SUBA-Jet

A New Coherent Jet Energy Loss Model For Heavy Ion Collisions

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arXiv:2404.14579 [hep-ph]

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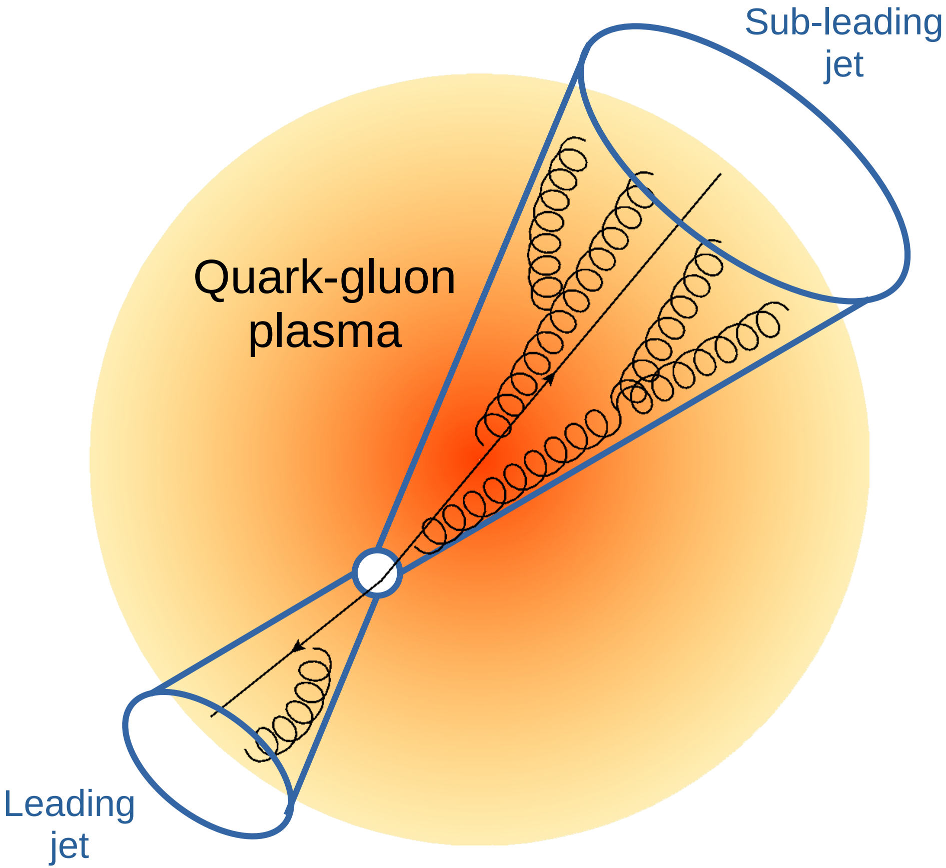
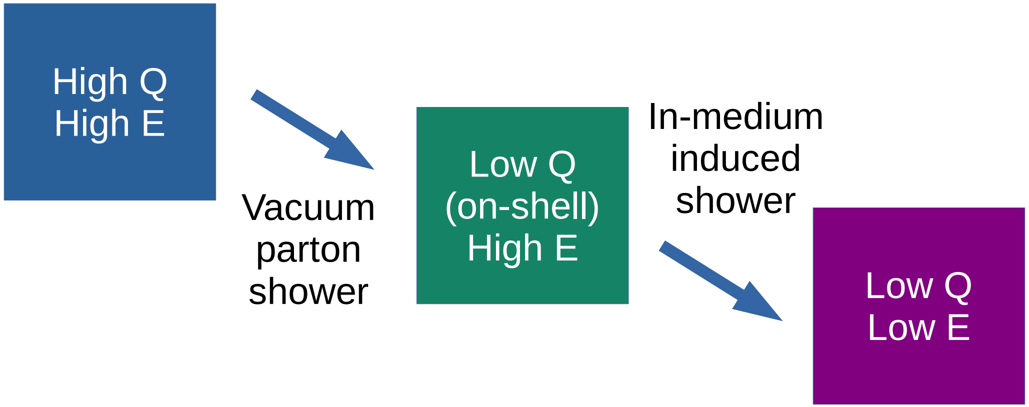
Jets in Heavy Ion Collisions



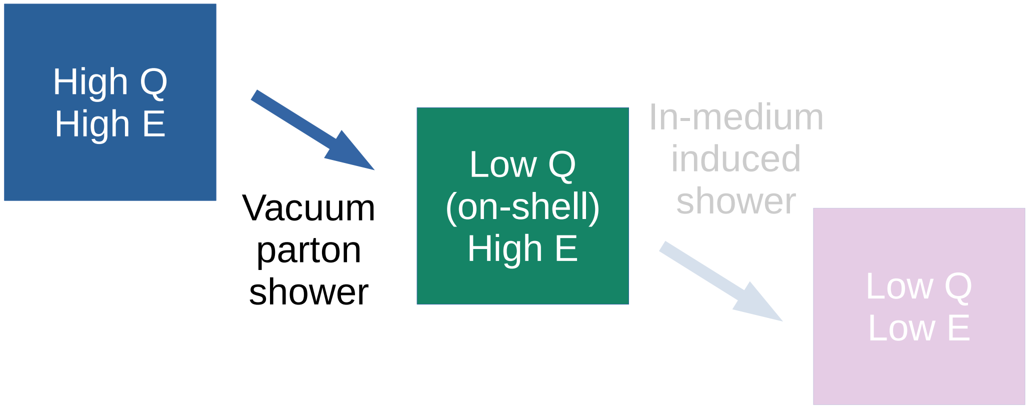
- Interactions between jet partons and the QGP medium leads to modifications of jet properties

  → **Jet Energy Loss / Quenching**

- **SUBA-Jet:**
  Monte Carlo for jet energy loss in heavy ion collisions

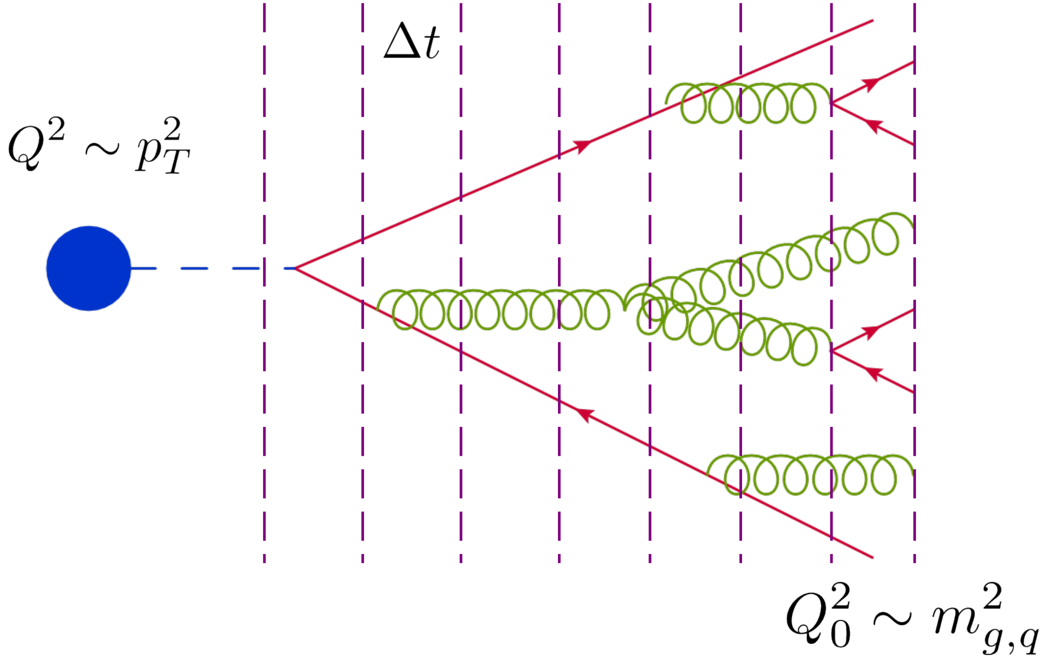


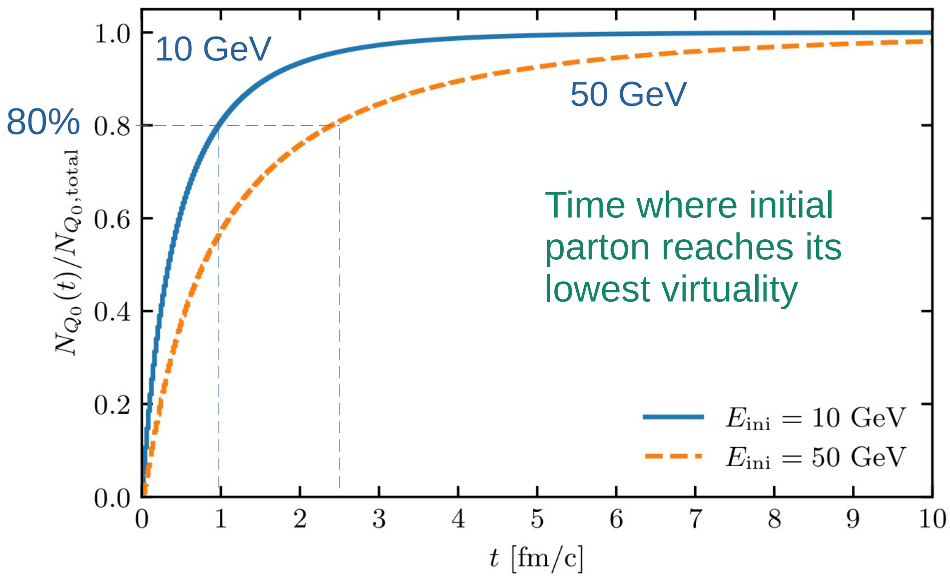
High Virtuality Regime



Vacuum Parton Shower

- Monte Carlo of a vacuum parton shower originally developed by Martin Rohrmoser
- Evolution according to the DGLAP equations from high virtuality $Q_{max} \sim p_T$ to low virtuality Q_0
- Time evolution split into time steps, mean life time $$\Delta t = \tau = \frac{E}{Q^2}$$

