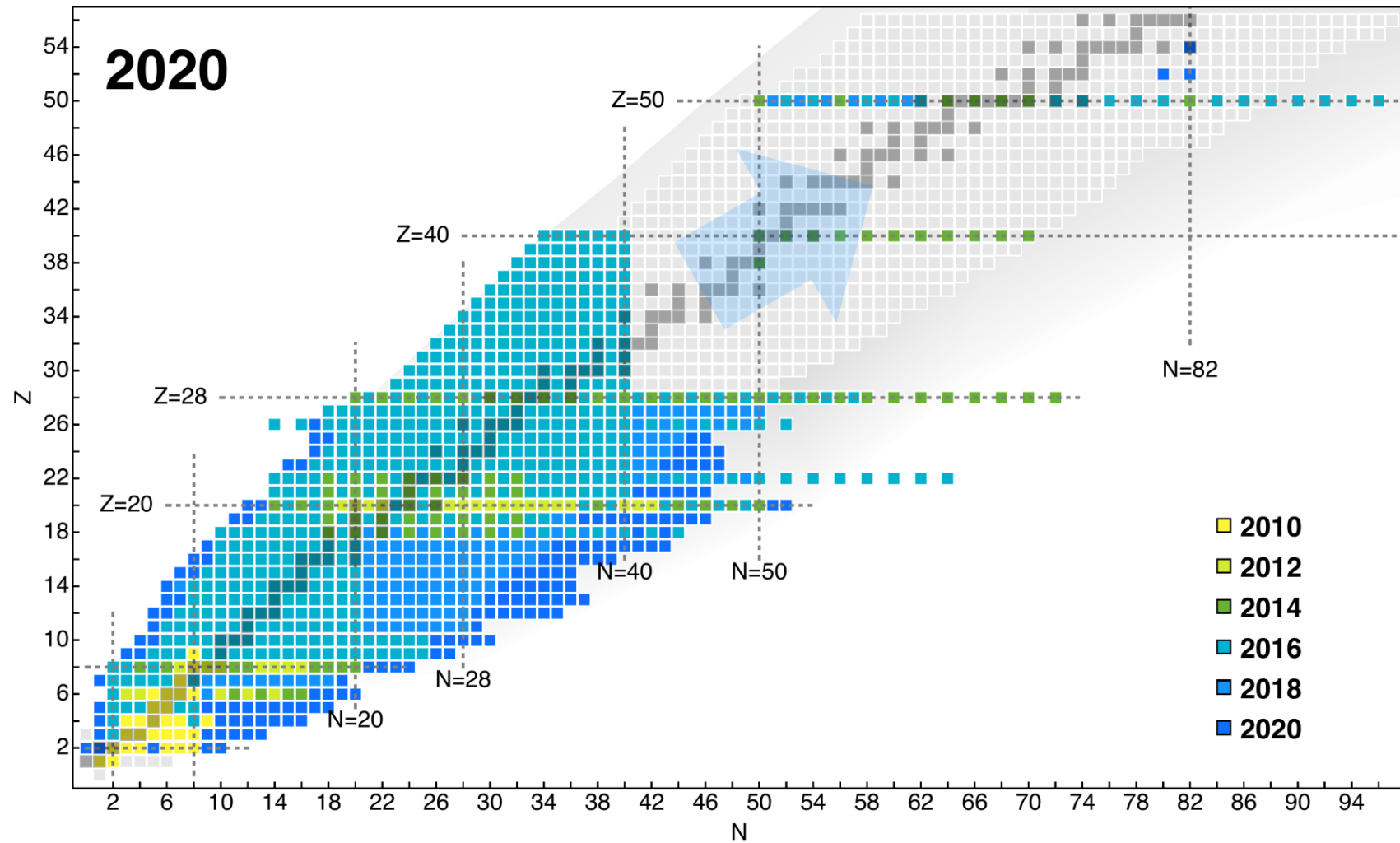


Ab Initio Nuclear Structure with Machine Learning

Javi Rozalén Sarmiento

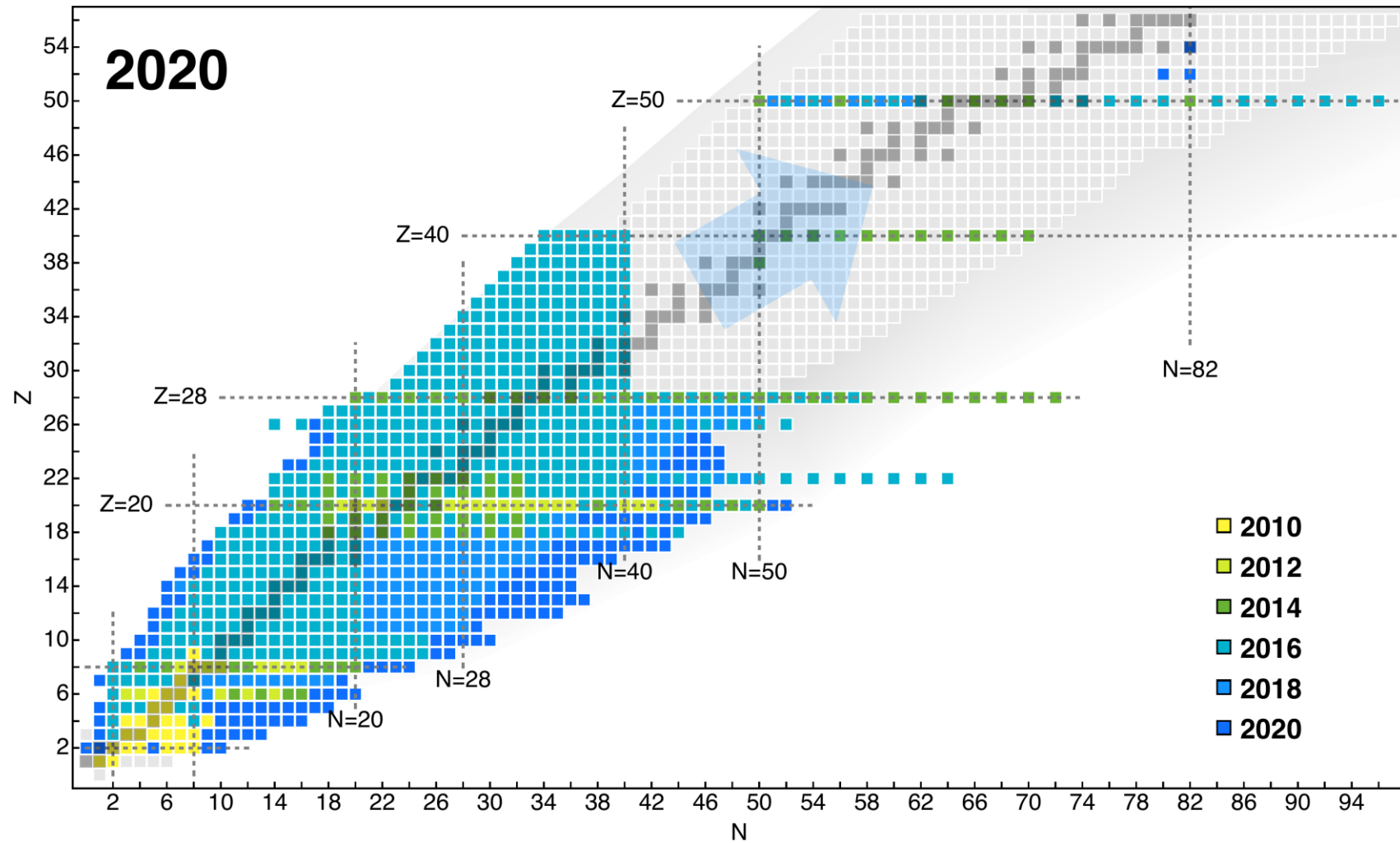
Ab Initio Nuclear Structure



[1] H. Hergert, *A Guided Tour of ab initio Nuclear Many-Body Theory*, 2020

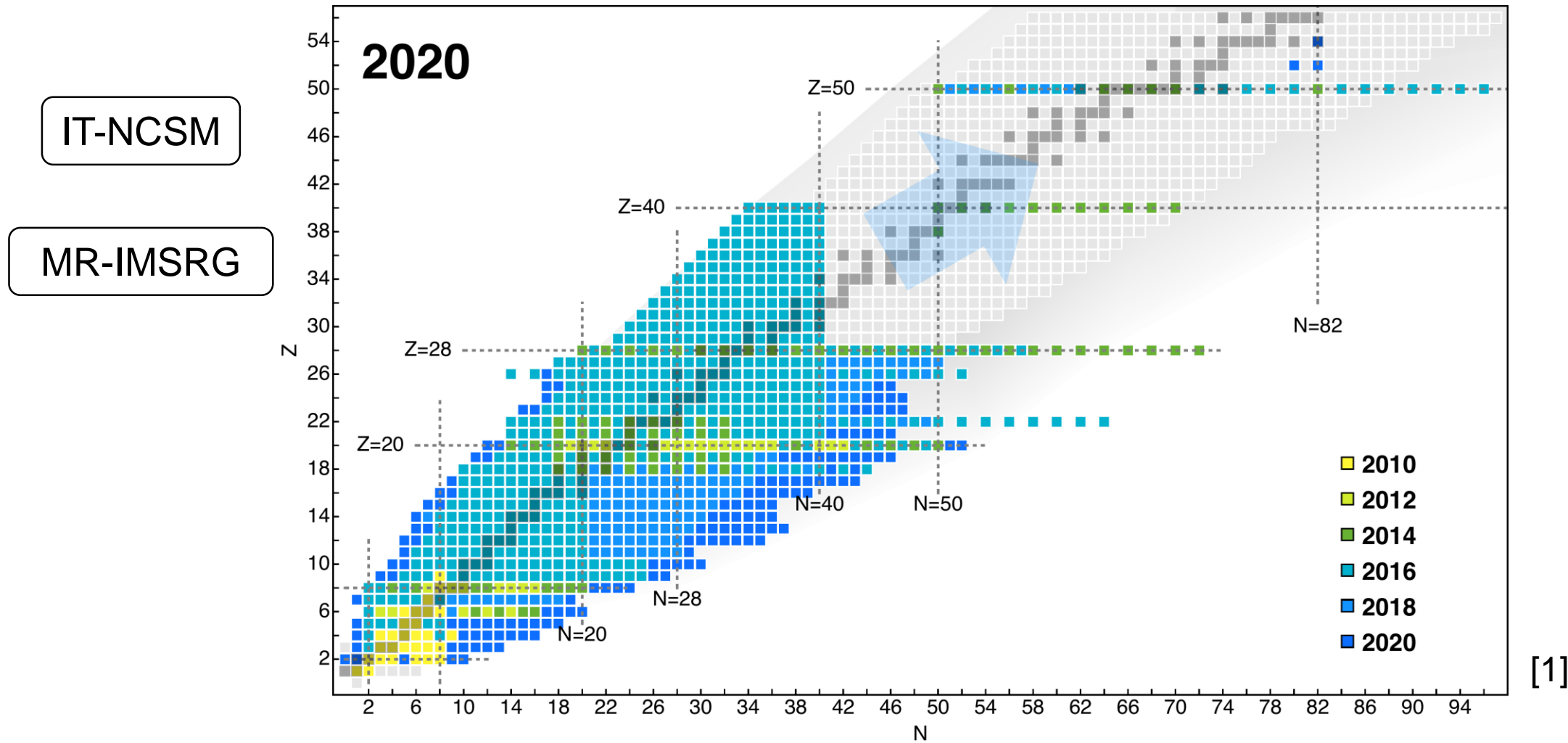
Ab Initio Nuclear Structure

IT-NCSM

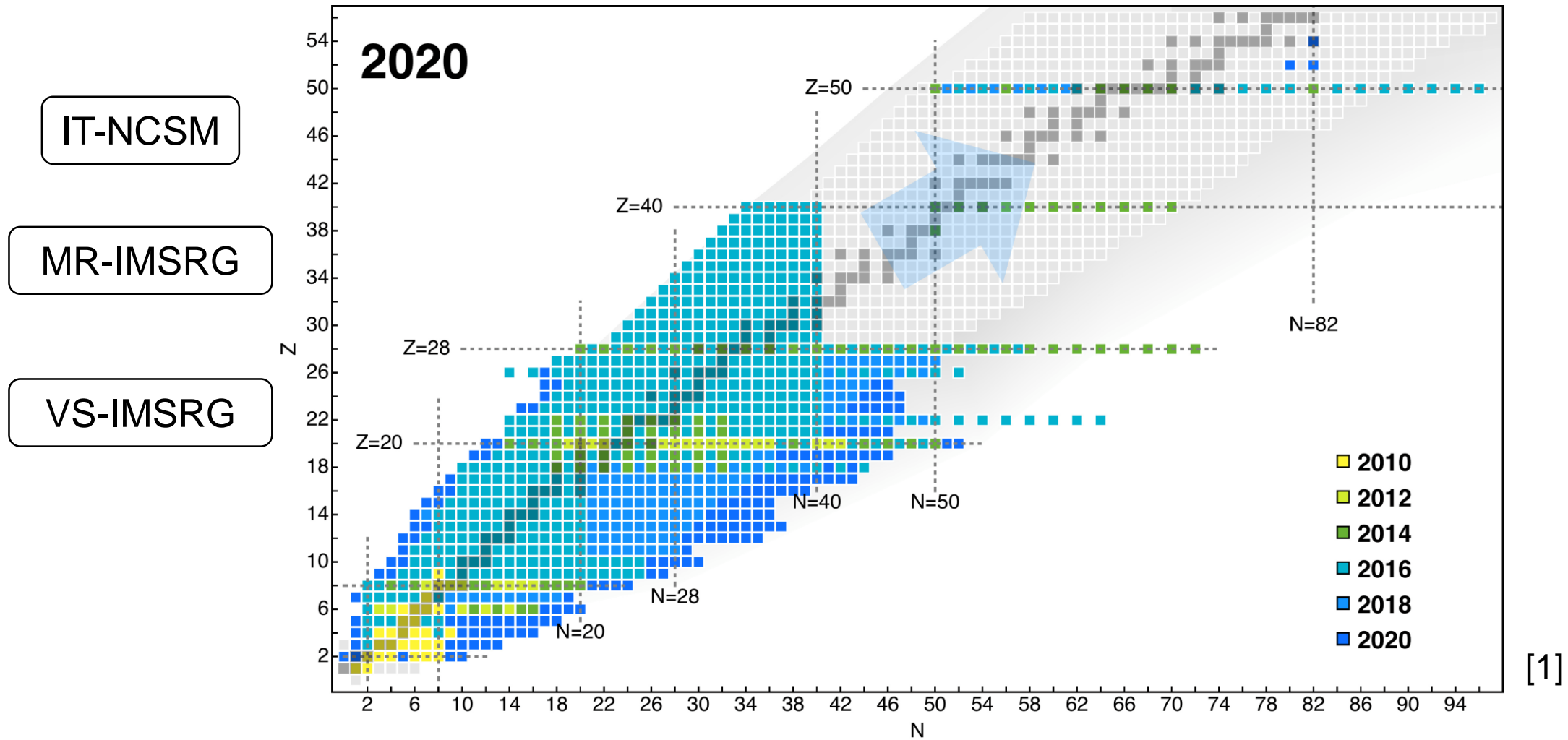


[1]

