



Origin of pulsar radio emission

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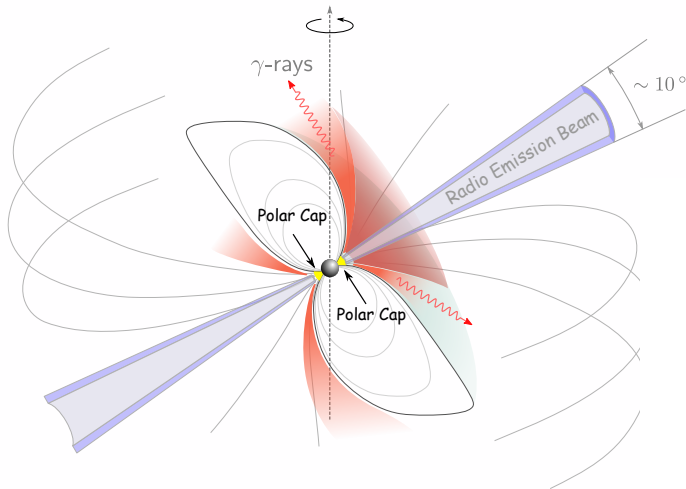
with A. Philippov (UMD) and E. Tolman (IAS)

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7th Heidelberg Symposium

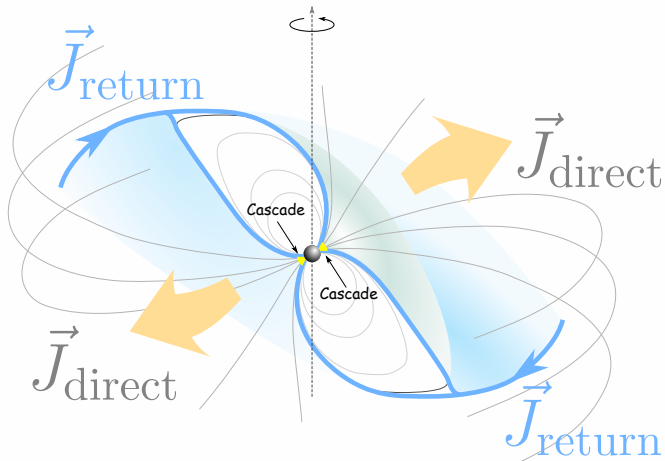
Pulsar: rapidly rotating magnet surrounded by plasma

“Electric lighthouse”



Magnetosphere is supported by electrical currents

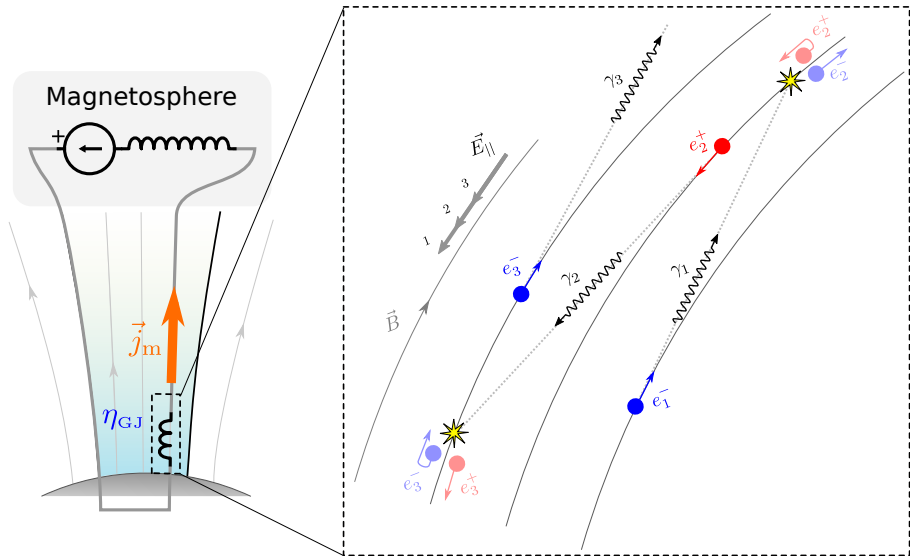
Electrical generator (NS) + Wires (plasma)