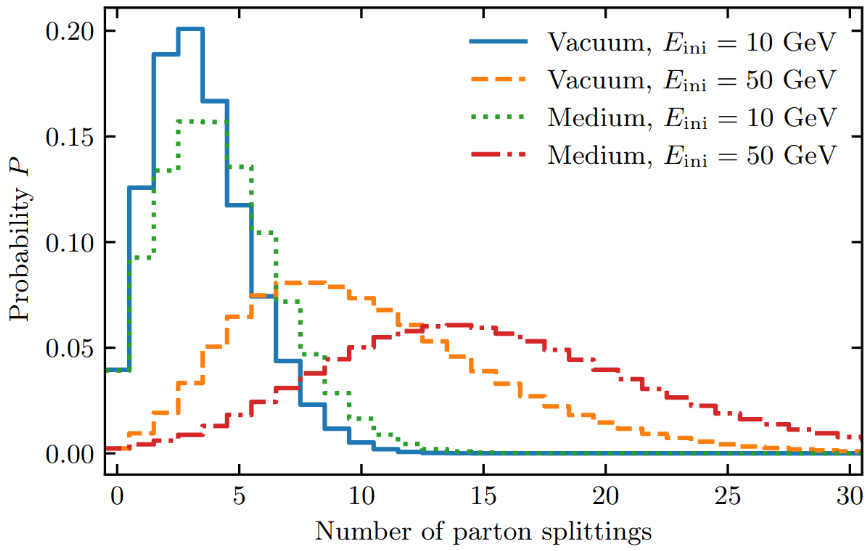


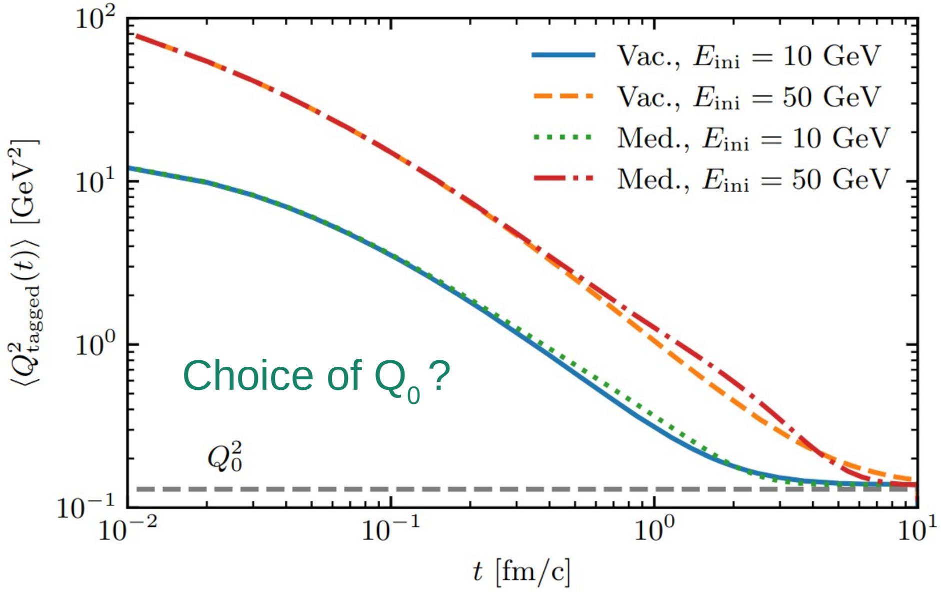


“Vacuum” Parton Shower in Medium

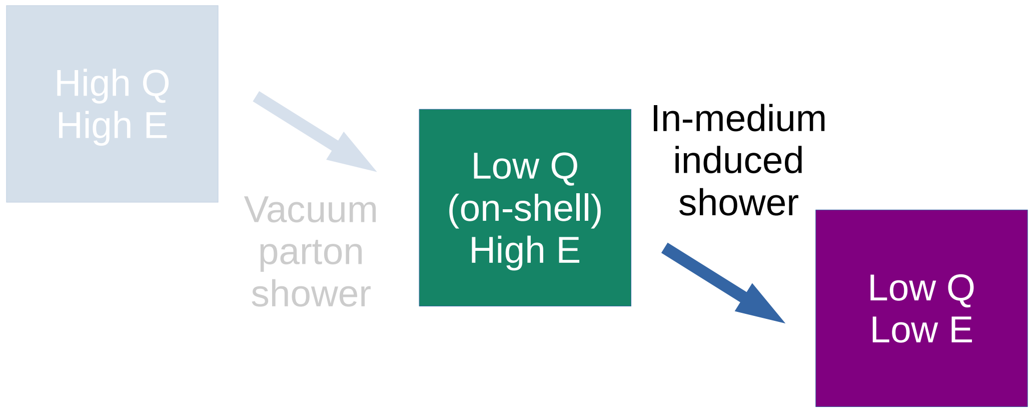
- Medium interactions for high Q regime resulting in virtuality increase, similar to YaJEM (T. Renk, 2008)

$$\frac{\mathrm{d}Q^2}{\mathrm{d}t} = \hat{q}(T)$$



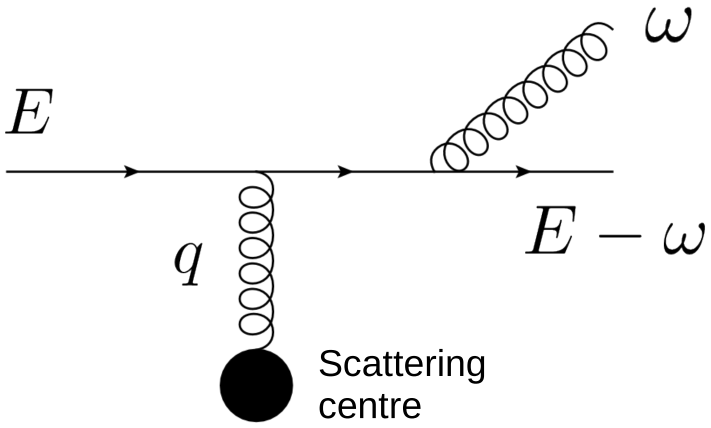


Low Virtuality Regime



Medium-Induced Single Radiation

- **Inelastic collision:**
  Single gluon emission from single medium scattering



- Original result from Gunion-Bertsch (1982)
  Generalised to massive case by Aichelin, Gossiaux, Gousset (2014)

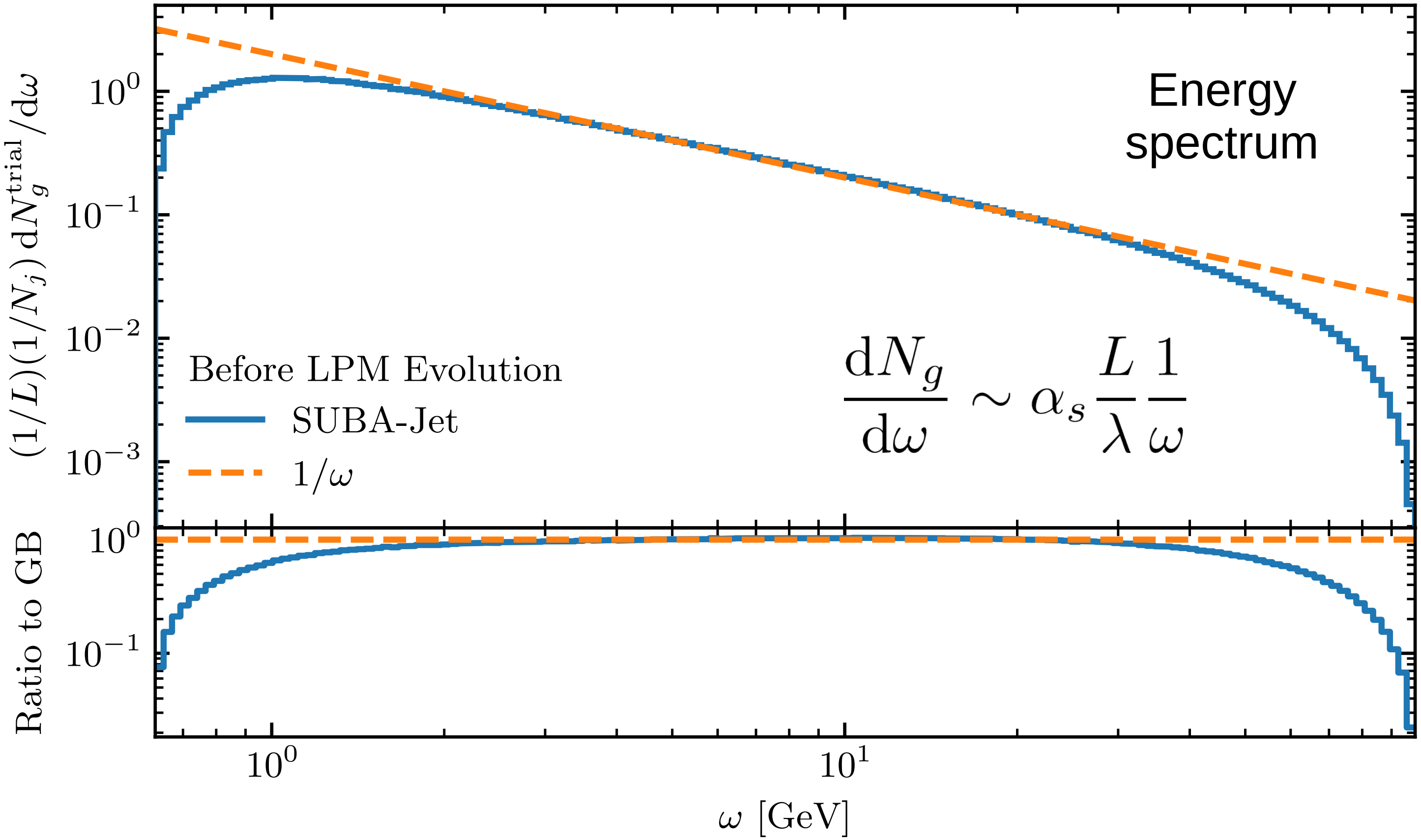
- Initial Gunion-Bertsch seed: i.e. radiation of a **preformed gluon** from a single scattering
  (Each parton can generate a number of preformed gluons)

