On the potential of bright, young pulsars to power ultra-high gamma-ray sources

Emma de Oña Wilhelmi¹, Ruben Lopez-Coto^{2,3,} Elena Amato^{4,5} and Felix Aharonian^{6,7}

¹ DESY-Zeuthen, Germany ² INFN Padova, Italy ³ IAA/CSIC, Granada, Spain ⁴ INAF/Arcetri, Florence, Italy ⁵ DIAS, Dublin, Ireland ⁶ MPIK, Heidelberg, Germany

RESEARCH FOR HELMHOLTZ **GRAND CHALLENGES** DESY.

7th Heidelberg International Symposium on High-Energy Gamma-Ray Astronomy **Barcelona 2022**

dOW et al 2022 ApJL 930 L2













 Spectrum of Cosmic rays: Mainly protons (+electrons and heavier nuclei) coming from beyond the Solar system and extending >10 decades in energy



- Spectrum of Cosmic rays: Mainly protons (+electrons and heavier nuclei) coming from beyond the Solar system and extending >10 decades in energy
- Up to the knee, thought to be of Galactic origin, but where are these particles accelerated?



- Spectrum of Cosmic rays: Mainly protons (+electrons and heavier nuclei) coming from beyond the Solar system and extending >10 decades in energy
- Up to the knee, thought to be of Galactic origin, but where are these particles accelerated?
- Now we can see these PeV sources thanks to LHAASO Cao et al, 2021



- LHAASO discovered 12 sources extending their spectrum up to PeV energies
- First observational signature of an overall population of gamma-ray sources accelerating particles to ~PeV energies



Cao et al, 2021



| Source name | RA (°) | dec. (°) | Significance above 100 TeV (× σ) | E _{max} (PeV) | Flux at 100 TeV (CU) |
|-------------------|--------|----------|--|------------------------|----------------------|
| LHAASO J0534+2202 | 83.55 | 22.05 | 17.8 | 0.88 ± 0.11 | 1.00(0.14) |
| LHAASO J1825-1326 | 276.45 | -13.45 | 16.4 | 0.42 ± 0.16 | 3.57(0.52) |
| LHAASO J1839-0545 | 279.95 | -5.75 | 7.7 | 0.21 ± 0.05 | 0.70(0.18) |
| LHAASO J1843-0338 | 280.75 | -3.65 | 8.5 | $0.26 - 0.10^{+0.16}$ | 0.73(0.17) |
| LHAASO J1849-0003 | 282.35 | -0.05 | 10.4 | 0.35 ± 0.07 | 0.74(0.15) |
| LHAASO J1908+0621 | 287.05 | 6.35 | 17.2 | 0.44 ± 0.05 | 1.36(0.18) |
| LHAASO J1929+1745 | 292.25 | 17.75 | 7.4 | $0.71 - 0.07^{+0.16}$ | 0.38(0.09) |
| LHAASO J1956+2845 | 299.05 | 28.75 | 7.4 | 0.42 ± 0.03 | 0.41(0.09) |
| LHAASO J2018+3651 | 304.75 | 36.85 | 10.4 | 0.27 ± 0.02 | 0.50(0.10) |
| LHAASO J2032+4102 | 308.05 | 41.05 | 10.5 | 1.42 ± 0.13 | 0.54(0.10) |
| LHAASO J2108+5157 | 317.15 | 51.95 | 8.3 | 0.43 ± 0.05 | 0.38(0.09) |
| LHAASO J2226+6057 | 336.75 | 60.95 | 13.6 | 0.57 ± 0.19 | 1.05(0.16) |

DESY.