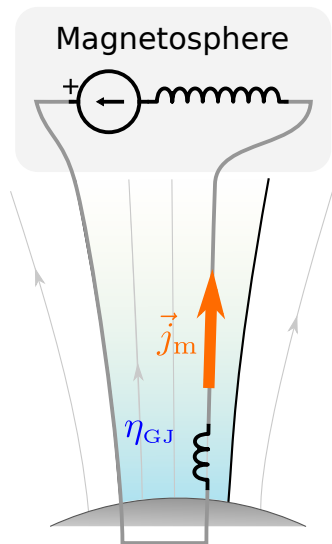
The magnetosphere is **charged**
characteristic charge density – “Goldreich-Julian” density η_{GJ} .

Plasma creation in the polar cap

Particle acceleration is regulated by pair production



Polar Cap Electrodynamics



- Rotation of the NS

$$\nabla \cdot \mathbf{E} = 4\pi(\eta - \eta_{GJ})$$

- Twist of magnetic field lines

$$\nabla \times \mathbf{B} = \frac{4\pi}{c} \mathbf{j} + \frac{1}{c} \frac{\partial \mathbf{E}}{\partial t}$$

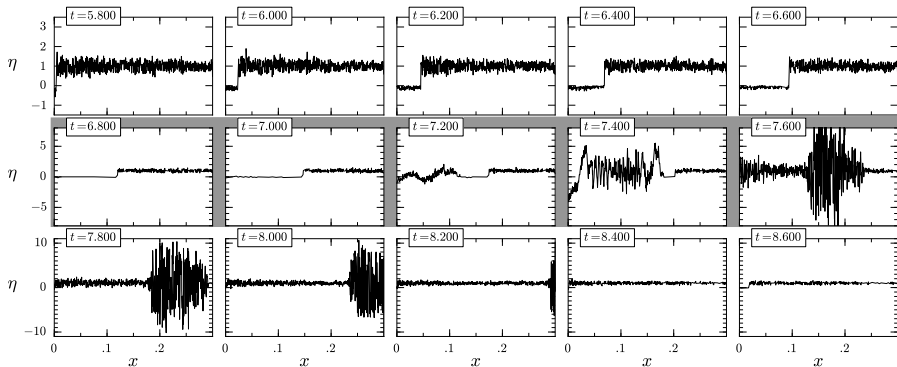
- $\mathbf{E} = 0$ if **both**