- Gunion-Bertsch cross-section from scalar QCD

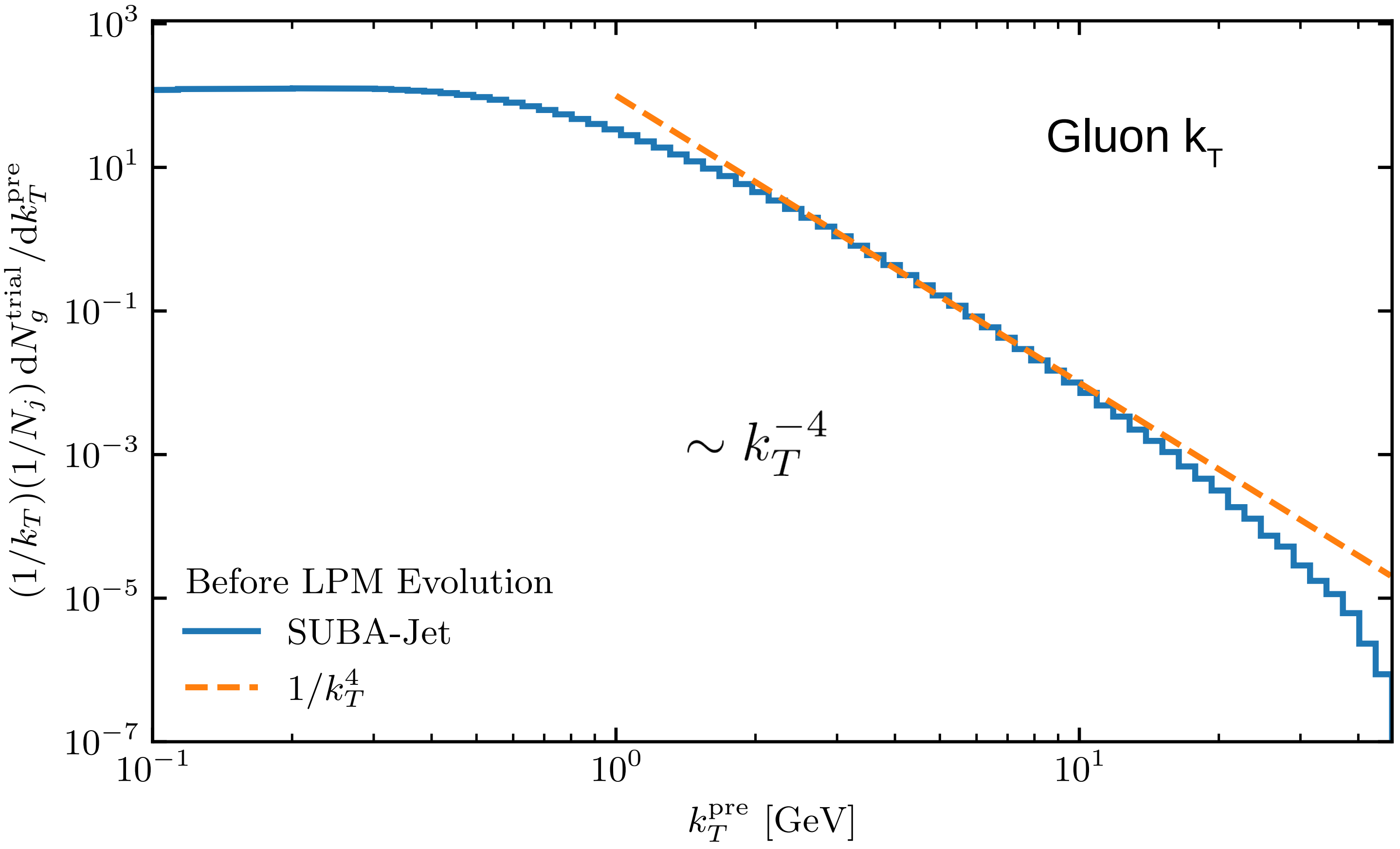
$$\frac{\mathrm{d}\sigma^{Qq\to Qqg}}{\mathrm{d}x\,\mathrm{d}^2k_T\,\mathrm{d}^2l_t} = \frac{\mathrm{d}\sigma_{\mathrm{el}}}{\mathrm{d}^2l_t} P_g(x, k_T, l_T)\theta(\Delta)$$

$$\frac{\mathrm{d}\sigma_{\mathrm{el}}}{\mathrm{d}^2l_t} \sim \frac{8\alpha_s^2}{9(l_T^2+\mu^2)^2}$$

Medium-Induced Single Radiation



Medium-Induced Single Radiation



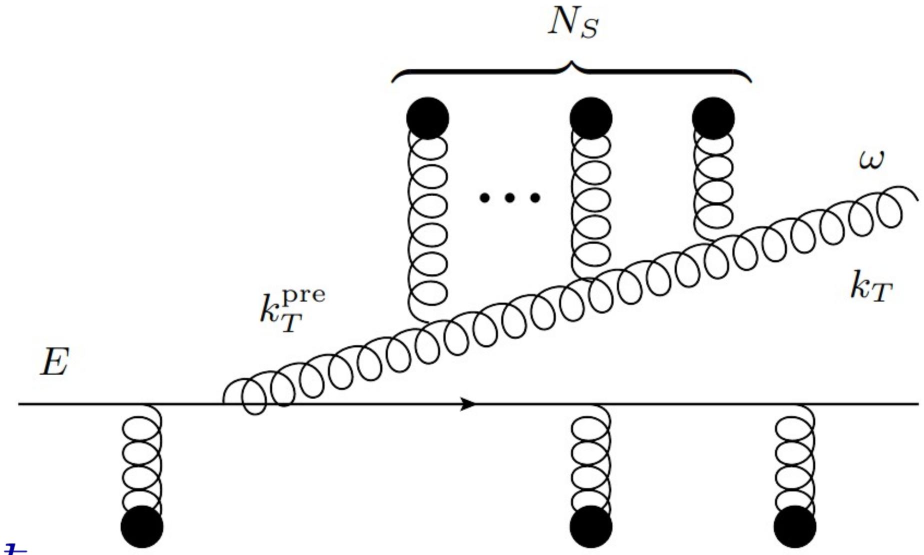
Coherency and the LPM Effect

$\omega =$ gluon energy

$\hat{q} =$ medium modifications

- The formation of the radiated gluon is a quantum mechanical process

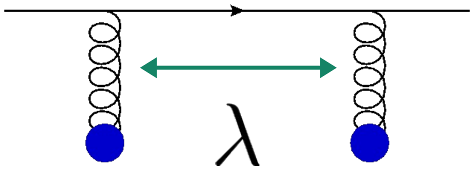
  Formation time: $t_f \sim \sqrt{\frac{\omega}{\hat{q}}}$



- Coherence effects:
  Landau-Pomeranchuk-Migdal (LPM) effect

- Have to take into account multiple scatterings with the medium during the formation time

$$N_s = \frac{t_f}{\lambda}$$



$$\lambda \simeq \frac{\hbar c}{\alpha_s T}$$

$L =$ path length of medium

Implementation of the LPM Effect

- At each timestep:
  - Elastic scattering with prob. $\Gamma_{\mathrm{el}}\Delta t$

$$\Gamma^q_{\mathrm{el}} = \left(1 + \frac{N_f}{N}\right) \frac{(N^2-1)T^3}{\pi\hbar c} \frac{4\alpha_s^2}{\mu^2}$$

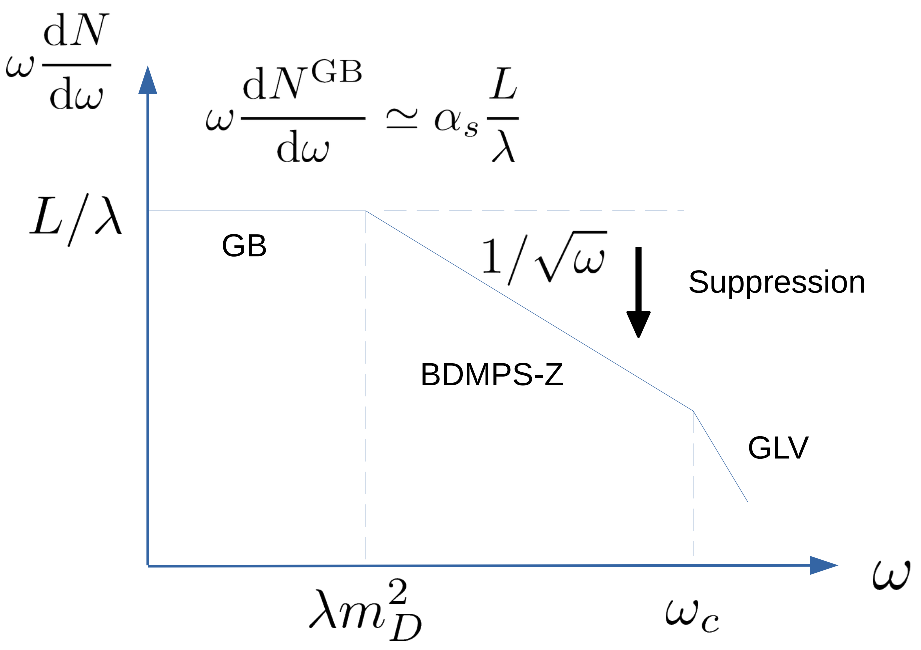
  - Radiation of preformed gluon with prob. $\Gamma_{\mathrm{inel}}\Delta t$