Ab Initio Nuclear Structure

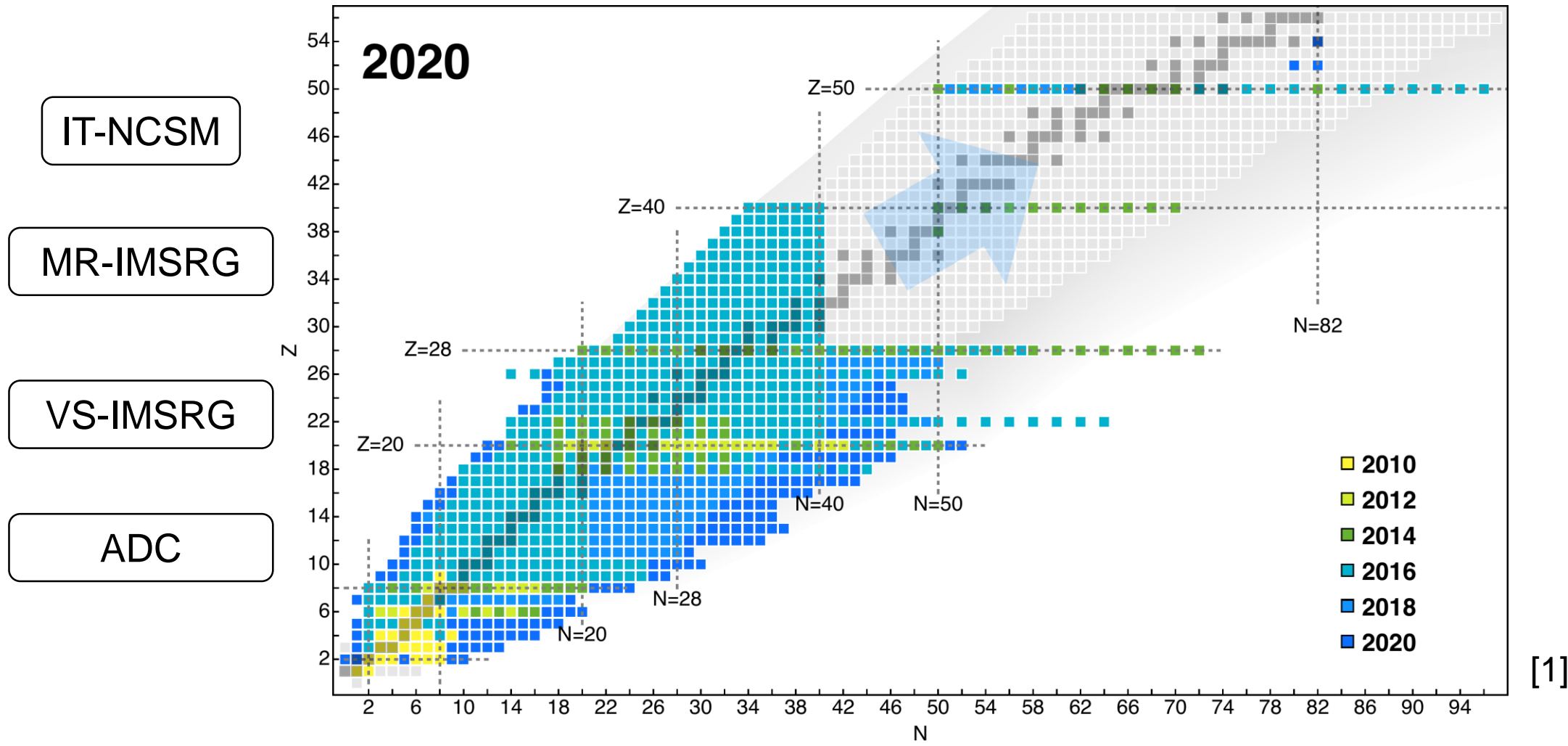


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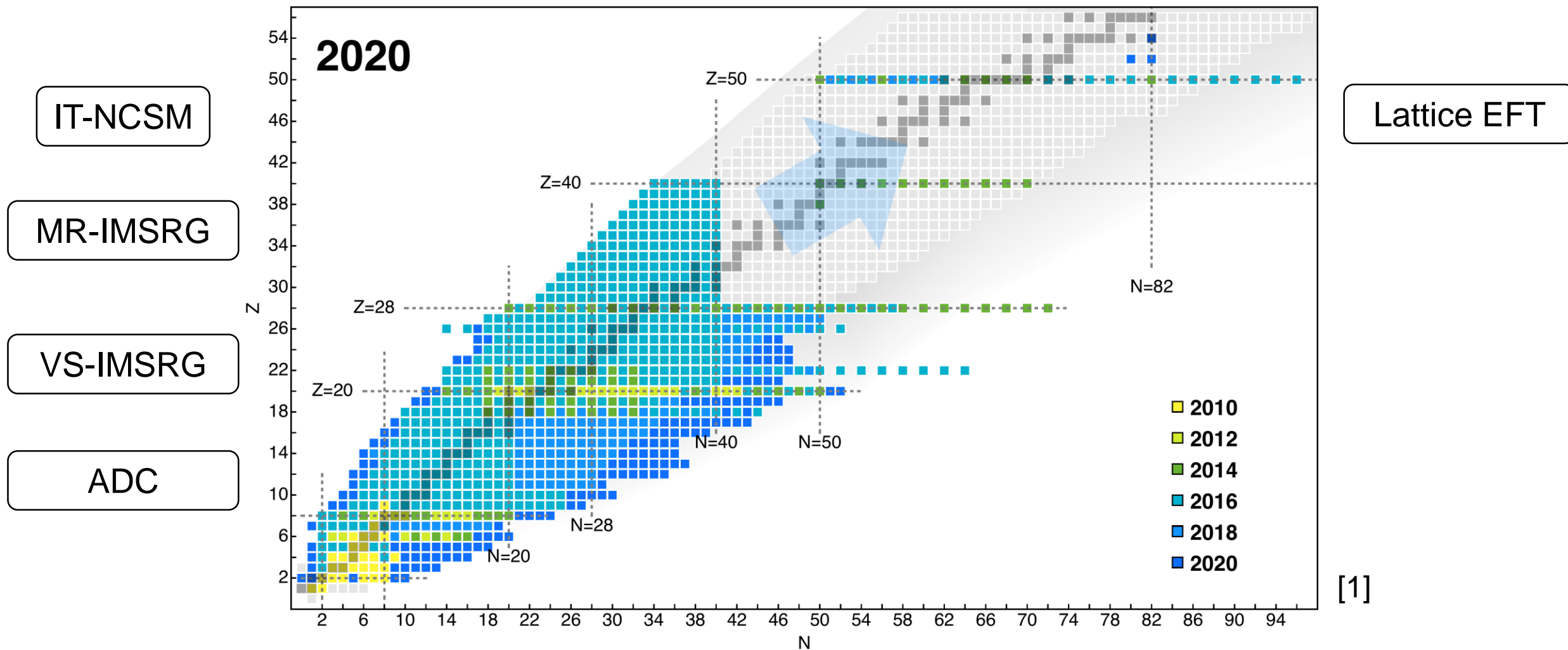
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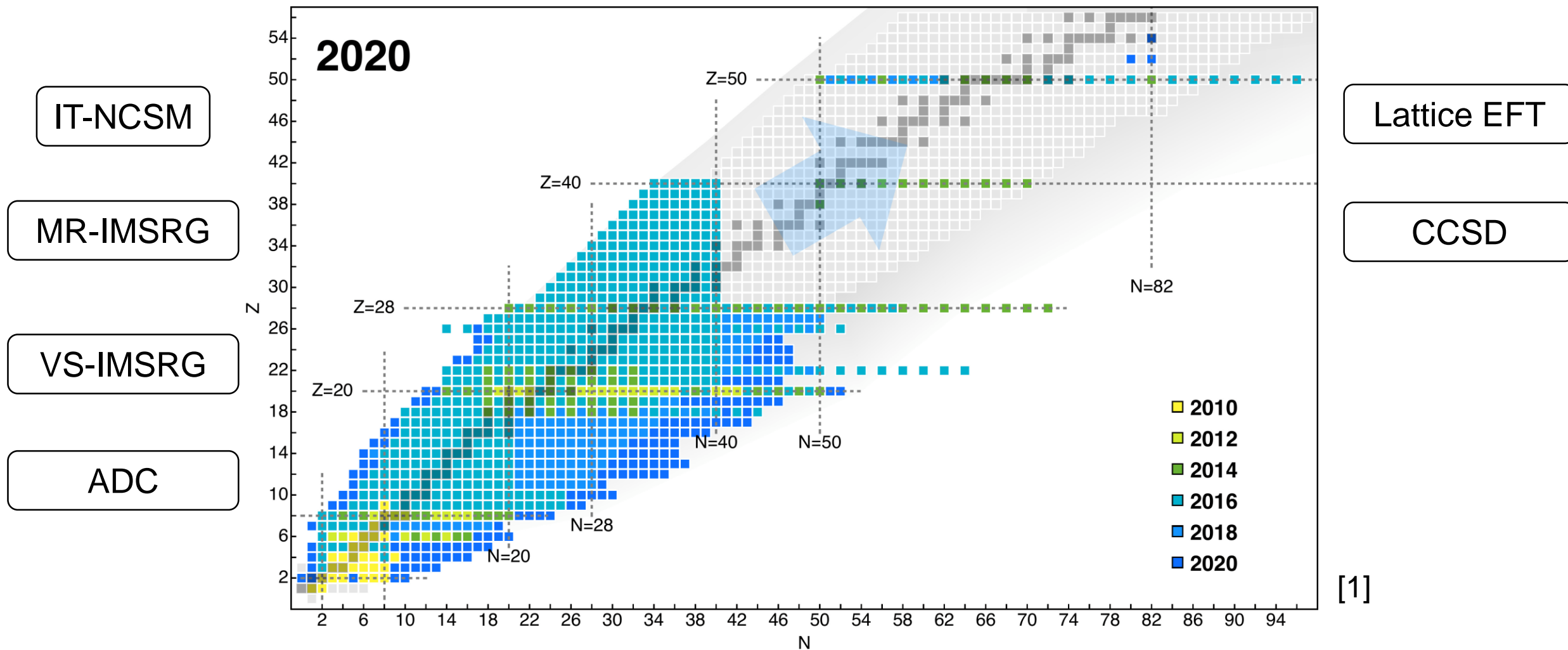
