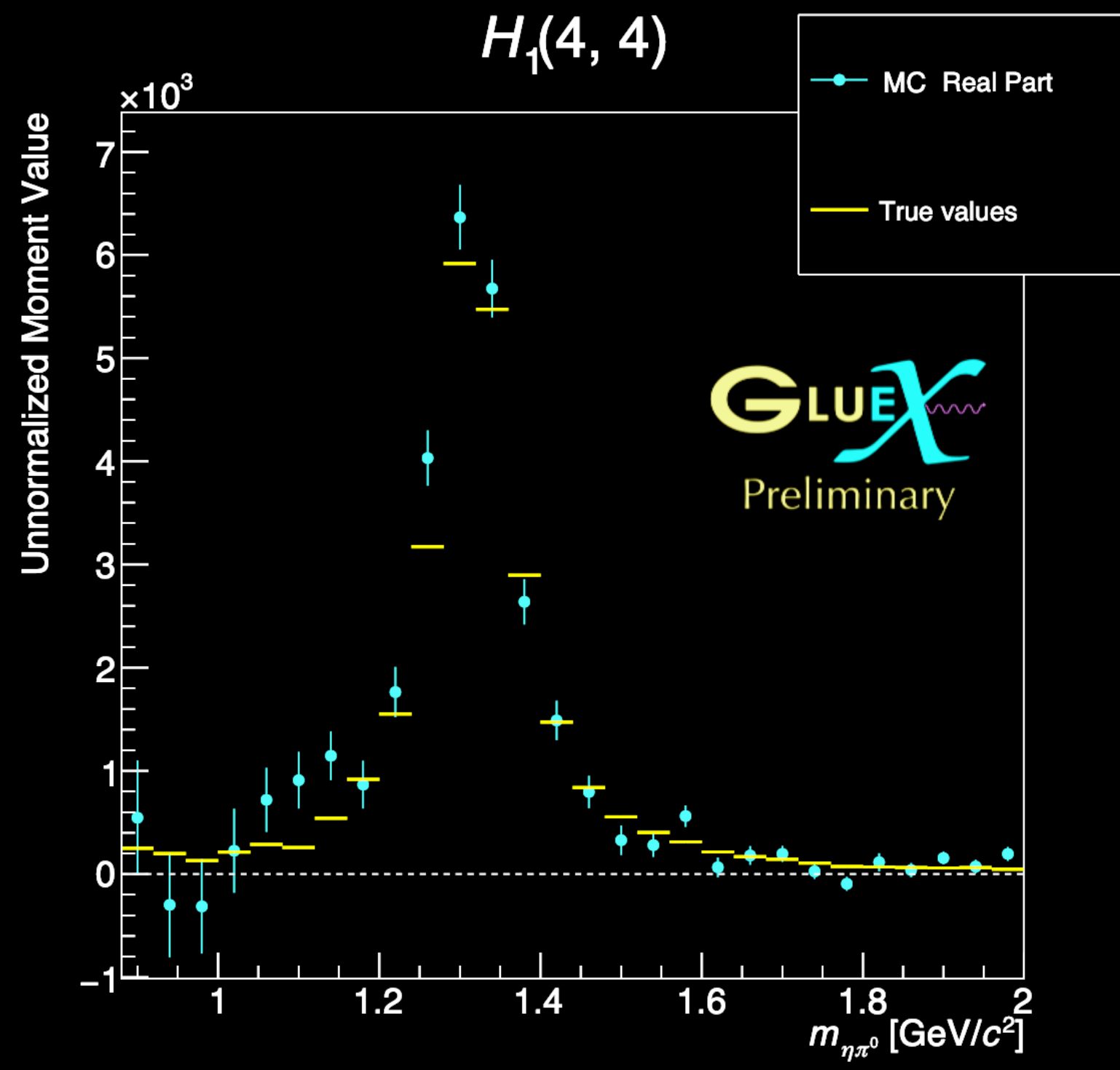
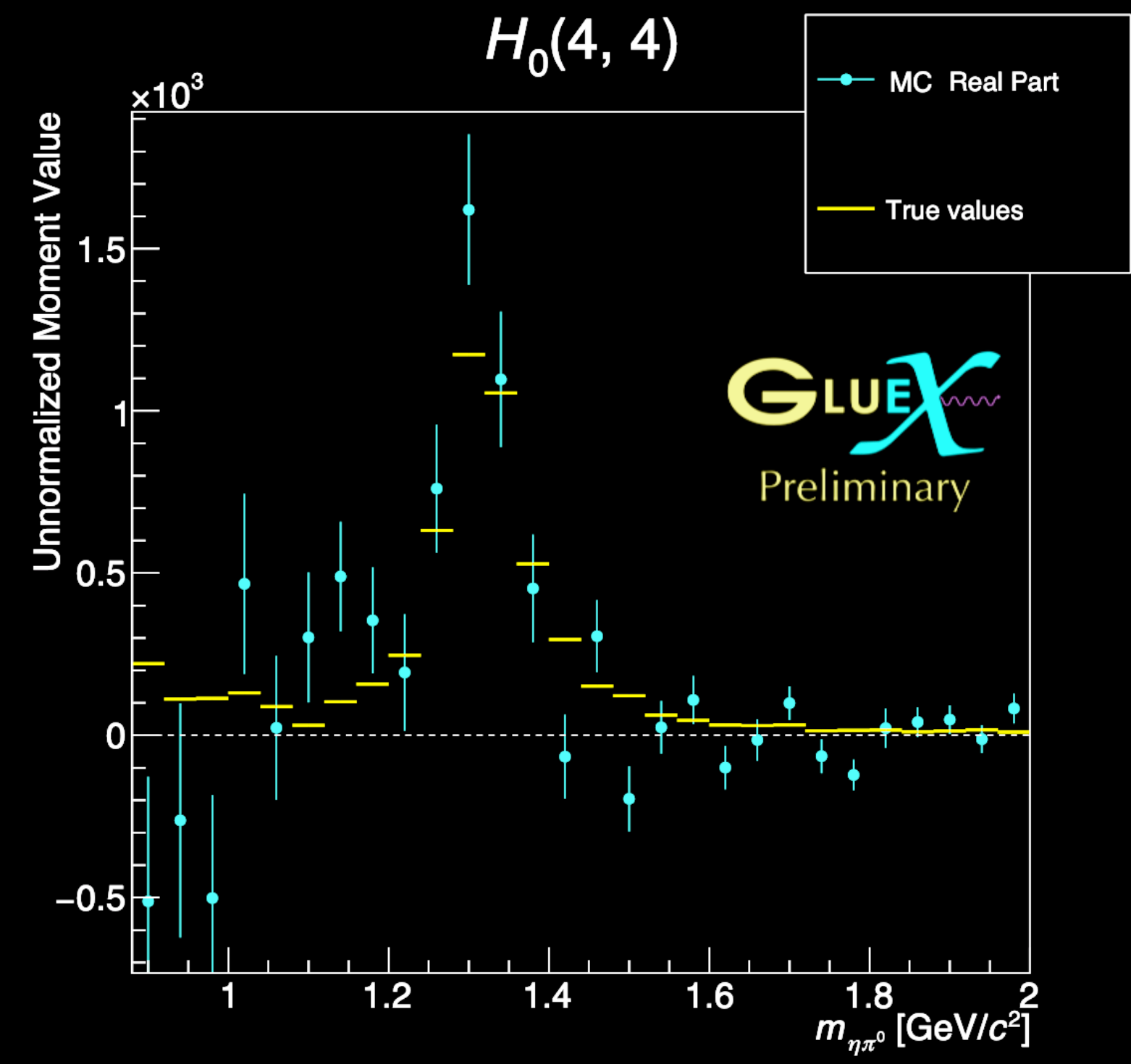
- **Assume $a_2(1320)$ and $a_2(1700)$ are text book Breit-Wigner resonances**
 - share only 1 common phase parameter for each in the D_{waves}
- S_{wave} contributions more complicated
 - define *mass independent* piecewise parameterization
- **Decent agreement between JPAC predictions**
 - systematics finalized — **Publication in the works!**