Cao et al, 2021



| Source name | RA (°) | dec. (°) | Significance above 100 TeV (× σ) | E _{max} (PeV) | Flux at 100 TeV (CU) |
|-------------------|--------|----------|--|----------------------------|----------------------|
| LHAASO J0534+2202 | 83.55 | 22.05 | 17.8 | 0.88 ± 0.11 | 1.00(0.14) |
| LHAASO J1825-1326 | 276.45 | -13.45 | 16.4 | 0.42 ± 0.16 | 3.57(0.52) |
| LHAASO J1839-0545 | 279.95 | -5.75 | 7.7 | 0.21 ± 0.05 | 0.70(0.18) |
| LHAASO J1843-0338 | 280.75 | -3.65 | 8.5 | 0.26-0.10 ^{+0.16} | 0.73(0.17) |
| LHAASO J1849-0003 | 282.35 | -0.05 | 10.4 | 0.35 ± 0.07 | 0.74(0.15) |
| LHAASO J1908+0621 | 287.05 | 6.35 | 17.2 | 0.44 ± 0.05 | 1.36(0.18) |
| LHAASO J1929+1745 | 292.25 | 17.75 | 7.4 | 0.71-0.07 ^{+0.16} | 0.38(0.09) |
| LHAASO J1956+2845 | 299.05 | 28.75 | 7.4 | 0.42 ± 0.03 | 0.41(0.09) |
| LHAASO J2018+3651 | 304.75 | 36.85 | 10.4 | 0.27 ± 0.02 | 0.50(0.10) |
| LHAASO J2032+4102 | 308.05 | 41.05 | 10.5 | 1.42 ± 0.13 | 0.54(0.10) |
| LHAASO J2108+5157 | 317.15 | 51.95 | 8.3 | 0.43 ± 0.05 | 0.38(0.09) |
| LHAASO J2226+6057 | 336.75 | 60.95 | 13.6 | 0.57 ± 0.19 | 1.05(0.16) |

DESY.



Cao et al, 2021





large source (~1 degree)

=> identification of the accelerator is complex very steep spectra

=> still compatible with the leptonic scenario (Breuhaus+2022)

Can pulsars power the PeV particles observed by LHAASO?



Cao et al, 2021



Lopez-Coto, dOW et al, 2022



Lopez-Coto, dOW et al, 2022

mechanism, powered by the rotational energy of the pulsar \dot{E}



Lopez-Coto, dOW et al, 2022

• Pulsars (and PWNe) are one of the most efficient sources in the gamma-ray regime via inverse Compton



- mechanism, powered by the rotational energy of the pulsar \dot{E}
- J1930+188, VER J2019+368...)
- The Crab Nebula is the only identified accelerator to PeV energies

• Pulsars (and PWNe) are one of the most efficient sources in the gamma-ray regime via inverse Compton

• Many of the LHAASO sources lie within likely TeV PWNe (eg. HESS J1825-137, HESS J1849-000, HESS



• Pulsars (and PWNe) are one of the most efficient sources in the gamma-ray regime via inverse Compton mechanism, powered by the rotational energy of the pulsar \dot{E}

| • | Many of the LHAASO sources lie within J1930+188, VER J2019+368) | likely Te |
|---|---|-----------------------------|
| • | The Crab Nebula is the only identified | 10 ³ |
| | accelerator to PeV energies | (s) 10 ³ |
| | | لق 10 ³⁰ × |
| | | |

1032

eV PWNe (eg. HESS J1825-137, HESS J1849-000, HESS





1. To have enough energy budget to explain the gamma-ray luminosity emitted via inverse Compton

=> Limit on the fraction of the spin-down energy transferred to high energy particles

 $W_{\rm e,\gamma} < W_{\rm e,PSR}$

The total energy in electrons responsible for IC radiation

 $W_{\rm e,\gamma} = L_{\gamma} \tau_{\rm IC}$

DESY.

The total energy made available by the pulsar in the form of gamma-ray emitting electrons

$$W_{\rm e,PSR} = \gamma_{\rm eff} \dot{\rm E} \tau_{\rm loss}$$



1. To have enough energy budget to explain the gamma-ray luminosity emitted via inverse Compton

=> Limit on the fraction of the spin-down energy transferred to high energy particles

 $W_{\rm e,\gamma} < W_{\rm e,PSR}$

The total energy in electrons responsible for IC radiation

 $W_{\rm e,\gamma} = L_{\gamma} \tau_{\rm IC}$



DESY.