$$\eta = \eta_{GJ}$$

$$\mathbf{j} = \mathbf{j}_m \equiv \frac{c \nabla \times \mathbf{B}}{4\pi}$$

Pair creation is highly non-stationary

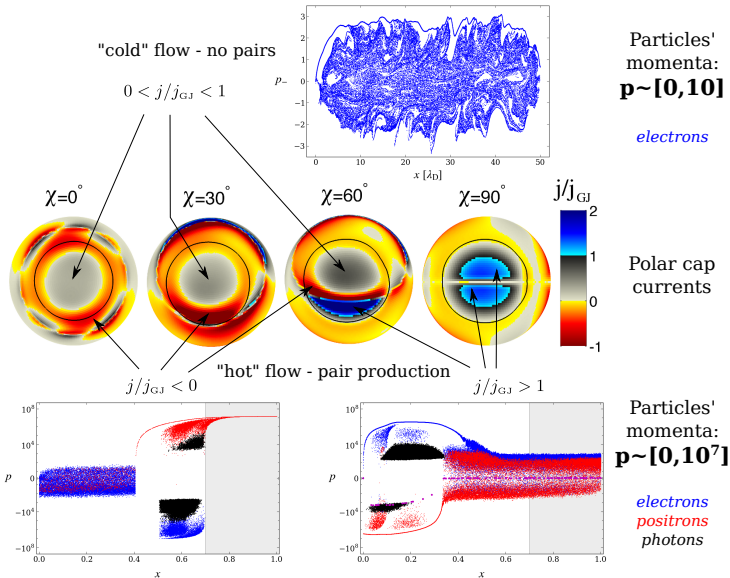
Limit cycle: series of discharges



Pair creation is non-uniform across the polar cap

Free particle extraction from the NS surface

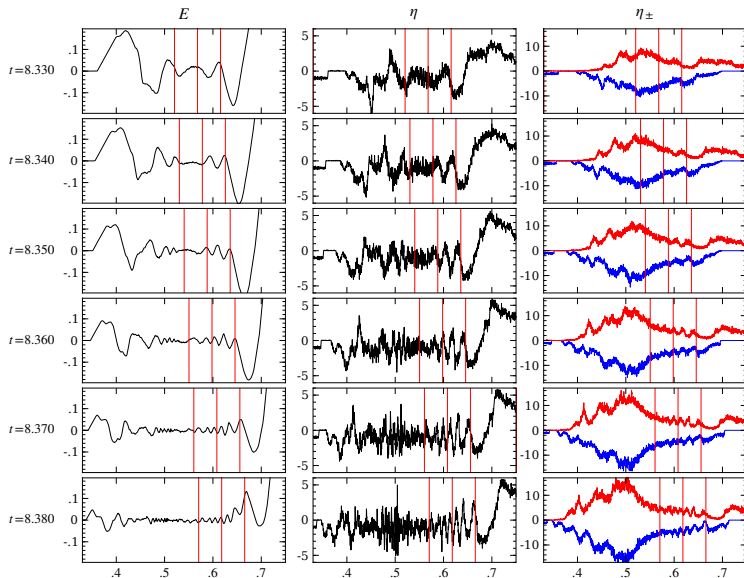
[AT & Arons'13]



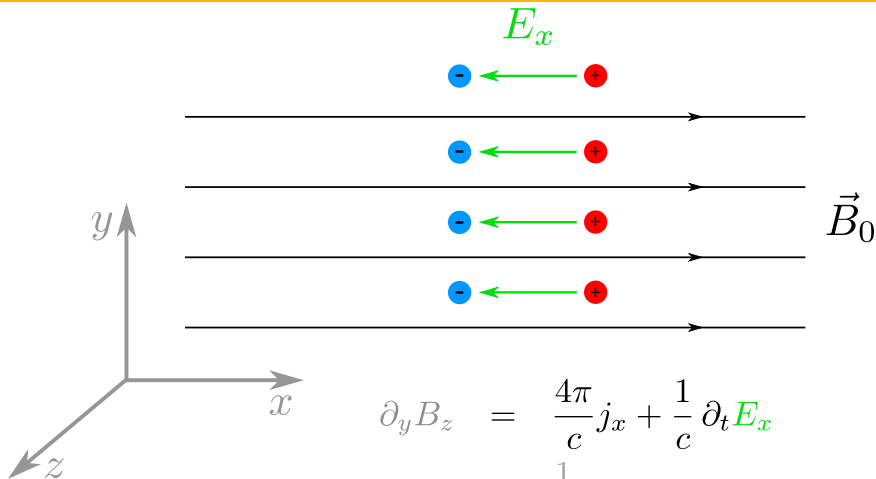
Waves during discharge: Coherent plasma motion (!)

In 1D only electrostatic Langmuir waves exist

[AT & Arons '13]