- BDMPS-Z spectrum at intermediate energies achieved by suppressing GB seed by

$$1/N_s$$

Like in Zapp, Stachel, Wiedemann, JHEP 07 (2011), 118

Radiation energy spectrum:



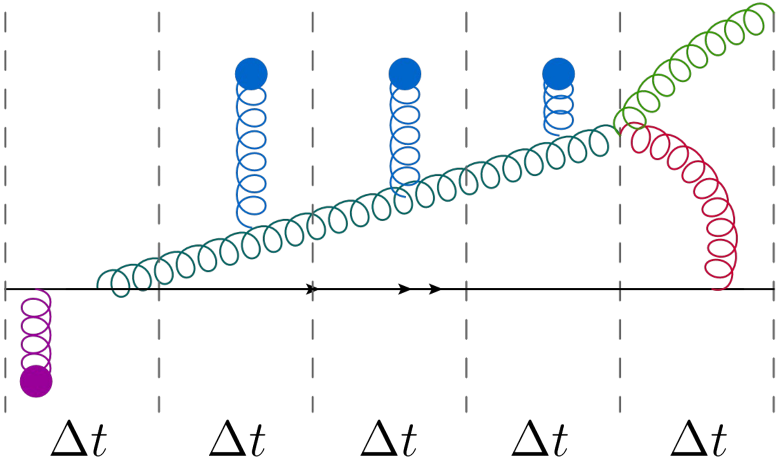
$$\omega \frac{\mathrm{d}N^{\mathrm{BDMPS\text{-}Z}}}{\mathrm{d}\omega} \simeq \alpha_s \sqrt{\frac{\hat{q}L^2}{\omega}}$$

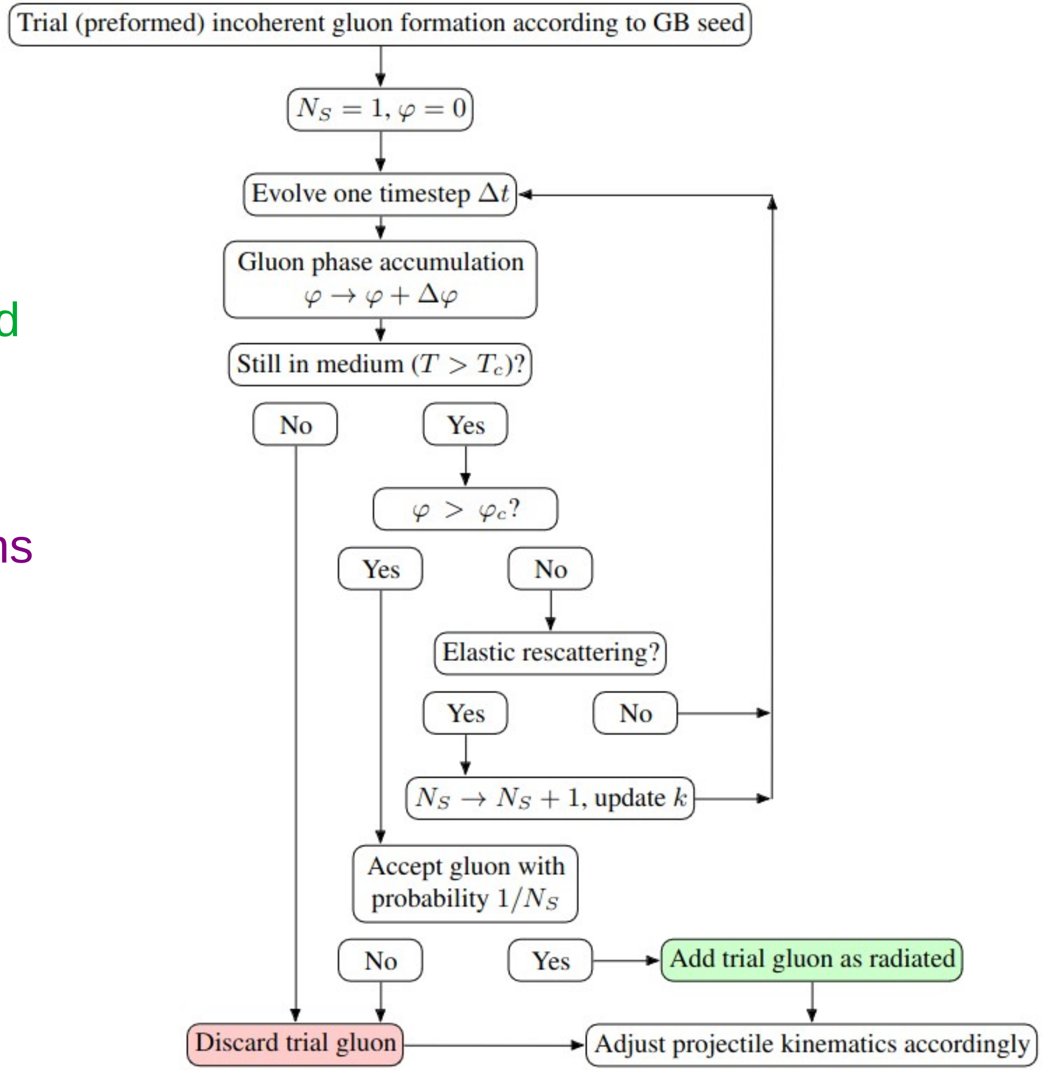
The Algorithm

Flow diagram:

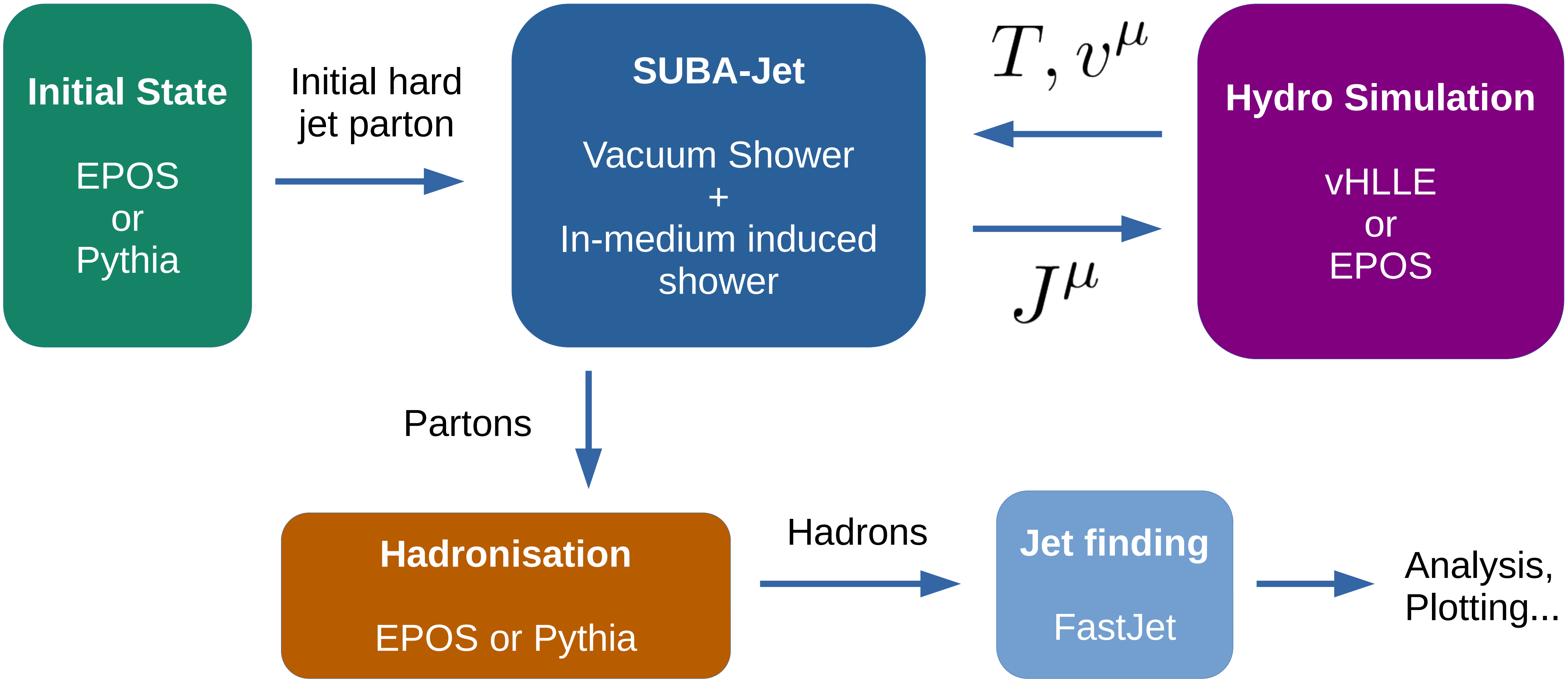
Algorithm for the coherent medium-induced gluon radiation in our model

Various parameters and settings can be changed and tuned to compare distributions