The total energy made available by the pulsar in the form of gamma-ray emitting electrons

$$W_{\rm e,PSR} = \gamma_{\rm eff} \dot{\rm E} \tau_{\rm loss}$$

$$1 + \frac{\tau_{\rm IC}}{\tau_{\rm syn}} = 10^{-4} \frac{L_{\gamma,32}}{\dot{E}_{36}} (1 + 260 \ E_{e,15}^{1.7} B_{-5}^2)$$



2. To be able to reach the maximum energy observed

Acceleration carried out by electric field (expect for non-ideal cases) => Maximum will be defined by the maximum potential drop

$$\Phi_{\rm PSR} = (\dot{E}/c)^{1/2} \qquad \rightarrow \qquad E_{\rm max} = q(\Phi_{\infty} = 0)$$

ratio between the electric and magnetic strength

 $\eta_e \leq 1$ in MHD flow

 $(\dot{E}/c)^{1/2} \rightarrow E_{\text{max}} \approx 2 \eta_e \eta_B^{1/2} \dot{E}_{36}^{1/2} \text{PeV}$

fraction of pulsar wind energy flux transferred to magnetic field $\eta_B \leq 1$ by energy conservation



2. To be able to reach the maximum energy observed

Acceleration carried out by electric field (expect for non-ideal cases) => Maximum will be defined by the maximum potential drop

$$\Phi_{\text{PSR}} = (\dot{E}/c)^{1/2} \qquad \rightarrow \qquad E_{\text{max}} = q(\Phi_{\infty} = 0)$$

$$E_{\rm max} \approx 2 \dot{E}_{36}^{1/2}$$

- ratio between the electric and magnetic strength $\eta_e \leq 1$ in MHD flow $(\dot{E}/c)^{1/2} \rightarrow E_{\text{max}} \approx 2 \eta_e \eta_B^{1/2} \dot{E}_{36}^{1/2} \text{PeV}$ fraction of pulsar wind energy flux transferred to magnetic field $\eta_B \leq 1$ by energy conservation PeV
- Independent on the particle nature!



2. To be able to reach the maximum energy observed

Acceleration carried out by electric field (expect for non-ideal cases) => Maximum will be defined by the maximum potential drop

$$\Phi_{\rm PSR} = (\dot{E}/c)^{1/2} \qquad \rightarrow \qquad E_{\rm max} = q(\Phi_{\infty} = 0)$$

$(\dot{E}/c)^{1/2} \rightarrow E_{\text{max}} \approx 2 \eta_e \eta_B^{1/2} \dot{E}_{36}^{1/2} \text{PeV}$



2. To be able to reach the maximum energy observed

Acceleration carried out by electric field (expect for non-ideal cases) => Maximum will be defined by the maximum potential drop

$$\Phi_{\rm PSR} = (\dot{E}/c)^{1/2} \rightarrow E_{\rm max} = q(\dot{E}/c)^{1/2} \rightarrow E_{\rm max} \approx 2 \ \eta_e \ \eta_{\rm B}^{1/2} \ \dot{E}_{36}^{1/2} \ {\rm PeV}$$

$$\Phi_{\infty} = 0$$

If electrons AND >100 TeV

photon field up scattered is 2.7K CMB => $E_{\rm e}$ =

$$E_{\gamma \max} \approx 0.9 \ \eta_e^{1.3} \ \eta_B^{0.65} \ \dot{E}_{36}^{0.65}$$
 PeV

$$\simeq 2.15 \ E_{\gamma,15}^{0.77} \ \text{PeV}$$

$$E_{\gamma \max} \approx 0.9 \dot{E}_{36}^{0.65}$$
 PeV





If electrons AND >100 TeV

photon field up scattered is 2.7K CMB => E_{e} \simeq

$$E_{\gamma \max} \approx 0.9 \ \eta_e^{1.3} \ \eta_B^{0.65} \ \dot{E}_{36}^{0.65} \ \text{PeV}$$