Pair formation uniform across $\vec{B}_0 \longrightarrow$ no EM mode

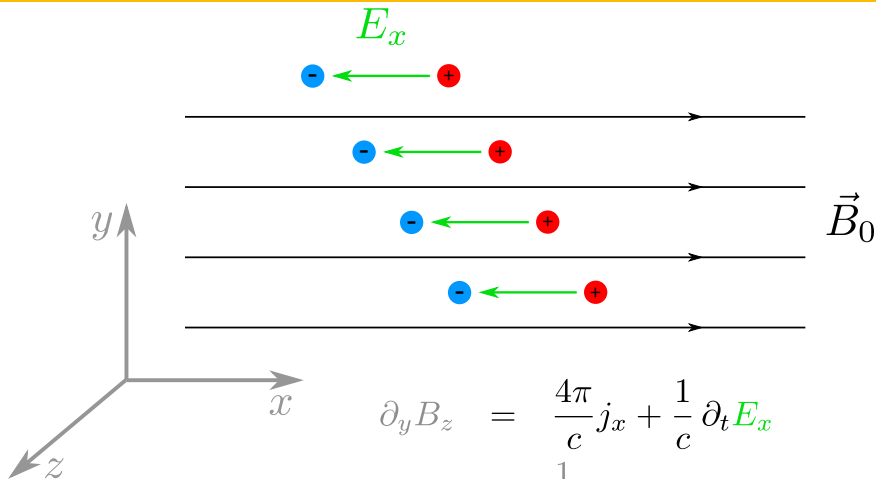


$$\partial_y B_z = \frac{4\pi}{c} j_x + \frac{1}{c} \partial_t E_x$$

$$-\partial_x B_z = \frac{1}{c} \partial_t E_y$$

$$\partial_x E_y - \partial_y E_x = -\frac{1}{c} \partial_t B_z$$

Inclined pair formation front: non-zero $\nabla \times \mathbf{E}$

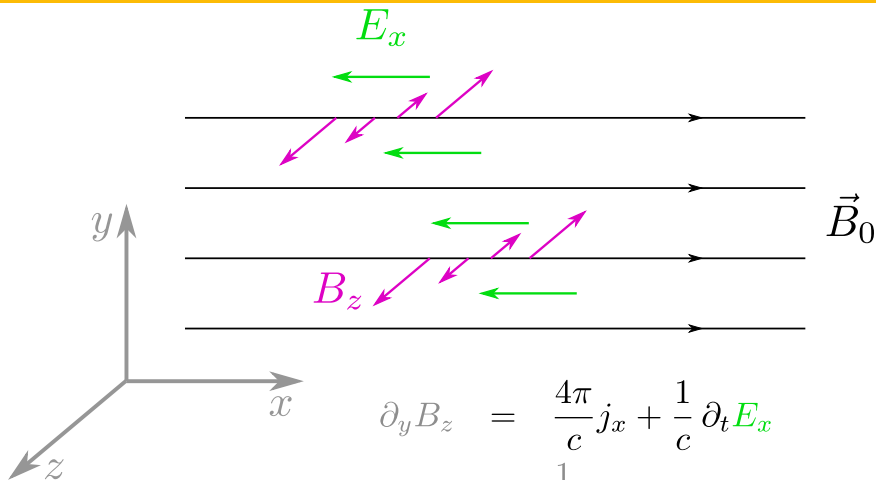


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Inclined pair formation front: $\nabla \times \mathbf{E} \longrightarrow B_z$

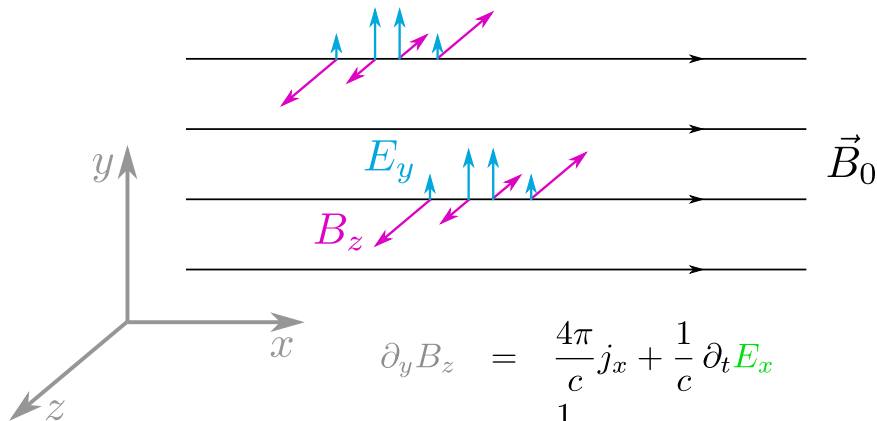


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Non-uniform B_z in x direction: $\nabla \times \mathbf{B} \longrightarrow E_y$