What can moments provide?

- **Decompose** angular components **into spherical harmonics**
 - determine sensitivity to exotic contributions
- **No direct access to partial wave amplitudes**
 - **BUT vice-versa** — can calculate moments from partial wave

$H_i(4,4)$ contains clear $a_2(1320)$ signal



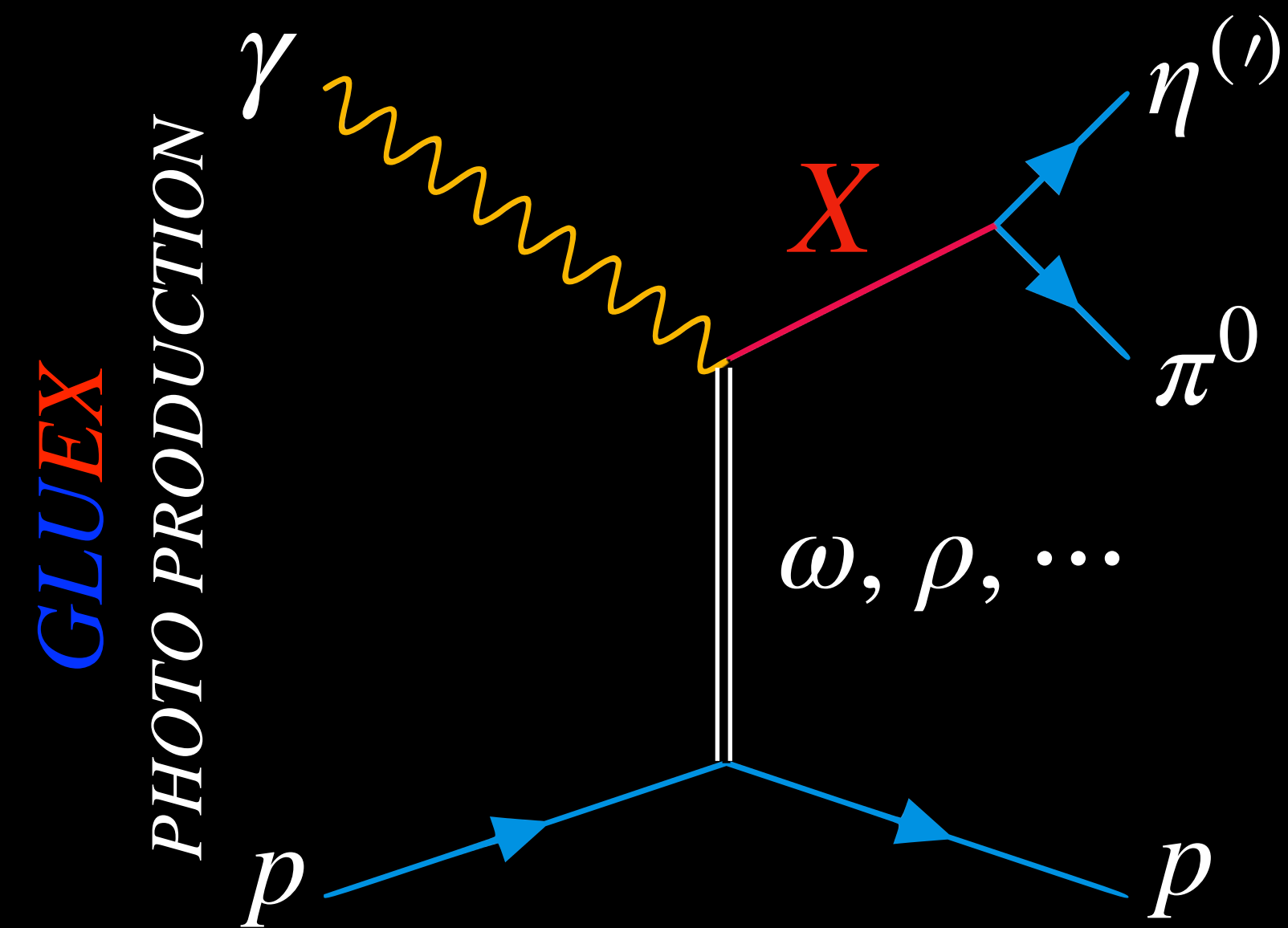
Monte Carlo Input-Output for $\eta^{(\prime)}\pi^0$
 — synthetic acceptance applied