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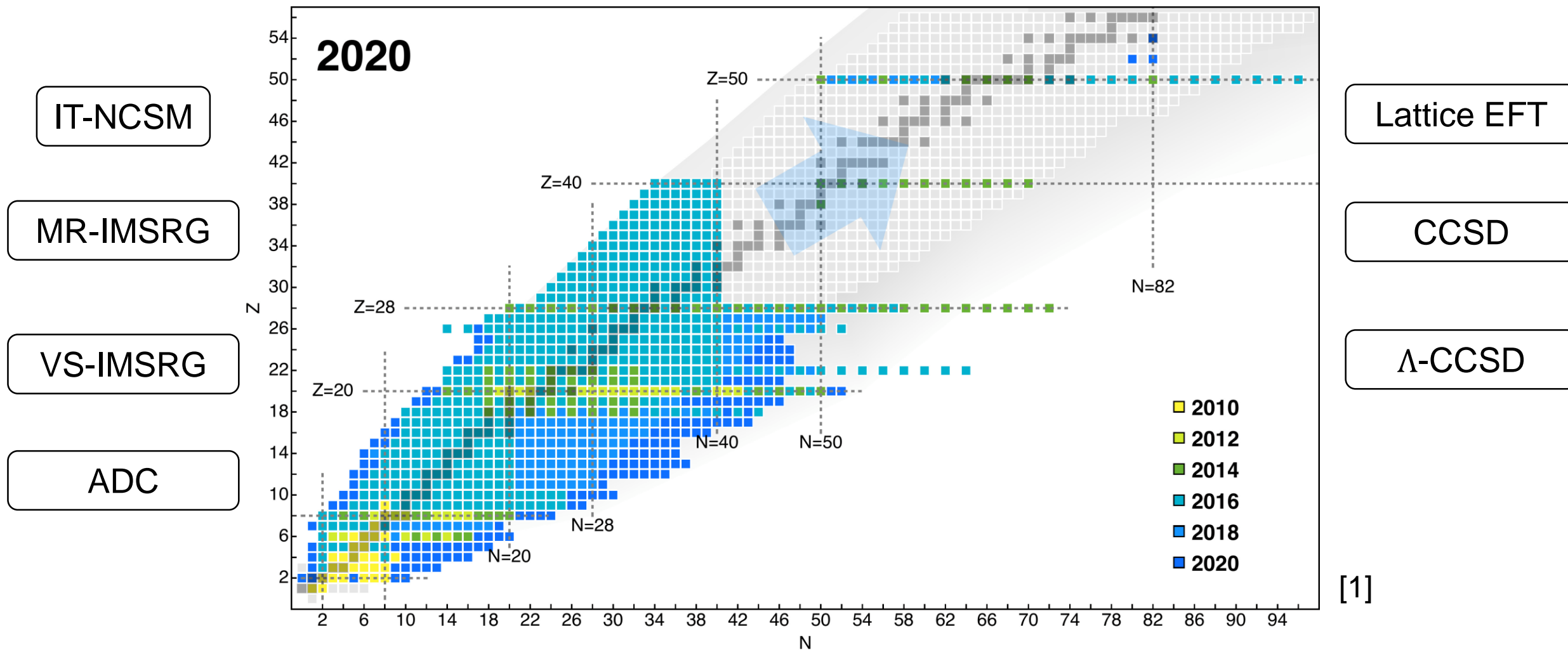
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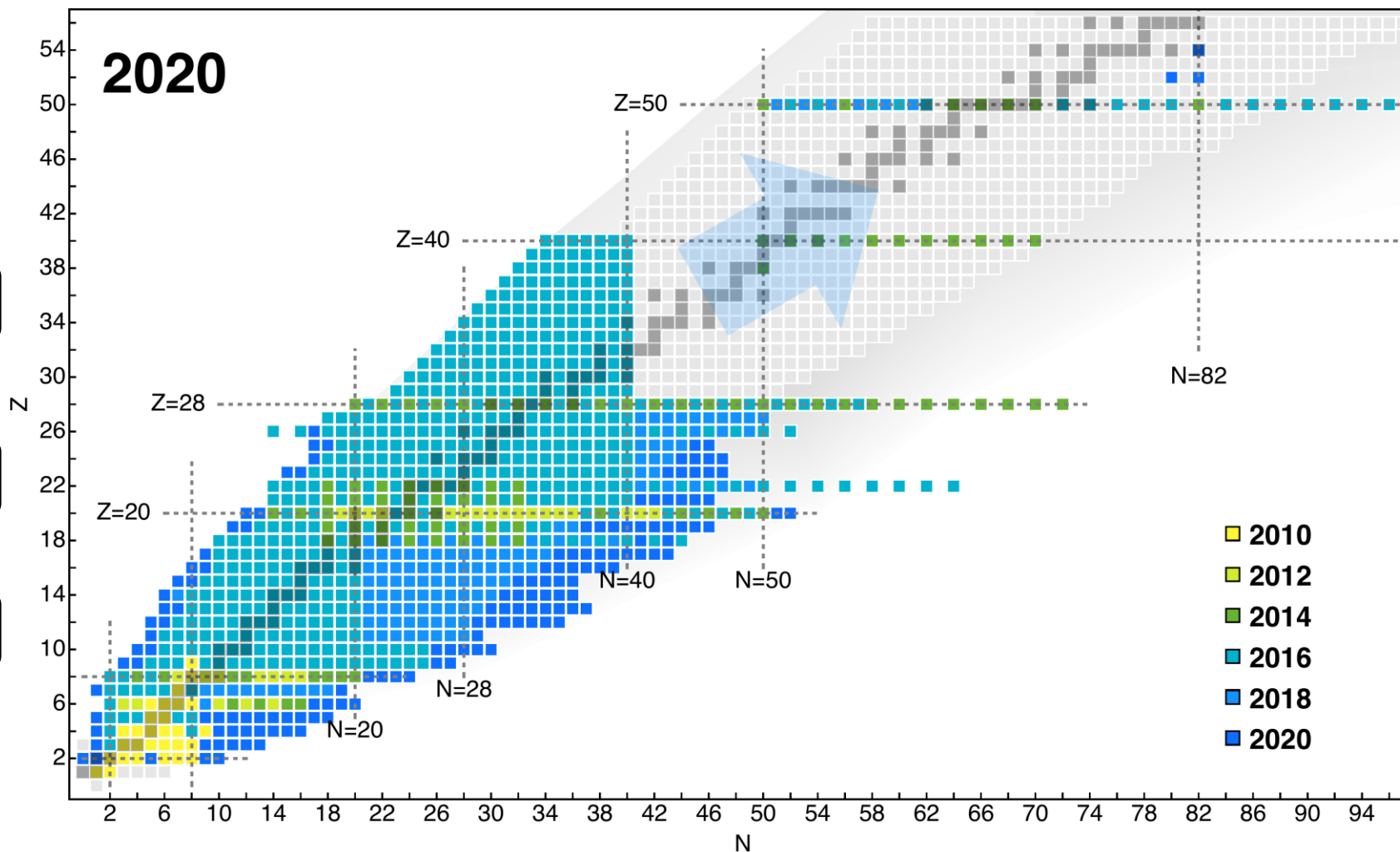
Ab Initio Nuclear Structure