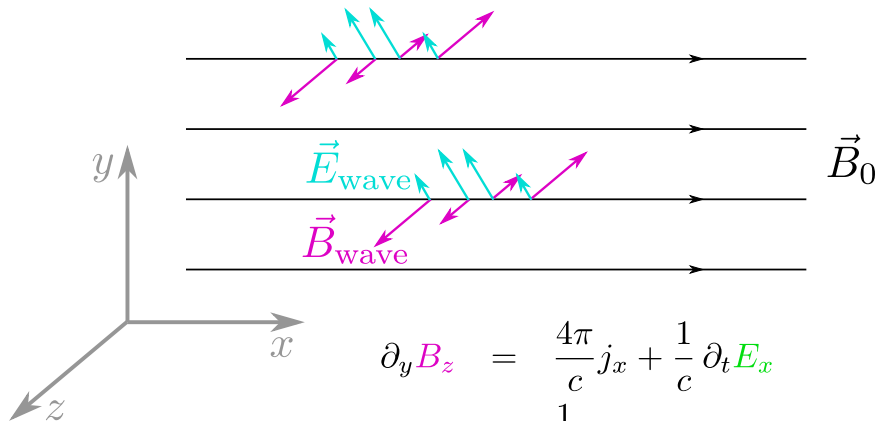


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Electromagnetic mode: $\vec{B}_{\text{wave}} \perp \vec{E}_{\text{wave}}$

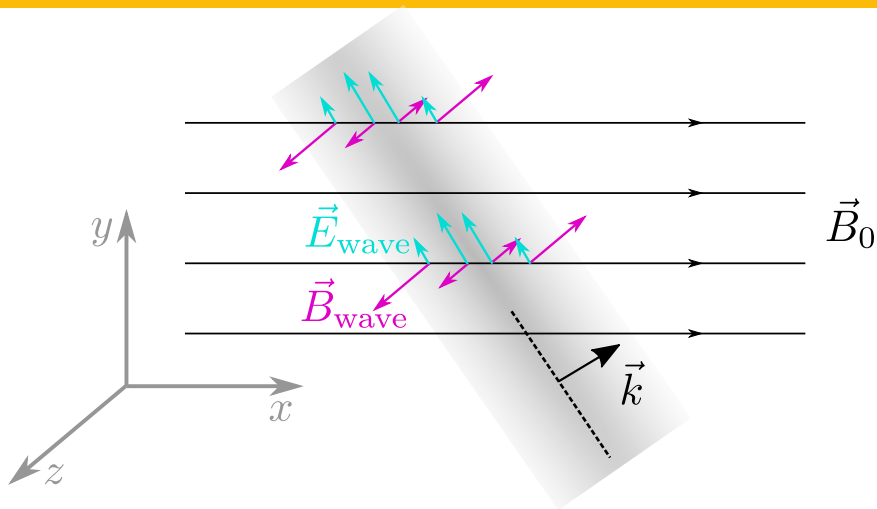


$$\partial_y B_z = \frac{4\pi}{c} j_x + \frac{1}{c} \partial_t E_x$$

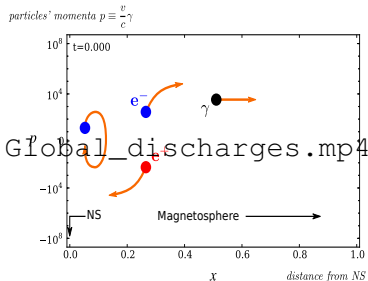
$$-\partial_x B_z = \frac{1}{c} \partial_t E_y$$

$$\partial_x E_y - \partial_y E_x = -\frac{1}{c} \partial_t B_z$$

Ordinary mode: \mathbf{E} lies in the plane ($\mathbf{k} \perp \mathbf{B}_0$)



2D PIC simulation of inclined discharge: it works!



- Oblique discharge directly excites a superluminal O-mode with non-zero k_{\perp} , which has EM component.
- As plasma density drops, these waves should become vacuum EM.
- Super-luminal O-mode can not be a result of a streaming instability.