DESY.

field (expect for non-ideal cases) e maximum potential drop

$$\dot{E}/c)^{1/2} \rightarrow E_{\rm max} \approx 2 \ \eta_e \ \eta_{\rm B}^{1/2} \ \dot{E}_{36}^{1/2} \ {\rm PeV}$$

$$\simeq 2.15 \ E_{\gamma,15}^{0.77} \ \text{PeV}$$

$$E_{\gamma \max} \approx 0.9 \dot{E}_{36}^{0.65}$$
 PeV



2. To be able to reach the maximum energy observed It is not sufficient to be able to accelerate the particles => we need to overcome the radiation loses:





$$4 \times 10^9 E_{e,15}^{-1} B_{-5}^{-2}$$
 s

$$max \approx 5\eta_e^{0.65}B_{-5}^{-0.65}$$
 PeV



2. To be able to reach the maximum energy observed It is not sufficient to be able to accelerate the particles => we need to overcome the radiation loses:

$$E_{\gamma \max} \approx 0.9 \ \eta_e^{1.3} \ \eta_B^{0.65} \ \dot{E}_{36}^{0.65} \ \text{PeV}$$

 $E_{\gamma \max} \approx 5 \eta_e^{0.65} B_{-5}^{-0.65} \ \text{PeV}$

The limit on the energy will be given by the interplay between the two expression above







• Search for pulsars (ATNF catalog) within 1 deg from the location of LHAASO sources.





- Search for pulsars (ATNF catalog) within 1 deg from the location of LHAASO sources.

• Select young ($\tau < 10^6$ yrs) and energetic ($\dot{E}/d_{\rm kpc}^2 > 10^{34}$ erg/s/kpc² or $\dot{E} > 10^{36}$ erg/s) pulsars

- Search for pulsars (ATNF catalog) within 1 deg from the location of LHAASO sources.
- Select young ($\tau < 10^6$ yrs) and energetic ($\dot{E}/d_{\rm kpc}^2 > 10^{34}$ erg/s/kpc² or $\dot{E} > 10^{36}$ erg/s) pulsars
- Obtain the maximum energy and compare with the observations

- Search for pulsars (ATNF catalog) within 1 deg from the location of LHAASO sources.
- Select young ($\tau < 10^6$ yrs) and energetic ($\dot{E}/d_{\rm kpc}^2 > 10^{34}$ erg/s/kpc² or $\dot{E} > 10^{36}$ erg/s) pulsars
- Obtain the maximum energy and compare with the observations





dOW, Lopez-Coto et al, 2022

| LHAASO Source | Pulsar | $E_{\gamma max}$ | E _{max} | B _{max} |
|---------------|------------|------------------|------------------|------------------|
| | | [PeV] | [PeV] | [µG] |
| J1825-1326 | J1826-1256 | 2.06 | 3.79 | 38 |
| | B1823-13 | 1.77 | 3.35 | 14 |
| J1839-0545 | J1837-0604 | 1.44 | 2.83 | 33 |
| | J1838-0537 | 2.78 | 4.90 | ≫100 |
| J1843-0338 | J1841-0345 | 0.41 | 1.04 | 12 |
| | J1844-0346 | 2.25 | 4.10 | ≫100 |
| J1849-0003 | J1849-0001 | 3.71 | 6.26 | ≫100 |
| J1908+0621 | J1907+0602 | 1.77 | 3.35 | 30 |
| | J1907+0631 | 0.63 | 1.46 | 9 |
| J1929+1745 | J1925+1720 | 0.91 | 1.95 | 9 |
| | J1928+1746 | 1.26 | 2.53 | 14 |
| J1956+2845 | J1954+2836 | 0.94 | 2.00 | 37 |
| | J1958+2846 | 0.47 | 1.17 | 22 |
| J2018+3651 | J2021+3651 | 1.99 | 3.69 | 102 |
| J2032+4102 | J2032+4127 | 0.28 | 0.77 | 7 |
| J2108+5157 | | | | |
| J2226+6057 | J2229+6114 | 5.89 | 9.38 | 64 |