IT-NCSM

MR-IMSRG

VS-IMSRG

ADC



Lattice EFT

CCSD

Λ -CCSD

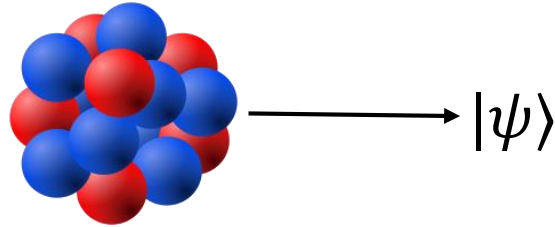
NQS

[1]

[1] H. Hergert, *A Guided Tour of ab initio Nuclear Many-Body Theory*, 2020

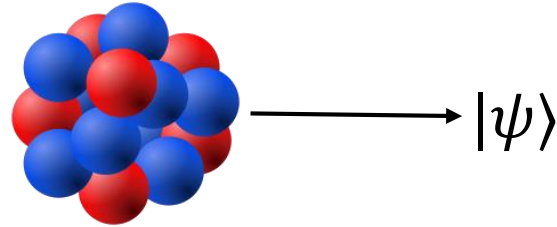
The Quantum Many-Body Problem: NQS

- **Goal:** Solve for the wave function of an atomic nucleus, ψ



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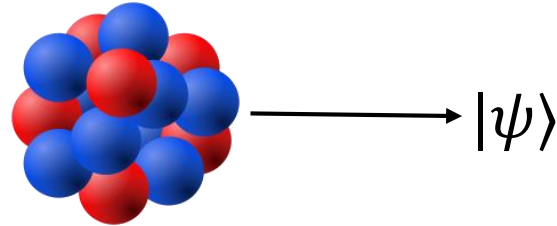
- **How?** Rayleigh-Ritz variational principle

$$\frac{\langle \psi_\theta | \hat{H} | \psi_\theta \rangle}{\langle \psi_\theta | \psi_\theta \rangle} \geq E_{GS}$$

$$|\psi_\theta\rangle = \int |p\rangle \psi_\theta(\mathbf{p}_1, \mathbf{p}_2, \dots, \mathbf{p}_N) d^{3N}p$$

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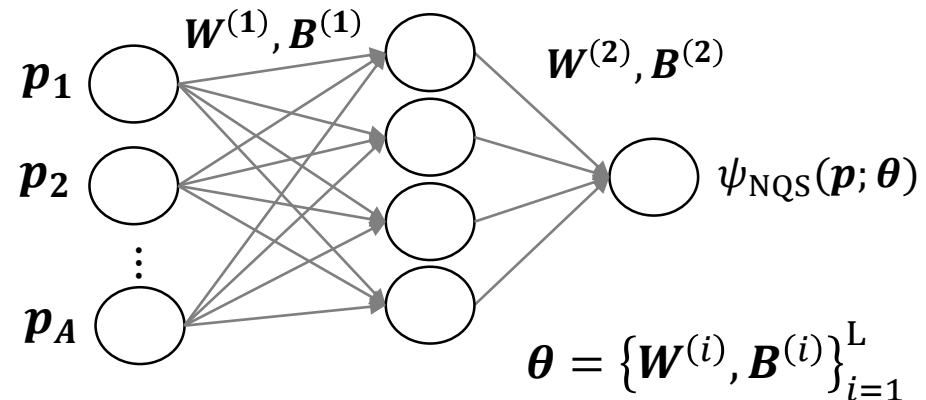


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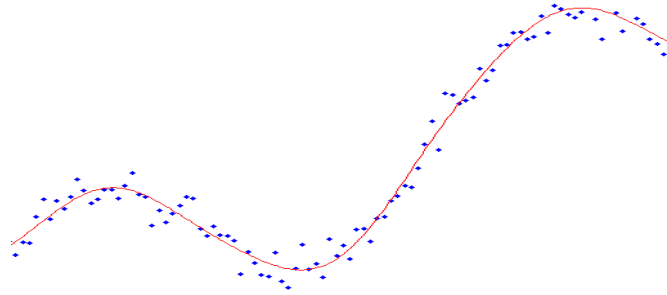
$$|\psi_\theta\rangle = \int |p\rangle \psi_\theta(\mathbf{p}_1, \mathbf{p}_2, \dots, \mathbf{p}_N) d^{3N}p$$

- **NQS Ansatz:** $\psi_{NQS}(\mathbf{p}_1, \mathbf{p}_2, \dots, \mathbf{p}_A; \boldsymbol{\theta})$



Why Neural Networks?