Summary

- **GlueX has collected large quantity of photoproduced data**
 - strong effort to look for *exotic hybrid* π_1 meson in $\eta^{(\prime)}\pi$ systems using partial wave and moment analyses
 - recent results extract a_2 cross sections in multiple channels which will be used as a reference signal — **Publication in the works!**
 - can analyze production mechanisms using polarization info

Future Work

- **Further analyze both neutral and charged $\eta^{(\prime)}\pi$**
 - extract a_2^- cross section
 - perform cross checks with other decay modes
- **Continue strong collaboration with theorists**
 - work with JPAC on coupled channel fits



Exciting time for exotic hybrid search at GlueX

GlueX acknowledges the support of several funding agencies and computing facilities

gluex.org/thanks



BACKUP SLIDES

Cross Section definition

$$N_{produced} = \sigma \mathcal{L}$$

$$\Rightarrow \sigma = \frac{N_{obs}}{\mathcal{B}_{tot} \epsilon_{tot} \mathcal{L}}$$

$$\hookrightarrow = \epsilon_{tag} \times \epsilon_{non-tag}$$

$$\rightarrow = \mathcal{B}_{a_2 \rightarrow \pi^0 \eta} \times \mathcal{B}_{\eta \rightarrow \pi^0 \pi^+ \pi^-} \times \mathcal{B}_{\pi^0 \rightarrow \gamma \gamma}$$

(0.145) (0.2302) (0.9882)

$$= \left(\frac{1}{\mathcal{B}_{tot}} \right) \left(\frac{N_{obs}}{\epsilon_{non-tag}} \right) \left(\frac{1}{\epsilon_{tag} \mathcal{L}} \right) \quad \bullet \text{ Calculated to be } \approx 107.2 \text{ pb}^{-1}$$