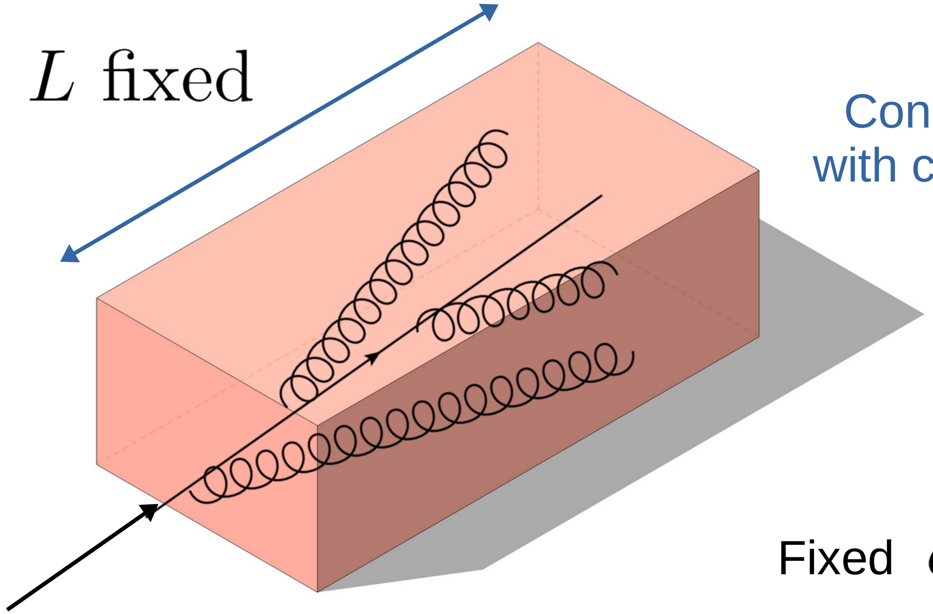




The Monte Carlo

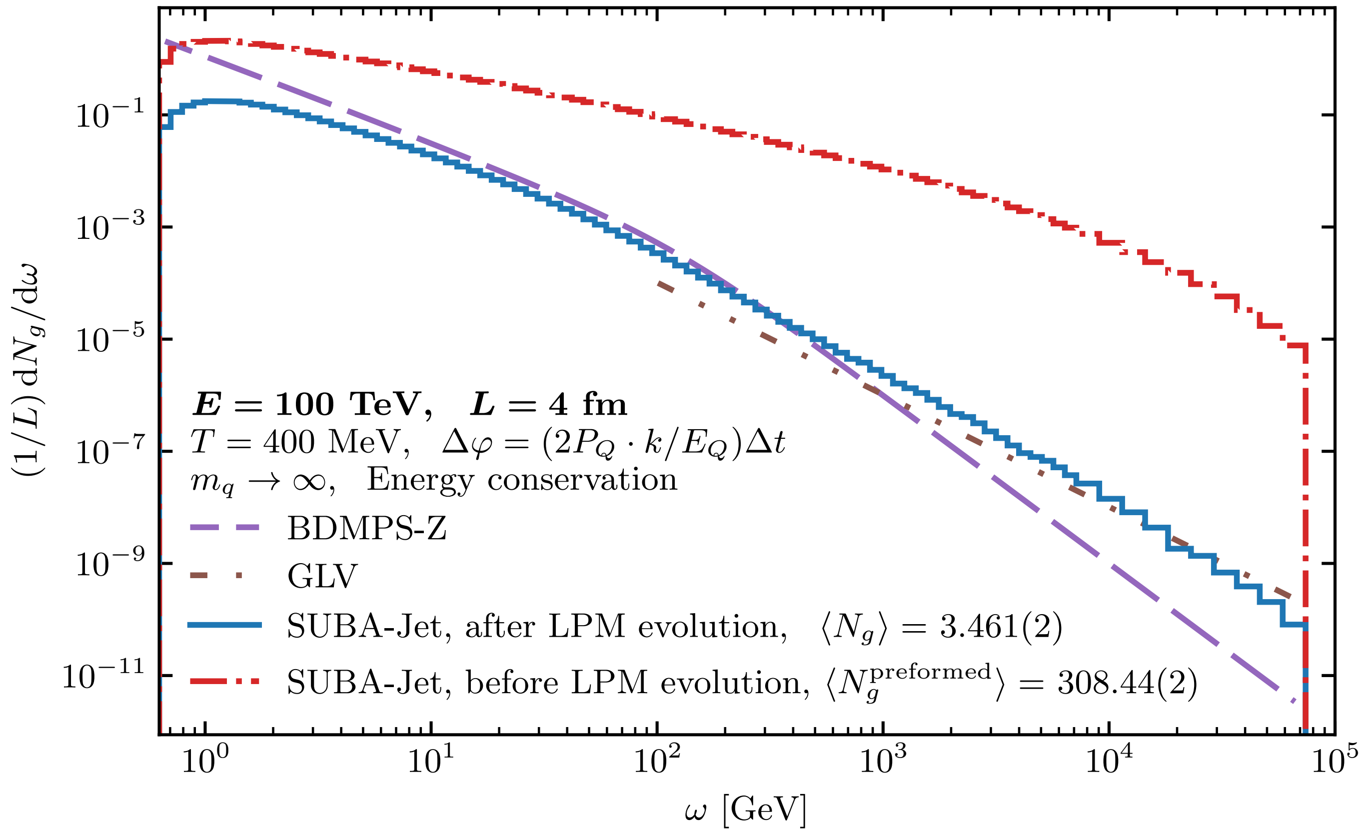


First Results



We want to reproduce theoretical expectations

Reproduction of the LPM and GLV regimes

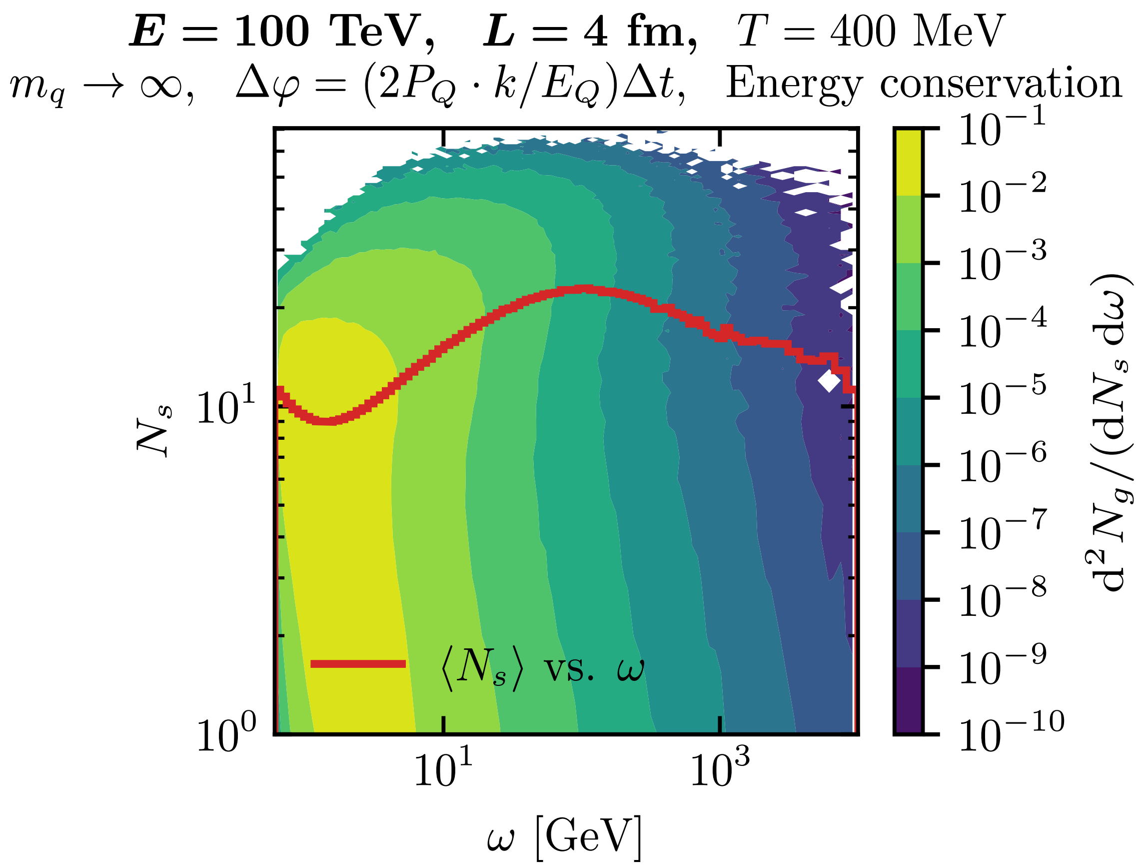


Gluon energy ω spectrum

$$\frac{\mathrm{d}N^{\mathrm{LPM}}}{\mathrm{d}\omega} \sim \frac{1}{\omega^{3/2}}$$

$$\frac{\mathrm{d}N^{\mathrm{GLV}}}{\mathrm{d}\omega} \sim \frac{1}{\omega^{2}}$$