- **NNs have “ ∞ power”**: a neural network can approximate any continuous function [2], [3].

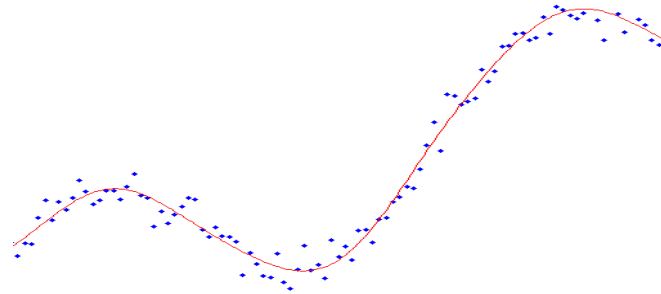


[2] G. Cybenko, *Approximation by superpositions of a sigmoidal function*, 1989

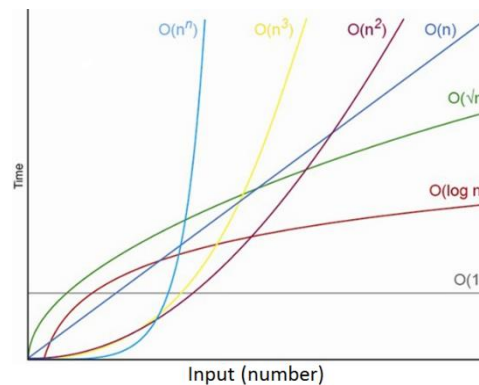
[3] K. Hornik M. Stinchcombe, H. White, *Multilayer feedforward networks are universal approximators*, 1989

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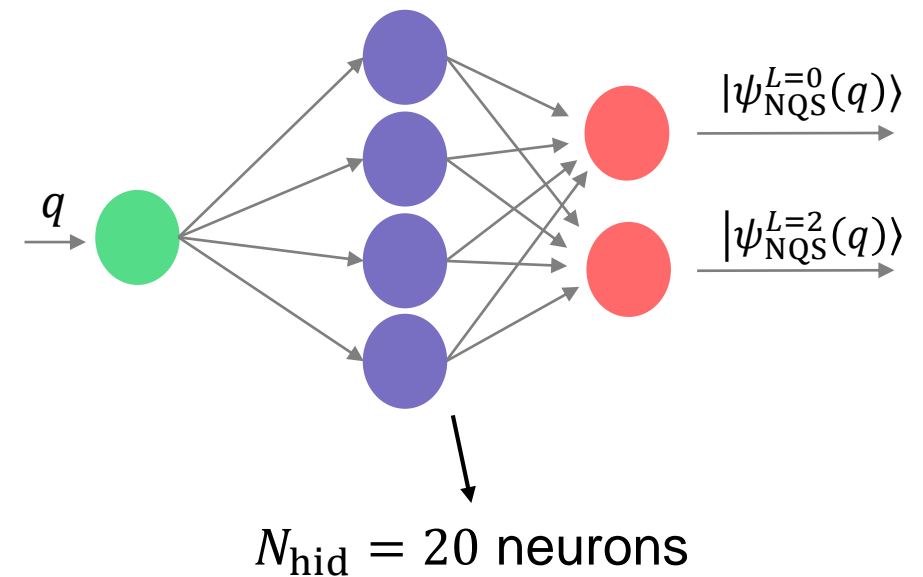
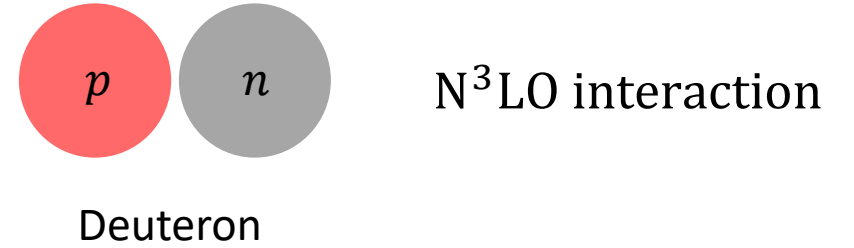
- **Space complexity**: polynomial scaling of memory resources... possibly!



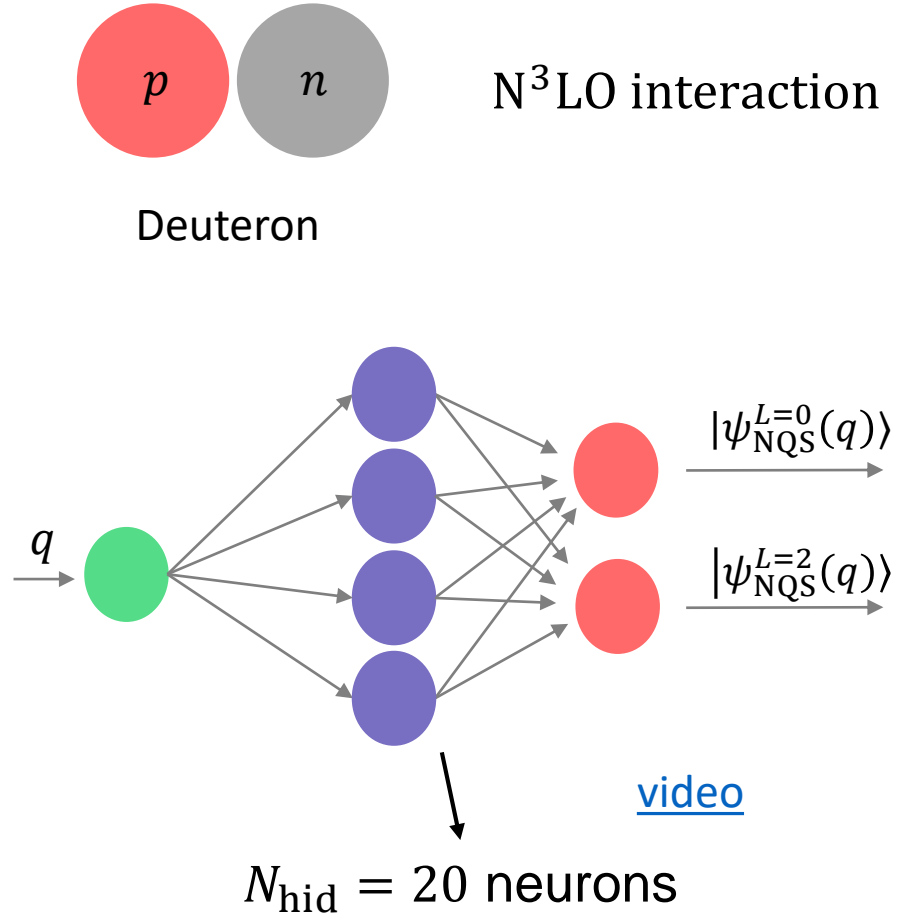
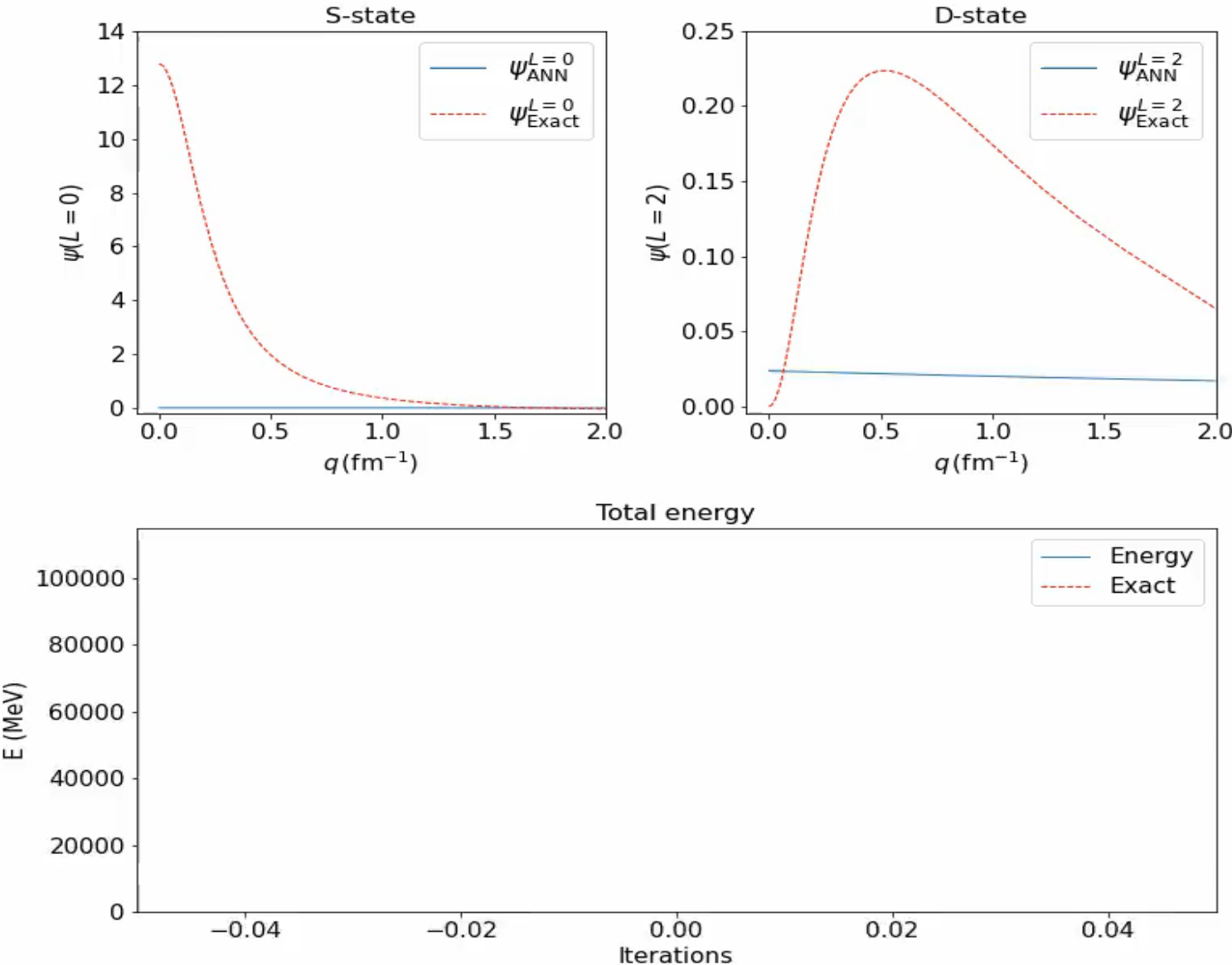
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Live Neural-Network Training

