High-mass gamma-ray binaries as very efficient accelerators

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Gamma 2022 Symposium Barcelona

July 7th, 2022



2 Most known HMGB reach \sim 10 TeV

Particle acceleration in non-accreting HMGB

4 Concluding

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Introduction

- 2) Most known HMGB reach \sim 10 TeV
- Particle acceleration in non-accreting HMGB

4 Concluding

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High-mass gamma-ray binaries (with CO)

- The phenomenological term Gamma-ray binary usually means star+CO and SED dominance of gamma rays (without the star).
- High-mass gamma-ray binaries (HMGB) are among the most powerful galactic sources: $L \sim 10^{36-37}$ (MeV), 10^{34-37} (GeV) and 10^{32-35} erg s⁻¹ (TeV).
- The great majority of the known HMGB are VHE emitters.

(Some reviews: Mirabel 2006; Romero 2009; B-R & Khangulyan 2009; Dubus 2015; Paredes & Bordas 2019; Chernyakova & Malyshev 2020...)



Main elements of a HMGB.



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Main elements of a HMGB.



Non-accreting (pulsar) vs accreting (HMMQ) HMGB

- A non-accreting HMGB consists of a young pulsar plus an OB star whose winds interact.
- A HMMQ consists of a CO plus an OB star in which the wind is accreted and jets form, which interact with the wind.
- In both cases, outflows interacting along the orbit are complex and emit radio, X- and gamma rays, likely through synchrotron and IC, plus γγ...
- (e.g., B-R, Khangulyan, Aharonian, Barkov, Perucho + ...;
- Bogovalov, + ...; Romero, + ...; Dubus, Lamberts, Cerutti + ...;
- Sierpowska-Bartosik, Torres, + ...; Bednarek, + ...; Reitberger,
- Reimer, Huber, Kissmann; Yoon, Heinz, + ...; Chernyakova,
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- (see e.g. Molina's talk -Wednesday- and Kefala's poster)







High-mass microquasar (Barkov & B-R 2021)

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High-mass microquasar (Barkov & B-R 2021)

Non-accreting HMGB



Pseudocolor Var: rho 3.0e+04 Var: Speed 0.98 1.6e+02 0 74 0.88 0.49 0.0048 0.25 2.6e-05 0.0011 30 10 30

(Zabalza, B-R, et al. 2013)

RHD simulations with PLUTO of 2-wind-orbit interactions (low *e*).

Fig. 2. Representation of the distribution of density in the XY-, XZ-, and YZ-planes for 3Dlf at t = 3.9 days (apastron). Streamlines are show in 3D.

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(LS 5039 at apastron; B-R, Barkov & Perucho 2015;see also Kissmann's talk from Wednesday)

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High-mass microquasar

RHD simulations with PLUTO of jet-wind-orbit interactions.



(HMMQ jet in a e = 0 orbit; Barkov & B-R 2021)



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VHE spectrum of LS I +61 303

Eccentric, relatively compact Be+CO? binary (end of 70s)



Figure 3: Spectral energy distribution (SED) for LS I +61°303 for two parts of the orbit (parts of the orbit shown on top panels). SED on the *left* is near apastron passage covering $\phi = 0.5 \rightarrow 0.8$ and SED on the *right* is for the rest of the orbit for $\phi = 0.8 \rightarrow 0.5$. The orbital parameters shown on top panel are used from [14]

(VERITAS: Kar et al. 2017)

Highly eccentric, wide Be+pulsar binary (beginning of 90s)



(HESS: Bordas et al. 2015)

Moderately eccentric, compact O+CO? binary (90s)



FIGURE 4. Left: SEDs obtained from monoscopic and a stereoscopic analyses of the H.E.S.S.-II and H.E.S.S.-I data sets, respectively. Results of fits with power-law functions are given in the inset. Also an SED obtained from a re-analysis of Fermi-LAT data is shown. *Right*: SEDs resulting from H.E.S.S.-I analyses for parts of the orbit corresponding to the inferior or superior conjunction. The corresponding orbital phase ranges are given for reference. Fit results are given in the main text.

(HESS: Bordas et al. 2015; also detected by HAWC)

VHE spectrum of HESS J0632+057

Eccentric, rather wide Be+CO? binary (00s)



Figure 7. Differential energy spectra of photons above 200 GeV obtained by H.E.S.S., MAGIC and VERITAS averaged over all available orbits. The figure shows the results for four different orbital phase bins: (a) orbital phases 0.2-0.4; (b) orbital phases 0.4-0.6; (c) orbital phases 0.6-0.8; (d) orbital phases 0.8-0.2. Vertical error bars show 1σ uncertainties; downwards pointing arrows indicate upper limits at the 95% confidence level.

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VHE spectrum of 1FGL J1018.6–5856

Moderately eccentric?, relatively compact O+CO? binary (10s)



(HESS 2015)

Fig. 1. SED of HESS J1018–589 A/1FGL J1018.6–5856 is shown in black (filled squares and circles for the LAT and HESS detection). For comparison, the SEDs of LS 5039 during superior (SUPC) and inferior conjunction (INFC) are also included (blue points from Hadasch et al. 2012; Aharonian et al. 2005a).

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Image: A matrix a

VHE spectrum of PSR J2032+4127

Extremely eccentric, very wide Be+pulsar binary (00s or 10s?)