Reproduction of the LPM and GLV regimes



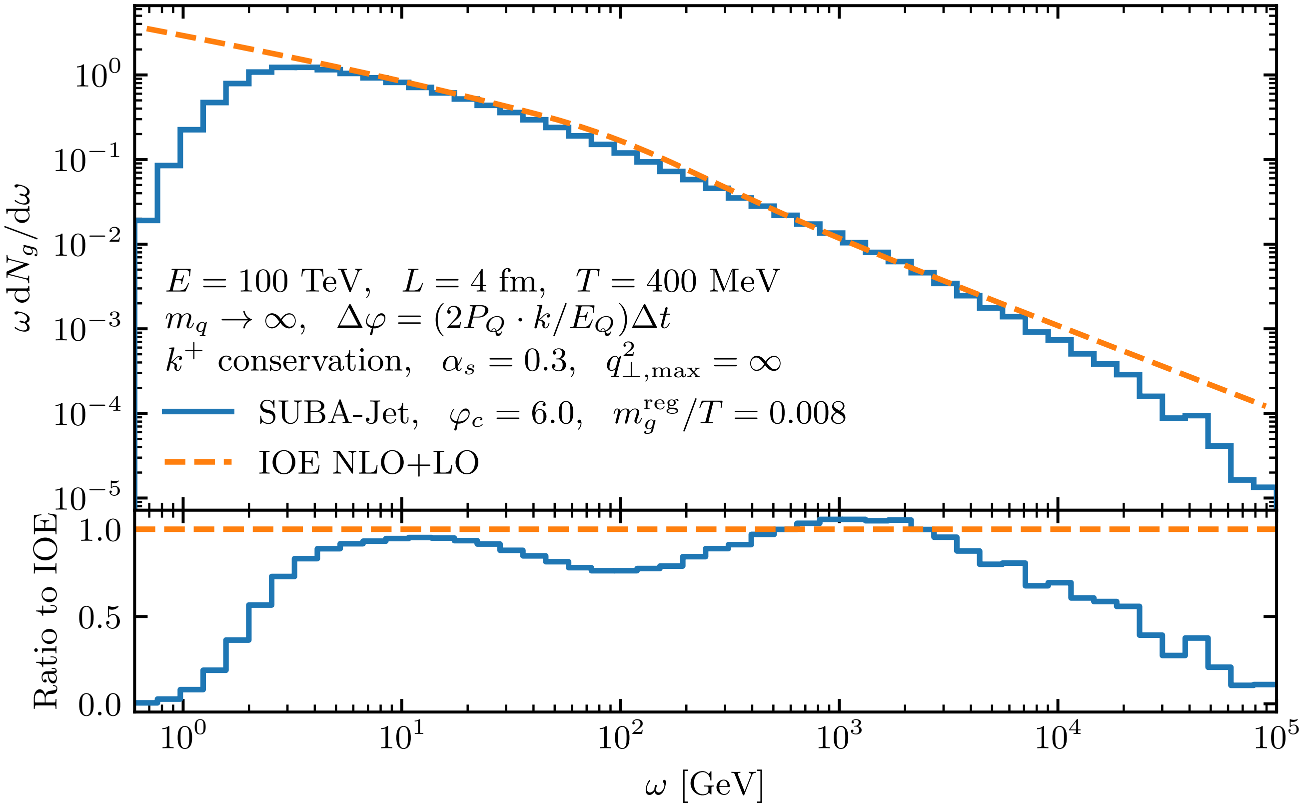
Double differential plot in N_s and ω

Red line: ‹N_s› vs. ω

$$N_S \sim t_f \sim \sqrt{\omega}$$

Convolution of different distributions

Reproduction of the IOE



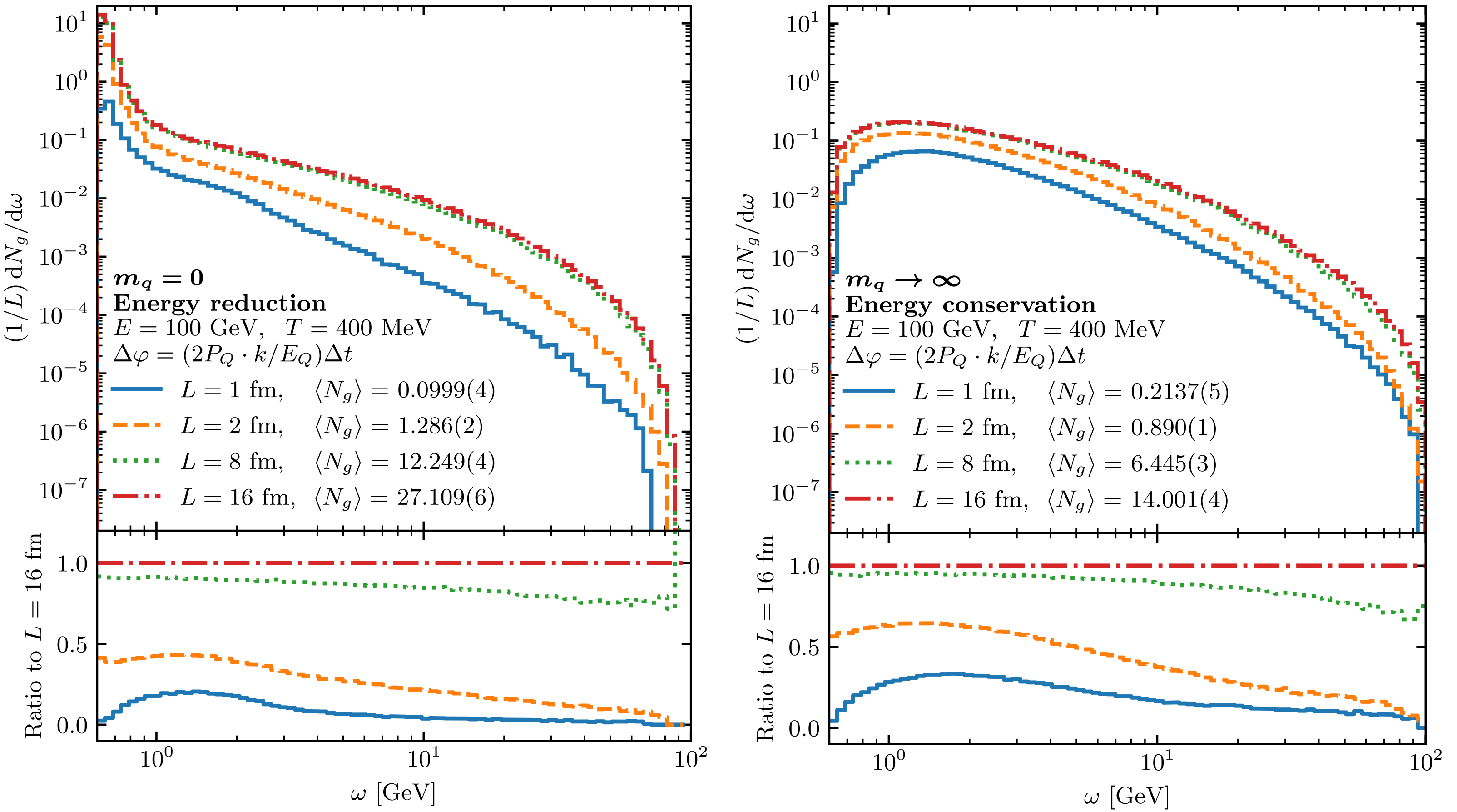
Improved Opacity Expansion

arXiv:1910.02032 [hep-ph]

$$\frac{\mathrm{d}N^{\mathrm{IOE}}}{\mathrm{d}\omega} = \frac{\mathrm{d}N^{\mathrm{IOE}}_{\mathrm{LO}}}{\mathrm{d}\omega} + \frac{\mathrm{d}N^{\mathrm{IOE}}_{\mathrm{NLO}}}{\mathrm{d}\omega}$$

$$\omega \frac{\mathrm{d}N^{\mathrm{IOE}}_{\mathrm{NLO}}}{\mathrm{d}\omega} \simeq \frac{1}{2}\frac{\alpha_s C_R}{\pi}\hat{q}_0\, \mathrm{Re}\int_0^L \mathrm{d}s\, \frac{1}{k^2(s)}\left[\ln\frac{k^2(s)}{Q^2} + \gamma_E\right]$$

The Role of the Path Length



$$\frac{\mathrm{d}N}{\mathrm{d}\omega} \overset{L\to\infty}{\approx} L$$

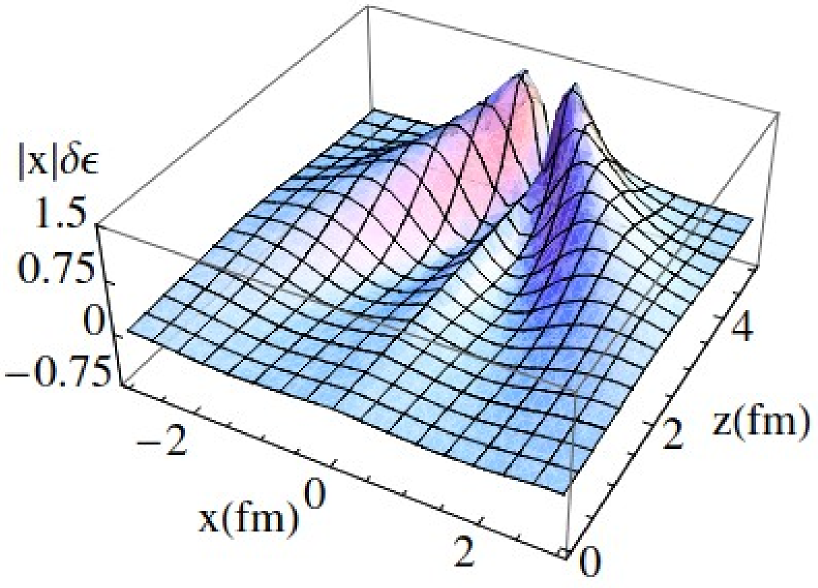
Looking Forward: Towards More Realism

Next step:

- **Implementation in two different MC frameworks:**
  - SUBA-Jet + EPOS4
  - SUBA-Jet + Pythia 8
- Realistic QGP evolution
- Pre-equilibirum stage
- Hadronisation
- Hadronic phase

Looking Forward: Effect on the Medium

The jet also affects the medium



G.-Y. Qin, A. Majumder, H. Song, U. Heinz
0903.2255 [nucl-th]

‘Wake wave’
in the medium
due to the jet

Summary

- We have presented a new model for jet energy loss in heavy ion collisions
- Implementation in a Monte Carlo framework

- **First step done**:
  - Reproduction of the BDMPS-Z and GLV limits

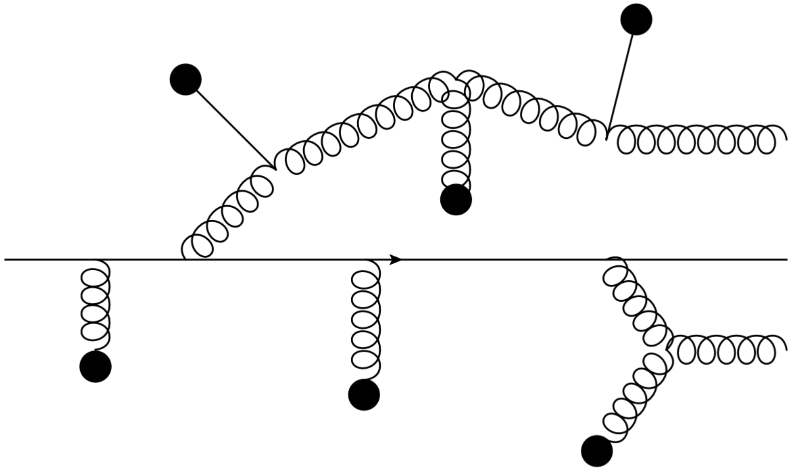
- **Next step:** Implementation within the new EPOS4
  - **EPOS4+JETS –** Initial state, hydro, and hadronisation from EPOS4

Thank you for your attention!