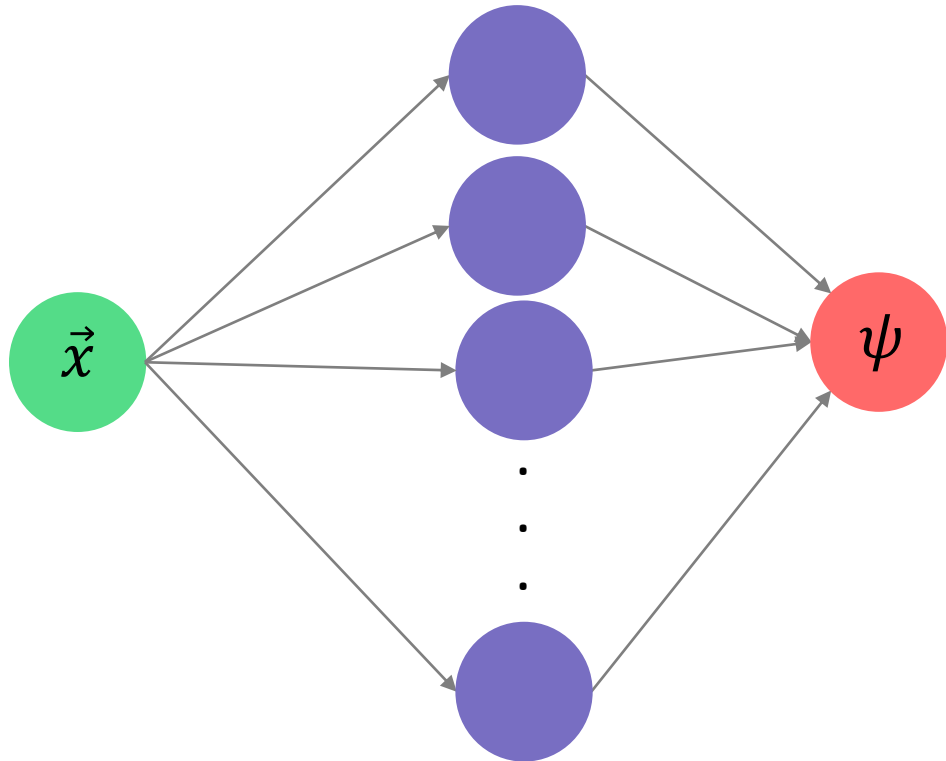


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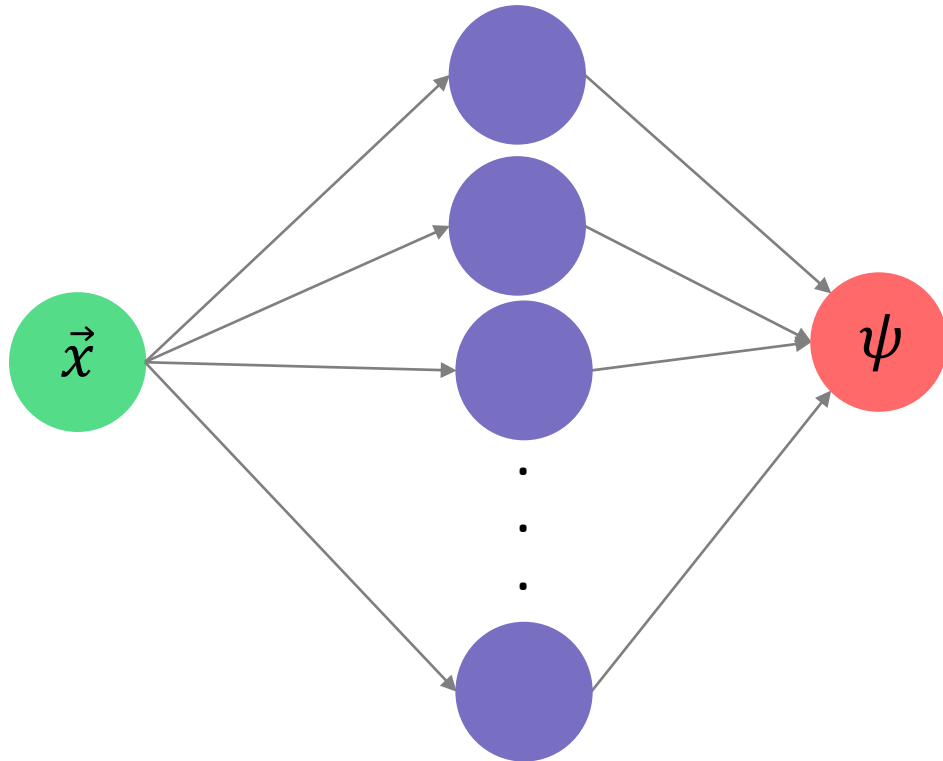


J. Keeble & A. Rios, *Machine Learning the Deuteron*, 2020
 J. Keeble & A. Rios & J. Rozalén Sarmiento, 2024

Physical properties in NNs



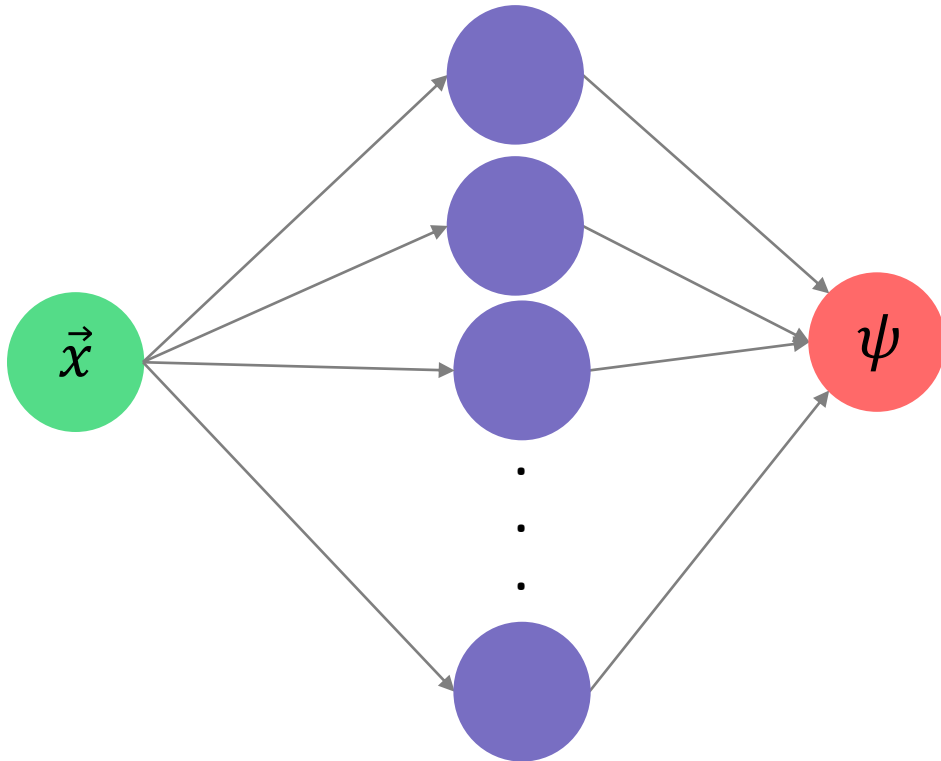
Physical properties in NNs



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$$\psi(x_1, x_2, \dots, x_N) = \pm \psi(x_2, x_1, \dots, x_N)$$