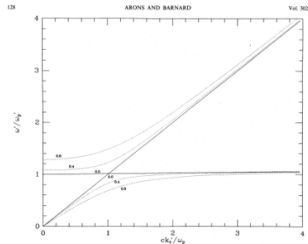
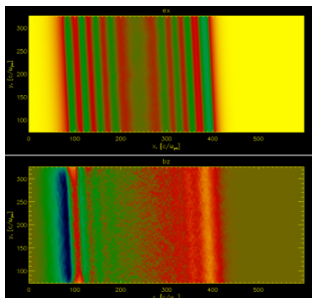


FIG. 1.—Dependence relation for a cold electron-positron plasma in the rest frame of the plasma. Each curve is for a constant value of α_0/cv_p labeled adjacent to the curve.

This mechanism generates the right kind of waves!

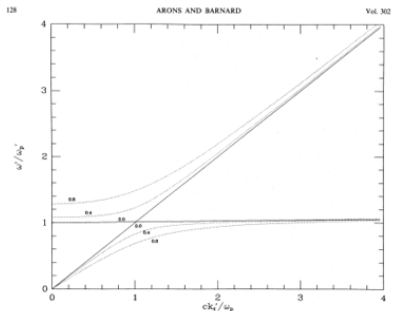
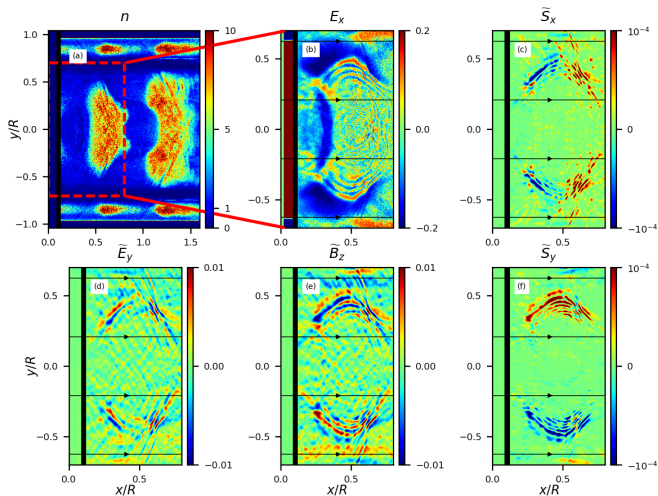


FIG. 1.—Dispersion relation for a cold electron-proton plasma in the rest frame of the plasma. Each curve is for a constant value of ω/ω_p labeled adjacent to the curve.

- Oblique discharge directly excites a superluminal O-mode with non-zero k_{\perp} , which has EM component.
- As plasma density drops, these waves should become vacuum EM.

Toy 2D PIC model of polar cap discharges

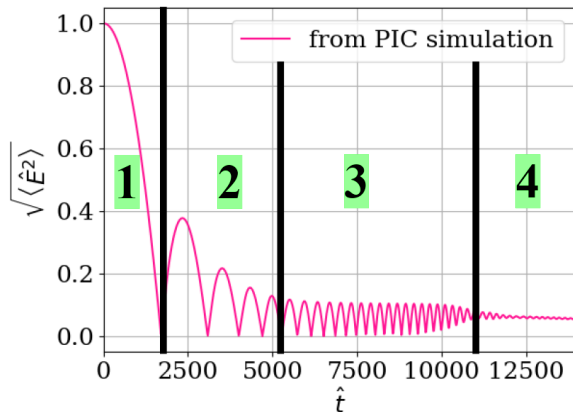
Pair formation threshold varies across parallel magnetic field lines



[Philippov, AT, Spitkovsky '20]

Damping of electrostatic wave becomes **linear**

Phases of electrostatic wave damping



- ① initial screening
- ② non-linear waves exponential damping
- ③ non-linear waves “frozen phase”
- ④ **linear damping**

[Tolman, Philippov, AT '22]

Conclusions

- Pulsars live when they produce e^\pm pairs in polar caps.
- Pair formation is non-stationary and non-uniform.
- Radio emission can be produced **directly** by the discharges
- The mechanism does not require any special conditions (besides those already existing in polar caps)
- It generates the right kind of waves
- The waves seem to survive.