- AmpTools calculation from “hybrid” fits

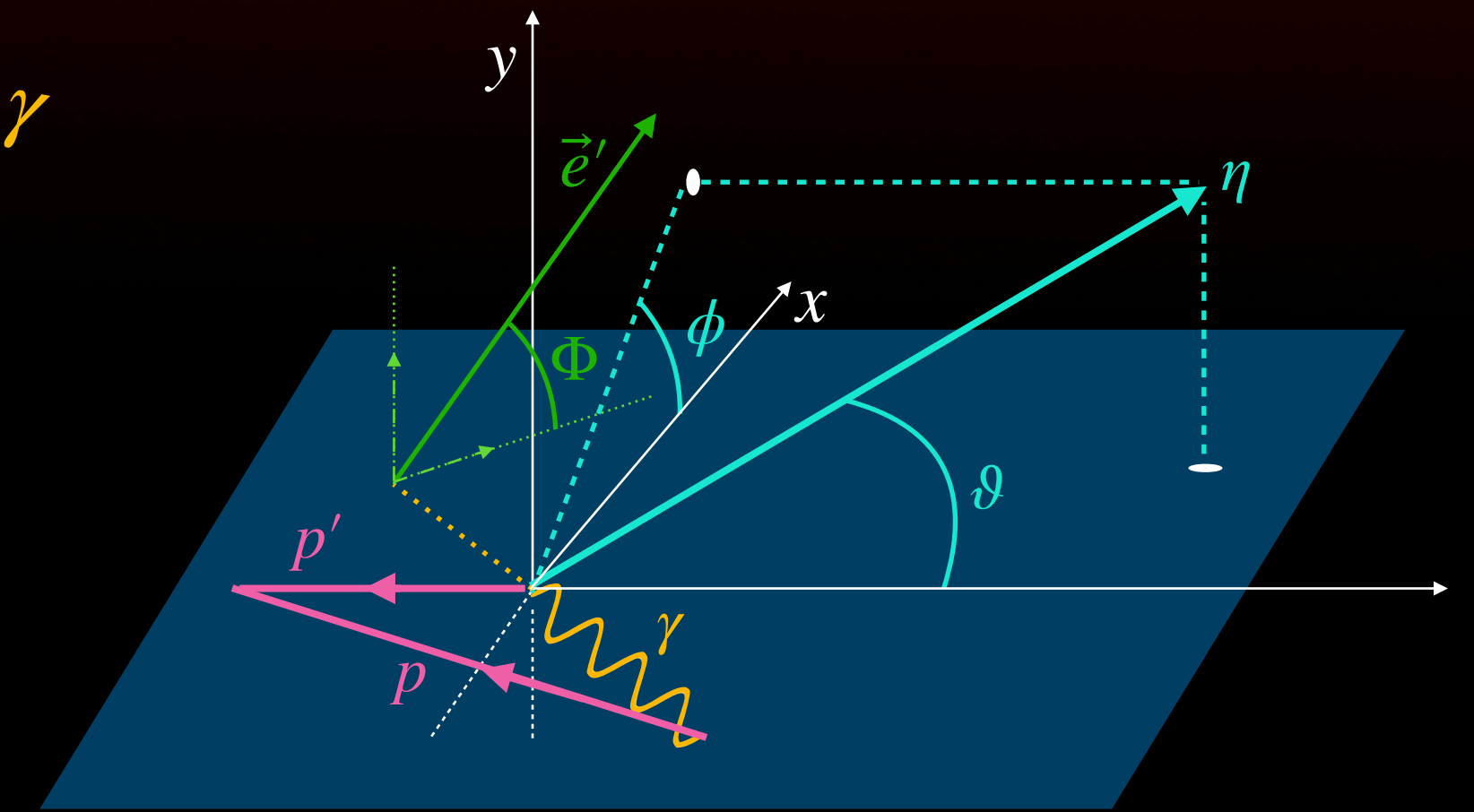
Moment definition

$$\mathcal{I}(\Omega, \Phi) = \sum_{\alpha=0}^3 \mathcal{I}_{\alpha}(\Omega) P_{\gamma}^{\alpha}(\Phi)$$

Linearly polarized γ

$$P_{\gamma}^{\alpha}(\Phi) = (1, -P_{\gamma} \cos(2\Phi), -P_{\gamma} \sin(2\Phi), 0)$$

$$\mathcal{I}(\Omega, \Phi) = \mathcal{I}_0(\Omega) - \mathcal{I}_1(\Omega) P_{\gamma} \cos(2\Phi) - \mathcal{I}_2(\Omega) P_{\gamma} \sin(2\Phi)$$



V. Mathieu et al. [JPAC], PRD 100, 054017 (2019)

Decompose intensity components into **spherical harmonics**

$$\mathcal{I}_0(\Omega) = \sum_{L,M}^{\infty} \sqrt{\frac{2L+1}{4\pi}} H_0(L, M) Y_l^m(\Omega)$$

$$\mathcal{I}_{1,2}(\Omega) = - \sum_{L,M}^{\infty} \sqrt{\frac{2L+1}{4\pi}} H_{1,2}(L, M) Y_l^m(\Omega)$$