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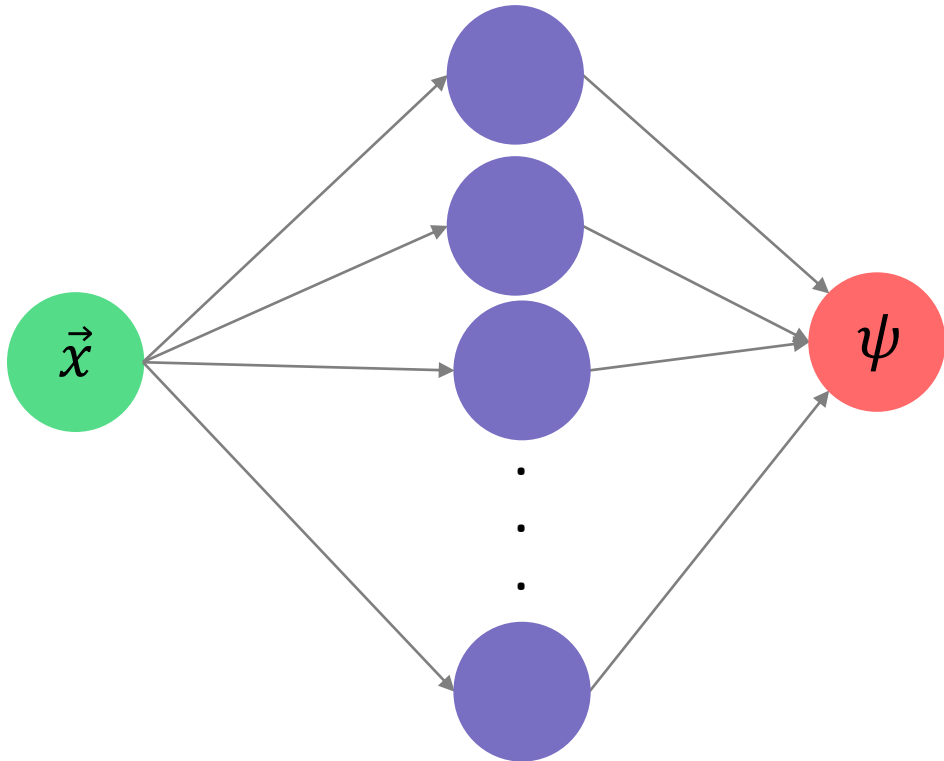
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- **Time-reversal symmetry**

$$\psi(\vec{x}, \sigma) = \psi(T(\vec{x}, \sigma))$$

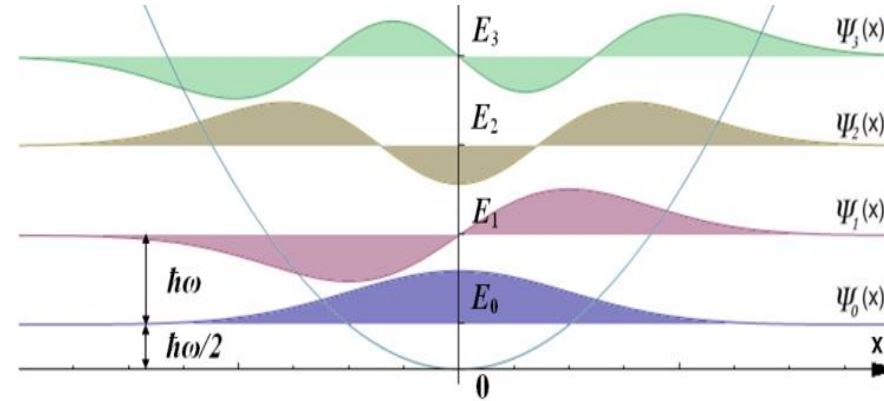
Symmetries in Quantum Mechanics

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- **1D Harmonic Oscillator**

$$\hat{H} = \frac{1}{2} \nabla^2 + \frac{1}{2} m \omega^2 x^2$$

$$\Rightarrow [\hat{H}, \hat{P}] = 0, \hat{P} \text{ parity}$$

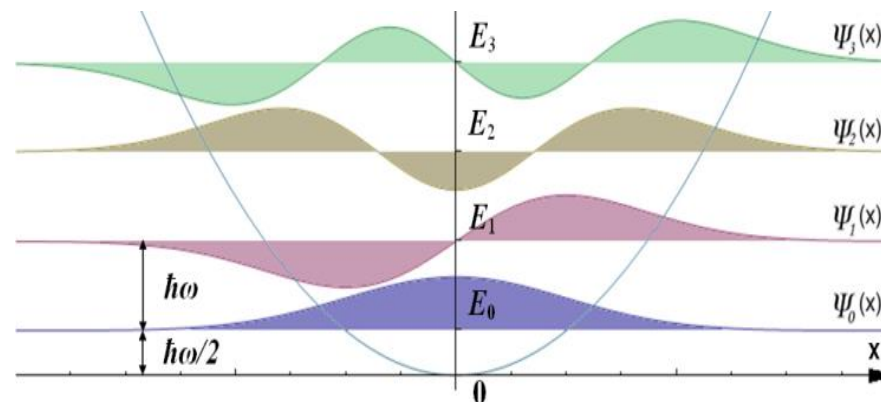


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- **Hydrogen Atom (Coulomb)**

$$\hat{H} = -\frac{\hbar^2}{2\mu} \nabla^2 - \frac{Ze^2}{r}$$

$$\Rightarrow [\hat{H}, \hat{L}^2] = [\hat{H}, \hat{L}_z] = 0$$

$$\psi_{nlm} \left\{ \begin{array}{l} \psi_{100}(r, \theta, \phi) = \left(\frac{Z}{a_0}\right)^{\frac{3}{2}} 2r \exp(-Zr/a_0) \times \sqrt{\frac{1}{4\pi}} \\ \psi_{110}(r, \theta, \phi) = \left(\frac{Z}{a_0}\right)^{\frac{3}{2}} 2r \exp(-Zr/a_0) \times \sqrt{\frac{3}{4\pi}} \cos \theta \end{array} \right.$$

Group theory (in a nutshell)

- **Group**

$$G = \{g \mid g_1, g_2 \in G \Rightarrow g_1 \cdot g_2 = g_3 \in G\}$$



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$$r_\theta \circ r_\varphi = r_{\theta+\varphi} \Rightarrow D(r_\theta)D(r_\varphi) = D(r_{\theta+\varphi})$$

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- **Schur's Lemma**

$$d_1(g) \text{ irrep, } H \text{ operator, } d_1(g)H = Hd_1(g)$$

$$\Rightarrow H = EI, E \text{ number } H = \begin{pmatrix} E_1 I & \dots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \dots & E_n I \end{pmatrix}$$

Convolution is all you need

- **Group Equivariance**

ψ is **equivariant / covariant** under G if:

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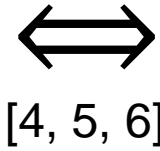
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[4] T S Cohen and M Welling, *Group Equivariant Convolutional Networks* (2016)

[5] T S Cohen and M Welling, *Steerable CNNs* (2016)

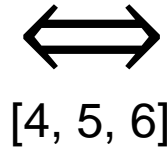
[6] R Kondor and S Trivedi, *On the Generalization of Equivariance and Convolution in Neural Networks to the Action of Compact Groups* (2018)

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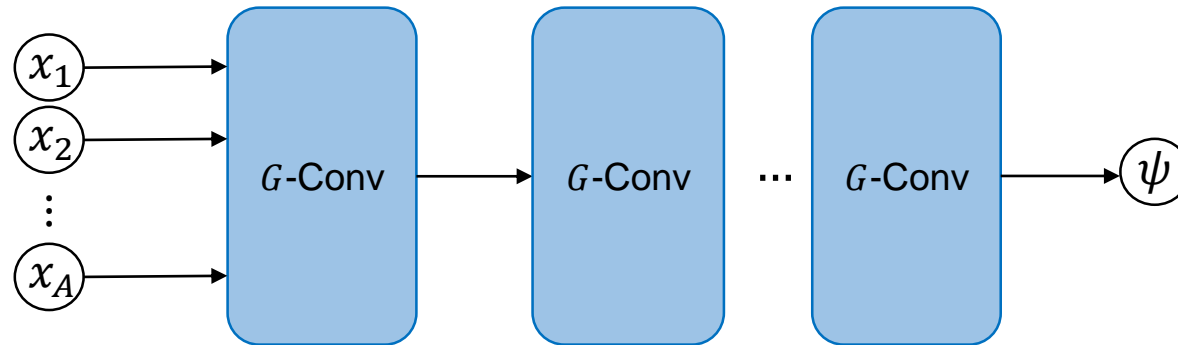
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Re-thinking previous work...