Backup Slides

The Role of the Phase Accumulation

Choice of phase accumulation of the preformed (trial) gluons:



- **More general formula:**

$$\Delta\varphi = \frac{2P_Q \cdot k}{E_Q}\Delta t$$

- **What is used in JEWEL:**

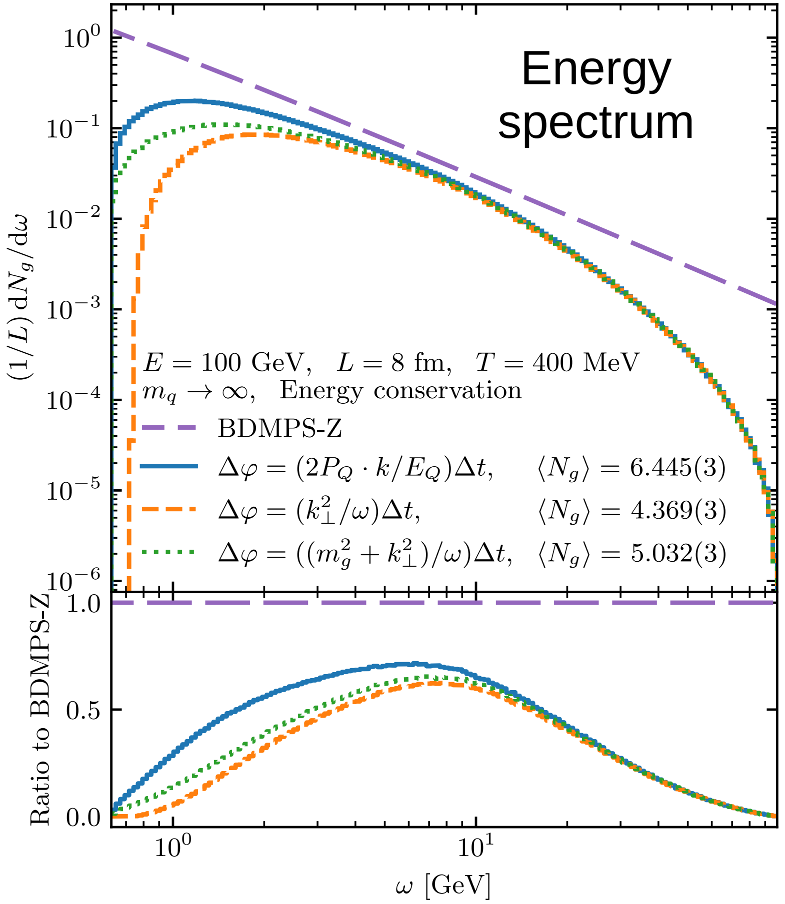
$$\Delta\varphi = \frac{k_T^2}{\omega}\Delta t$$

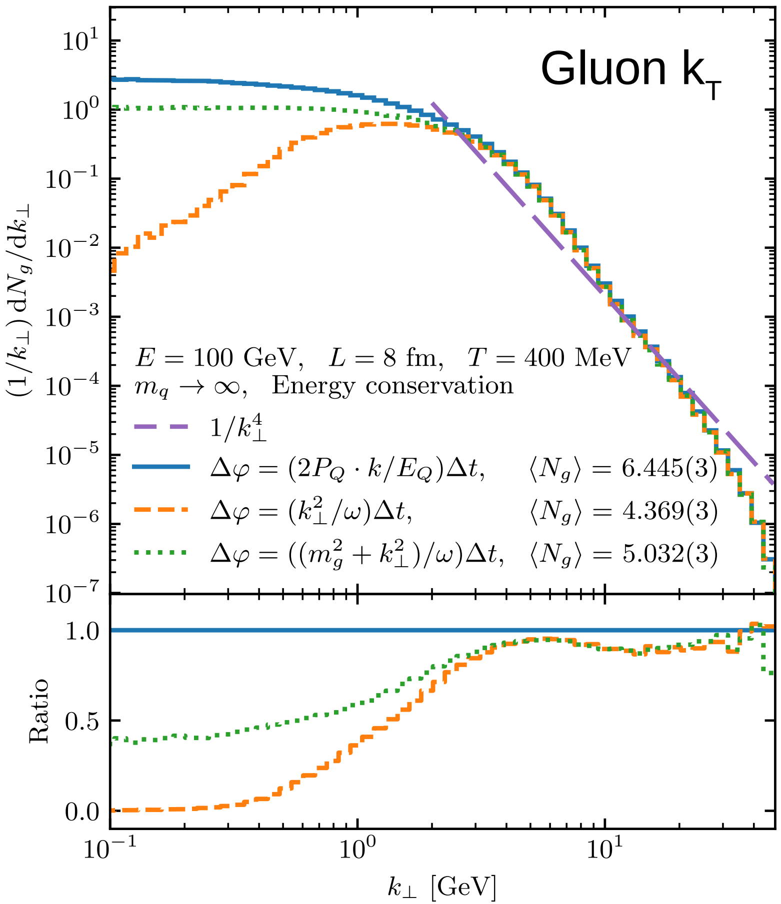
- **Including thermal gluon mass:**

$$\Delta\varphi = \frac{m_g^2 + k_T^2}{\omega}\Delta t$$

The Role of the Phase Accumulation

Effects at low energy & low k_T

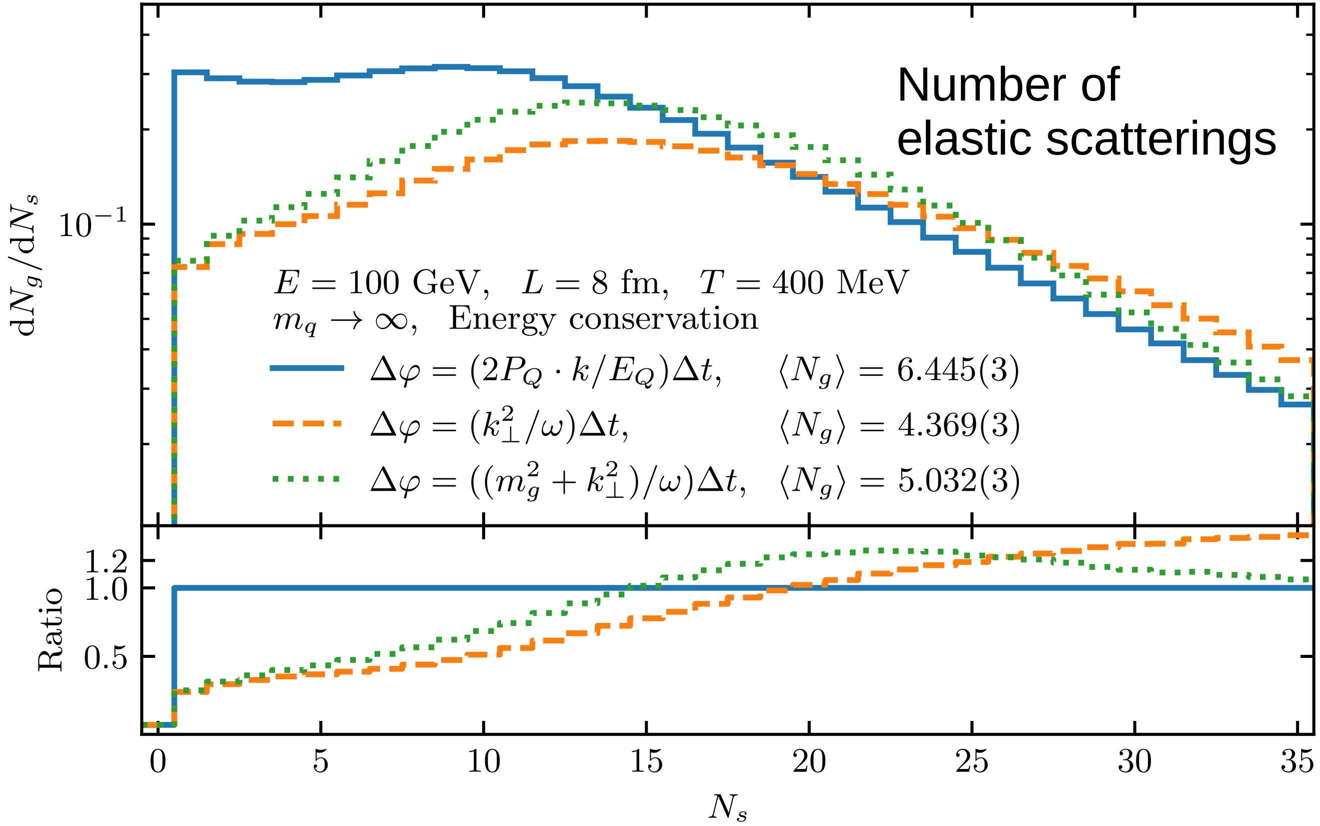




Number of radiated gluons per jet ~ 5 -- 6

The Role of the Phase Accumulation

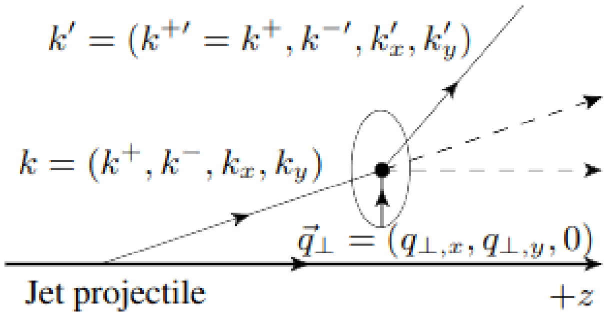
Effects at low N_S



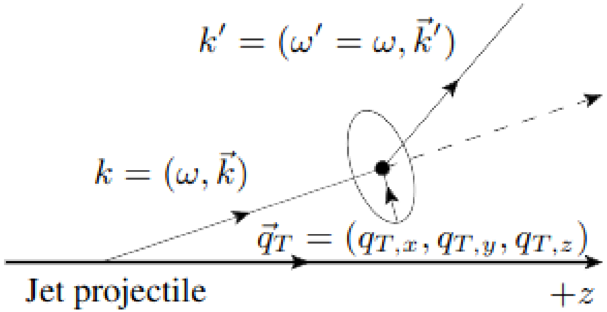
When neglecting the gluon mass in the phase accumulation, a larger path length is required to have a comparable overall number off radiations