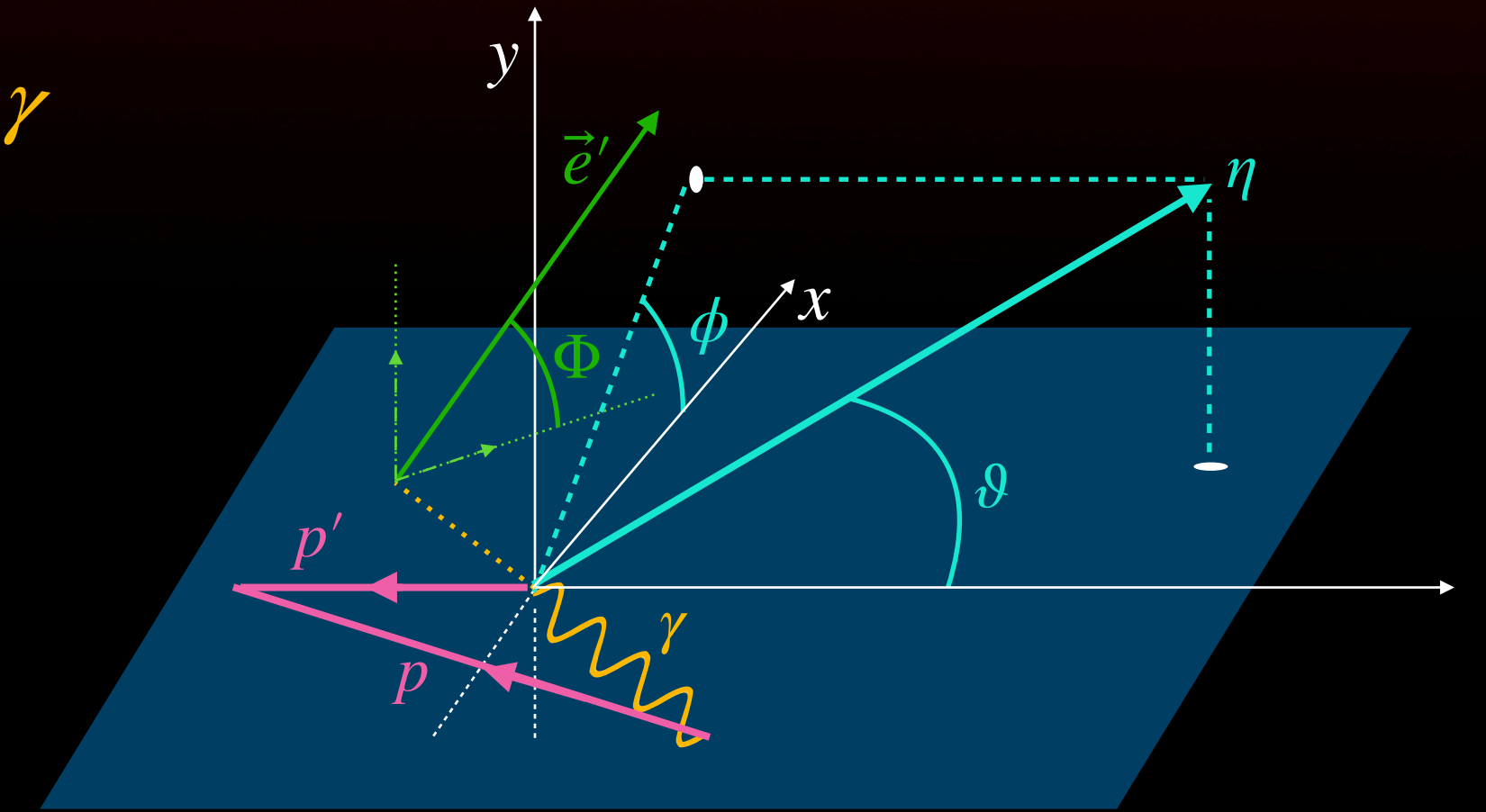
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V. Mathieu et al. [JPAC], PRD 100, 054017 (2019)

Corresponding moments:

$$\mathcal{I}_0(\Omega) \Rightarrow H_0(L, M) = \frac{1}{2\pi} \sqrt{\frac{4\pi}{2L+1}} \int_{4\pi} d\Omega \int_{-\pi}^{\pi} d\Phi \mathcal{I}(\Omega, \Phi) Y_L^{M*}(\Omega)$$

$$\mathcal{I}_{1,2}(\Omega) \Rightarrow H_{1,2}(L, M) = \frac{1}{\pi P_{\gamma}} \sqrt{\frac{4\pi}{2L+1}} \int_{4\pi} d\Omega \int_{-\pi}^{\pi} d\Phi \mathcal{I}(\Omega, \Phi) Y_L^{M*}(\Omega) \times \begin{cases} \cos(2\Phi) \\ \sin(2\Phi) \end{cases}$$