Parity, $\mathbb{Z}/2\mathbb{Z} = \{e, a\}$

- **Intuition**

$$\psi_{\text{NQS}}(x) := \psi_{\text{NN}}(x) \pm \psi_{\text{NN}}(-x)$$

- **Group convolution**

Irreps of $\mathbb{Z}/2\mathbb{Z}$:

$$D^t(e) = (1), \quad D^t(a) = (1) \quad (\text{Trivial rep})$$

$$D^s(e) = (1), \quad D^s(a) = (-1) \quad (\text{Sign rep})$$

$$\begin{cases} \psi_+(x) := \psi(x) + \psi(-x) \\ \psi_-(x) := \psi(x) - \psi(-x) \end{cases}$$

Fermionic particle exchange, S_N

- **Intuition**

$$\psi_{\text{NQS}}(x_1, x_2, x_3) = \begin{vmatrix} \psi_1(x_1) & \psi_2(x_1) & \psi_3(x_1) \\ \psi_1(x_2) & \psi_2(x_2) & \psi_3(x_2) \\ \psi_1(x_3) & \psi_2(x_3) & \psi_3(x_3) \end{vmatrix}$$

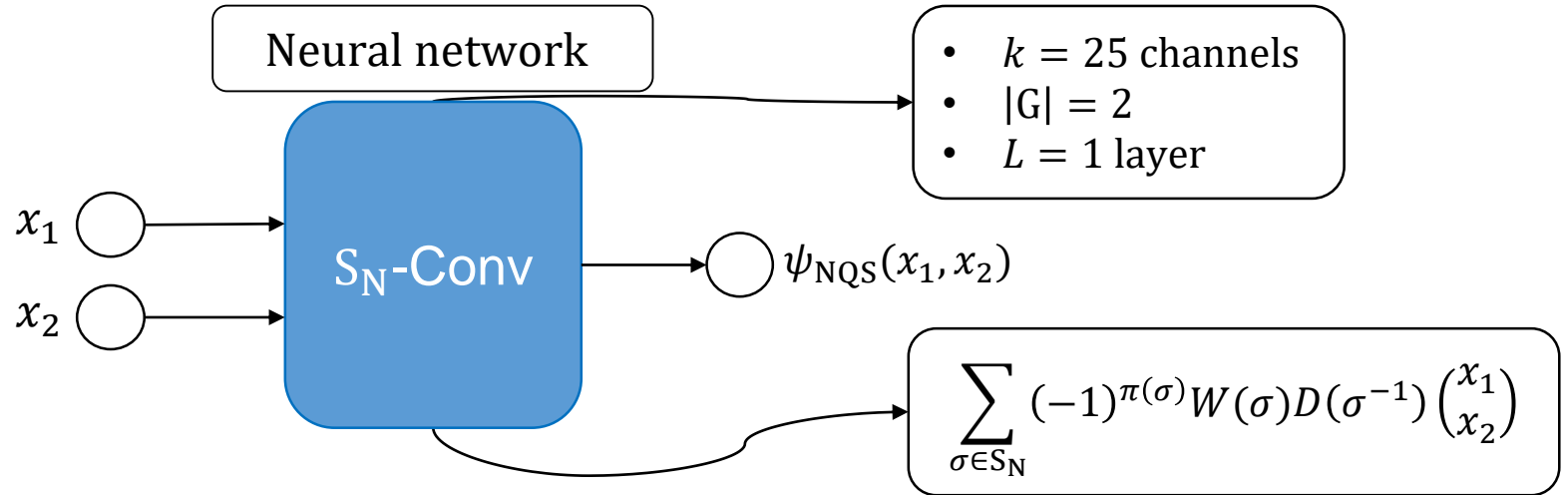
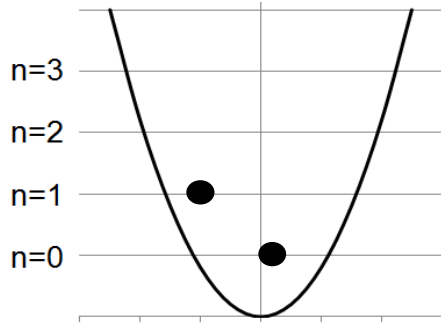
- **Group convolution**

$$\psi_{\text{NQS}}(x_1, x_2, x_3) = \sum_{\sigma \in S_N} (-1)^{\pi(\sigma)} \sigma(\psi_1(x_1)\psi_2(x_2)\psi_3(x_3))$$

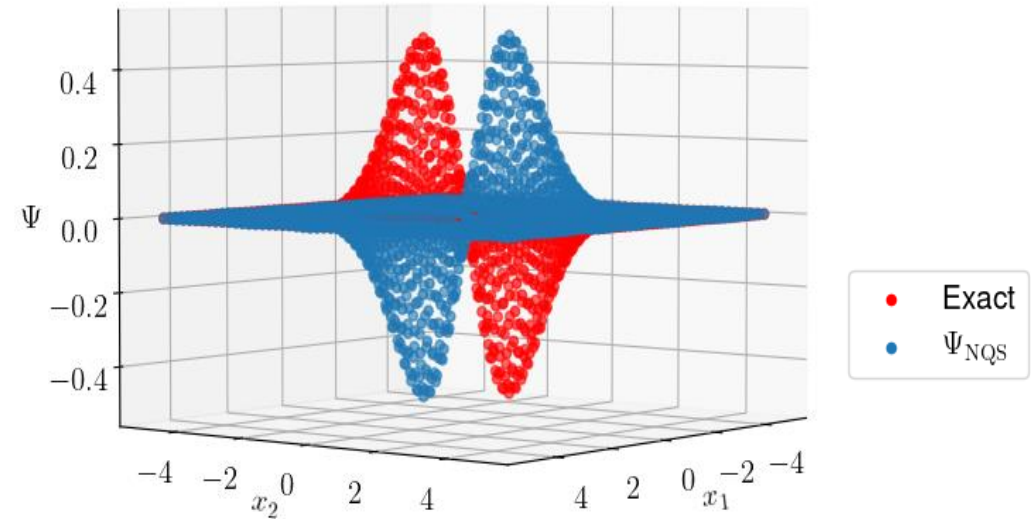
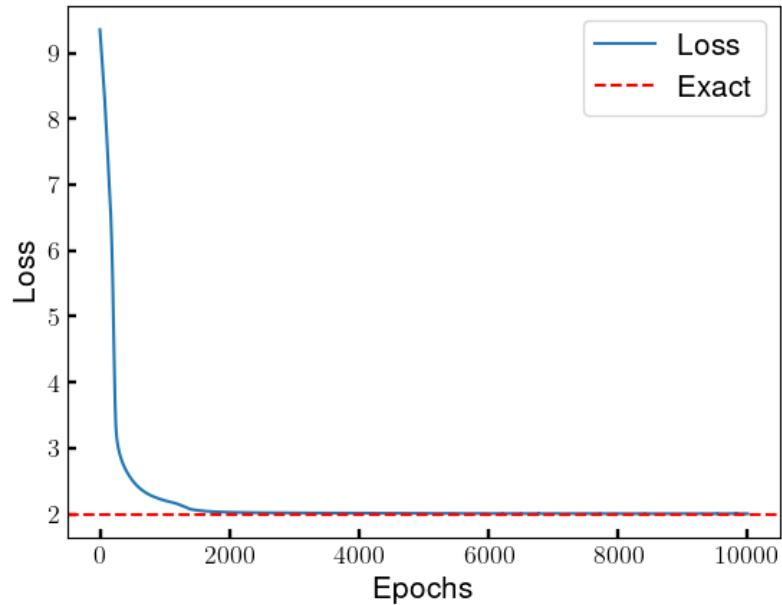
$\sigma \rightarrow$ "Alternating representation of S_N "

Toy example: $G = S_N$

$N = 2$ fermions, 1D HO

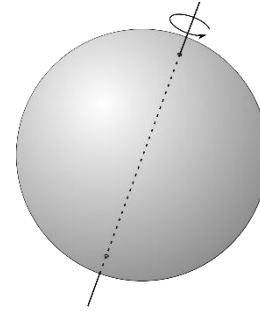


Loss evolution

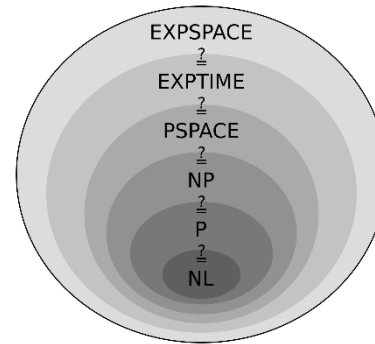


Conclusions and future outlook

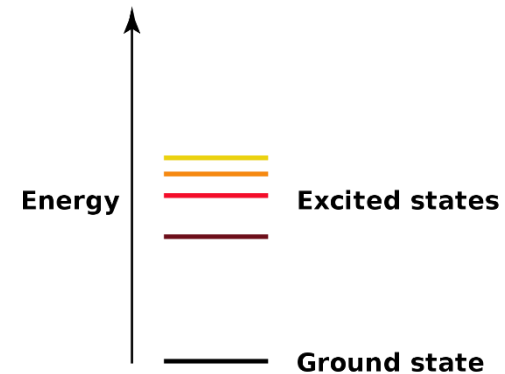
- What about continuous groups (SU(2))?



- What is the computational cost?



- Compute whole nuclear spectrum “at once”



Thank you