Figure 3. Spectral energy distributions for PSR J2032+4127/MT91 213 and TeV J2032+4130 from VERITAS (eff) and MAGIC (right). The blue butterflises are the spectral fits to TeV J2032+4130. The red butterflies in the upper plots are fits to the 2017 fail data: the sym a power-law fit to TeV J2032+4130 and a cutoff powerlaw fit to PSR J2032+4127/MT91 213. In the bottom plots, orange is the fit to the low-state data (PSR J2032+4127/MT91 213) is fit with a cutoff), while green represents the high-state data (PSR J2032+4127/MT91 213 is fit with a power law). The fit parameters are given in Table 1 and the time periods are defined in the text.

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VHE spectrum of LMC P3

Moderately eccentric, compact O+CO? binary (10s)



(HESS 2018)

Fig. 3. Spectral energy distribution averaged over the full orbit (green, squares) and for the on-peak orbital phase range (orbital phase from 0.2 to 0.4: blue, circles). The data points have 1σ statistical error bars, upper limits are for a 95% confidence level. The best fit and its uncertainty are represented by the solid lines and shaded areas, respectively.

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• The source is rather similar to LS 5039 and 1FGL J1018.6–5856.

It is not known so far if this source emits VHE.

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(see Corbet et al. 2019)

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VHE spectrum of HESS J1832–093

Moderately wide O?+CO? binary? (10s)



(Martí-Devesa & Reimer 2020 ↑, Tam et al. 2019, and ref. therein; HESS 2015, Eger et al. 2016)

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- Ultrarelativistic weak-B' flow.
- Perpendicular/oblique shocks of different speeds and strength (B'?).
- Flow reacceleration and further shocks, shear layers, turbulence and mass-loading... (*B*'?)

(e.g., Rieger et al. 2007; Khangulyan et al. 2008; B-R & Khangulyan 2009; Takahashi et al. 2009 B-R 2012; B-R & Rieger 2012, Derishev & Aharonian 2012)

(B-R et al. 2012 -2D-, 2015 -3D-; low e) \rightarrow



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(Barkov & B-R 2018, high e)

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Flow termination.(B-R & Barkov 2011) Mainly X-ray evidence.

(Paredes+2007 -LSI-; Durant+2011 -LS-; Pavlov+2015 -PSRB-, Williams+2015 -1FGL-; Kargaltsev+2021 -HESS-; Albacete-Colombo+2020 -PSRJ- ↓)



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• $E_{\rm max}$ for the most relevant processes ($t_{\rm acc} \sim \eta E/qBc$; $D \sim \chi D_{\rm Bohm}$; $RB \sim ct$?):

- Hillas limit (e^{\pm}, p) :
- Escape/adiabatic cooling (e^{\pm}, p) :
- Diffusion (e^{\pm} , p):
- Synchrotron (e^{\pm}):

(e.g., Rieger et al. 2007; Khangulyan et al. 2008; B-R & Khangulyan

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 - Hillas limit (e^{\pm} , p): $E_{\text{max}}^{\text{H}} \sim 300 R_{12} B_0$ TeV
 - Escape/adiabatic cooling (e^{\pm} , ρ): $E_{\text{max}}^{\text{dy}} \sim 90 R_{12} B_0 v_{10}^{-1} \eta_1^{-1} \text{ TeV}$
 - Diffusion (e^{\pm} , p): $E_{\text{max}}^{\text{diff}} \sim 40 R_{12} B_0 \eta_1^{-1/2} \chi_1^{-1/2}$ TeV
 - Synchrotron (e^{\pm}): $E_{\max}^{sy} \sim 20 \eta_1^{-1/2} B_0^{-1/2} Te^{-1/2}$

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Gamma-ray binaries as efficient accelerators

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Energy [MeV] ^Q C July 7th, 2022 21/24

10⁶ 10⁷

*E*_{max} for the most relevant processes (*t*_{acc} ~ η*E*/*qBc*; *D* ~ χ*D*_{Bohm}; *RB* ~ct?):

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- Derishev & Aharonian (2012) showed that the converter mechanism (Derishev et al. 2003; Stern 2003) can operate in compact HMGB via e[±]-creation in the unshocked pulsar wind.
- The wind loads and slows down while providing a Γ²-boost to the new e[±], which cool little due to the KN effect (IC) and B' ≈ B/Γ (sync.) until reaching the shock. (Derishev & Aharonian 2012)
- This mechanism may solve several misteries in LS 5039, and perhaps also in LS I +61 303, 1FGL J1018.6–5856...?

(Khangulyan et al., B-R et al. 2008; Cerutti et al. 2008; Collmar, W.; B-R 2021...)

- Pairs accelerate close to t_{acc} = E/qBc, with γ_{peak} ~ Γ² ~ 10⁸.
- A 1–30 MeV synchrotron (seen) component would arise naturally from postshock synchrotron.
- The (unseen) unshocked pulsar wind SED component would be also smoothed out.
- External cascading effects are closely linked to this process.

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• HMGB are perfect sites for multi-TeV particle acceleration and radiation.

- However, leptonic and hadronic CR injection from the system may be inefficient due to adiabatic losses.
- On the other hand, large-scale outflow-medium interactions might be a suitable site for PeV CR production (large scale X-rays).
- Finally, as LS 5039, LS I 61 040, 1FGL J1018.6–5856, 4FGL J1405.1-6119... are known within $\sim 1/4$ of the disk.
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