- The majority of the source can be, in principle, powered by energetic pulsars in a 1degree region
- We derive upper limits to the magnetic field (the efficiency provides the most constraining limit)

$$\gamma_{\text{eff}} = \frac{L_{\gamma}}{\dot{E}} \left(1 + \frac{\tau_{\text{IC}}}{\tau_{\text{syn}}} \right) = 10^{-4} \frac{L_{\gamma,32}}{\dot{E}_{36}} (1 + 260 \ E_{e,15}^{1.7} B_{-5}^2)$$



Summary

- => Two of the LHAASO sources cannot be explained by pulsars in the FoV
- For young pulsars with magnetic field of $\sim 100 \ \mu$ G, the maximum energy is constraint by the synchrotron losses, whereas for older ones (~few μ G) by the potential drop
- (N157B) in the LMC could also be an efficient accelerator.
- The synchrotron counterpart of the 100 TeV IC nebula should be bright in the X-ray domain =>eROSITA
- More constraints should come from the size of the gamma-ray emission (in preparation)

DESY.

• Bright pulsars can in theory accelerate particles (electrons and protons) to the observed PeV energies

• We constrain the maximum photon energy in Geminga to less than 200 TeV, whereas the Crab twin



The search of sources accelerated particles up to PeV energies continues https://indico.desy.de/event/34265/



DESY.

ST PeVatrons and their environments

Enter your search term

Q

Chairs

Ruben Lopez-Coto (IAA-CSIC) Emma de Oña Wilhelmi (DESY Zeuthen)

Scientific Organizing Committee

Felix Aharonian (DIAS / MPIK) Elena Amato (INAF / Osservatorio Astrofisico di Arcetri) Tony Bell (University of Oxford) David Berge (DESY Zeuthen) Juan Cortina (CIEMAT) Stefano Gabici (APC) Petra Huentemeyer (Michigan Tech University) Jim Hinton (MPIK) Sarah Recchia (INFN Torino) Roberta Zanin (CTA Observatory) Cao Zhen (IHEP)

LHAASO J2032+4102





- Compact 0.15° TeV PWN
- Modulated gamma-ray emission => Binary system
- In the Cygnus region => Cocoon?

LHAASO J2108+5157

LHAASO Col., 2021



- No counterparts yet with ground-based instruments
- Detected above 25 TeV
- Relatively compact ~0.26d source





Pulsars Maximum energy **Hillas criterium**

What is the maximum energy a particle can reach in a Galactic source? Acceleration is always (expect for non-ideal cases) carried out by an electric field For a particle with charge q, moving a distance L

- $E = q | \vec{E} | L$
- We can define the acceleration efficiency as: • $\eta = \overrightarrow{B} / \overrightarrow{E}$

then:

 $E = q\eta BL$

- L = Size of the source
- B = Magnetic field in the source



This two value will define the maximum energy

Pulsars Maximum energy Hillas criterium

The maximum energy:

 $E_{\rm max} = q\eta_e B_{\rm TS} R_{\rm TS},$

The magnetic density is a fraction of the pulsar wind energy flux:

$$\frac{B_{\rm TS}^2}{8\pi} = \eta_{\rm B} \frac{\dot{E}}{(4\pi R_{\rm TS}^2 c)}$$

Then the Emax:

$$E_{\rm max} \approx 2 \,\eta_e \,\eta_{\rm B}^{1/2} \,\dot{E}_{36}^{1/2} \,\,{\rm PeV}$$



In astrophysical environment :

 $\overrightarrow{E} = v/c\overrightarrow{B}$ if $\overrightarrow{E} \sim \overrightarrow{B} =$ > relativistic plasmas: PWNe

FOR PSRs Φ IS MEASURED DIRECTLY:



Amato et al, 2012