The Role of the Elastic Scatterings

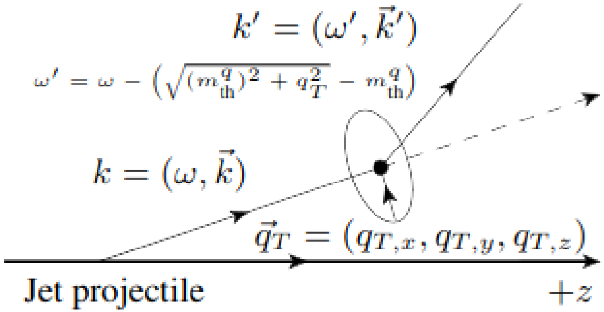
Choice of prescription in elastic scatterings:



k+ conservation

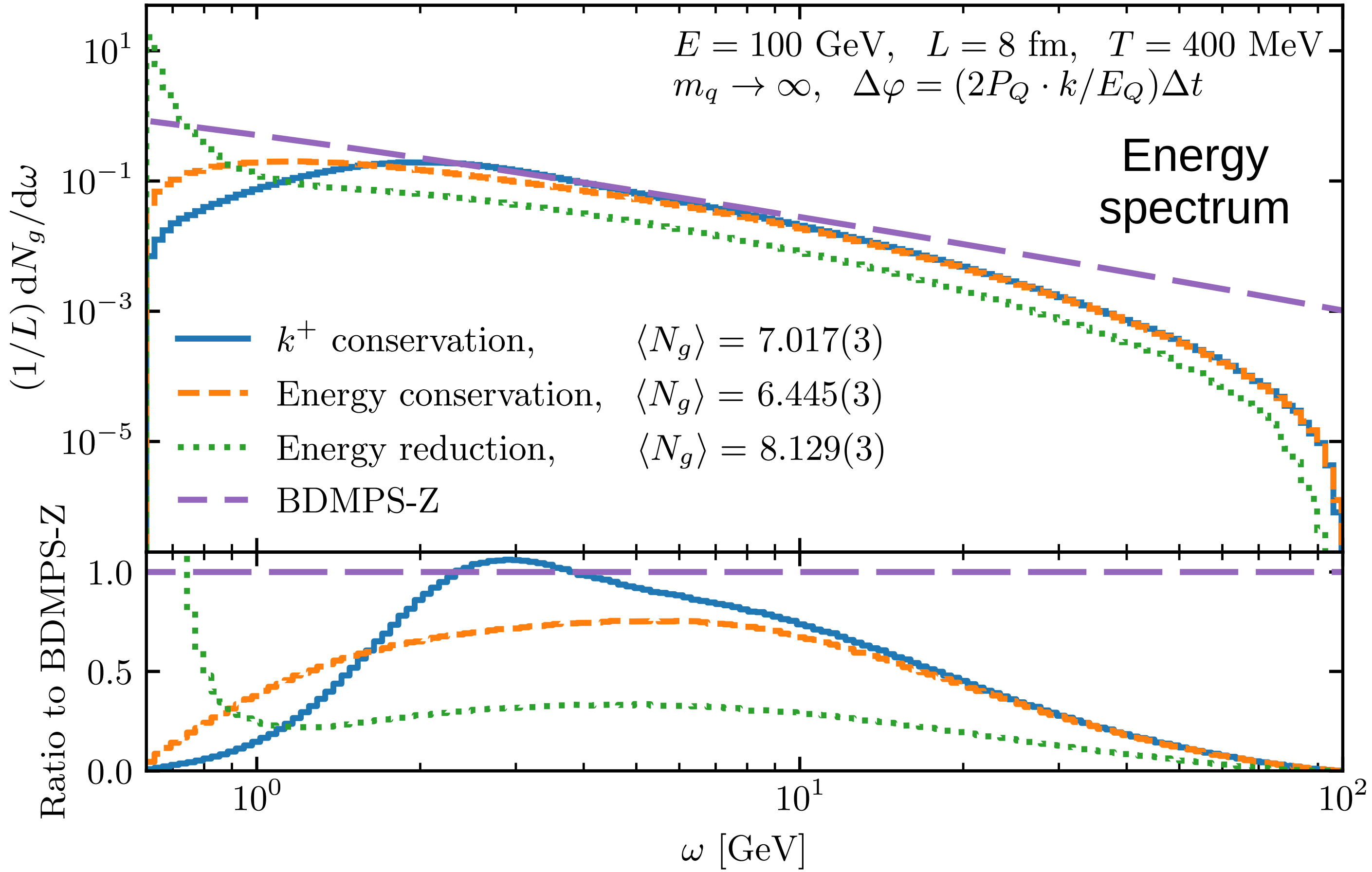


Energy conservation



Energy reduction

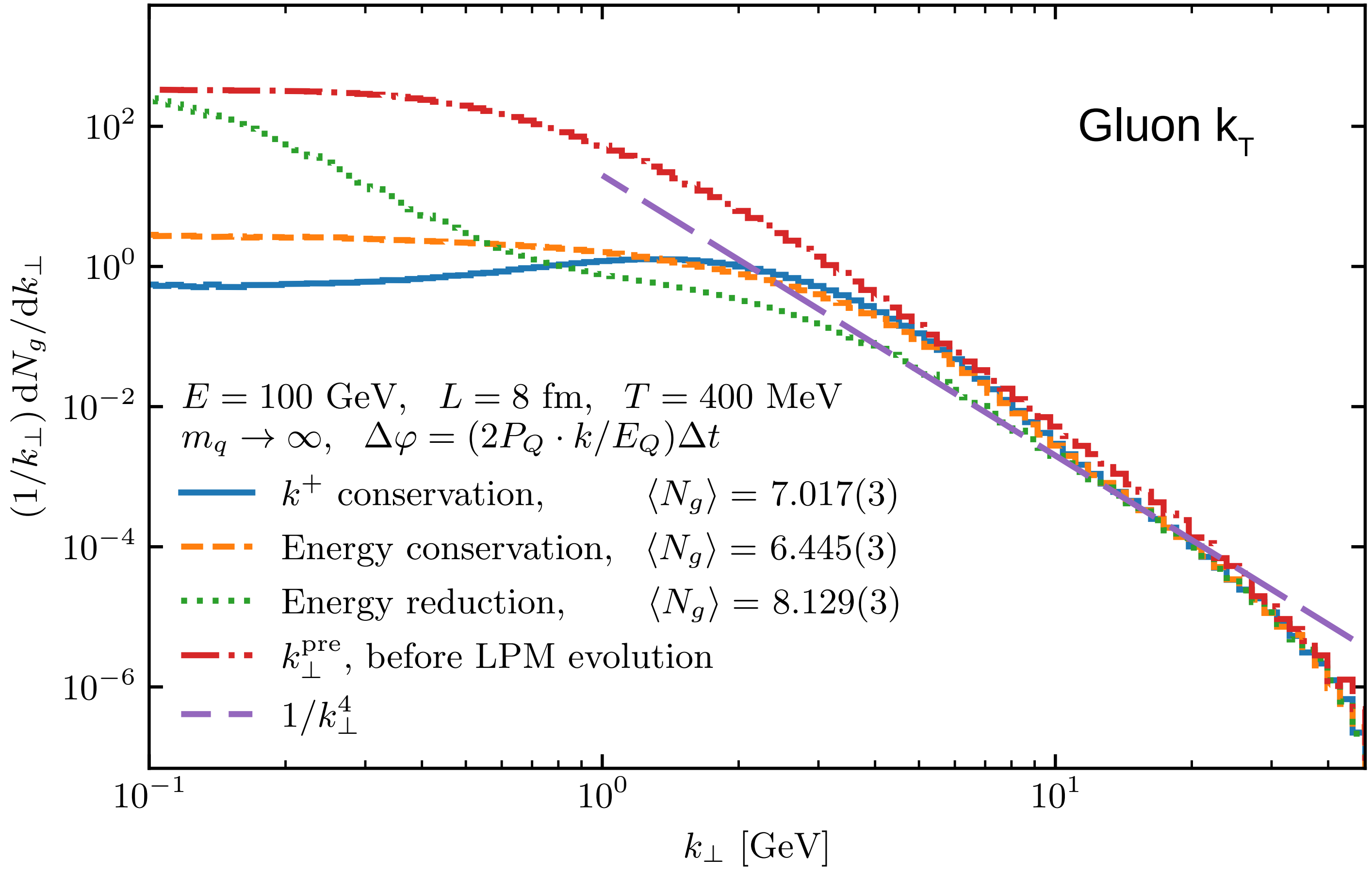
The Role of the Elastic Scatterings



Same BDMS behaviour at intermediate energies

Difference at small energies

The Role of the Elastic Scatterings



Large difference at small k_T