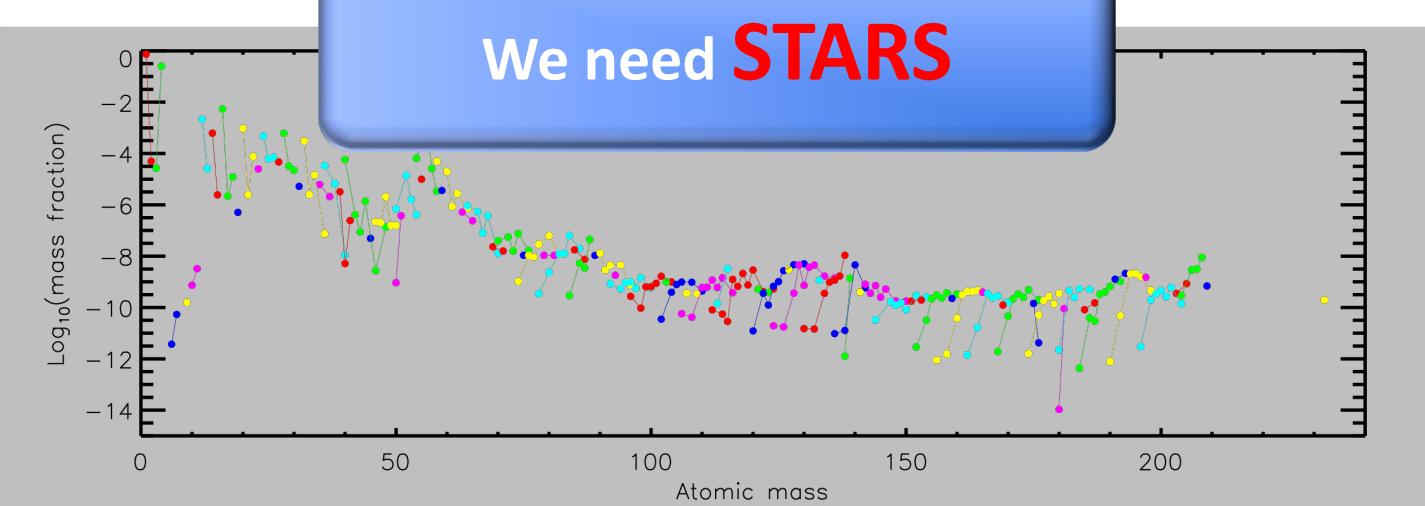
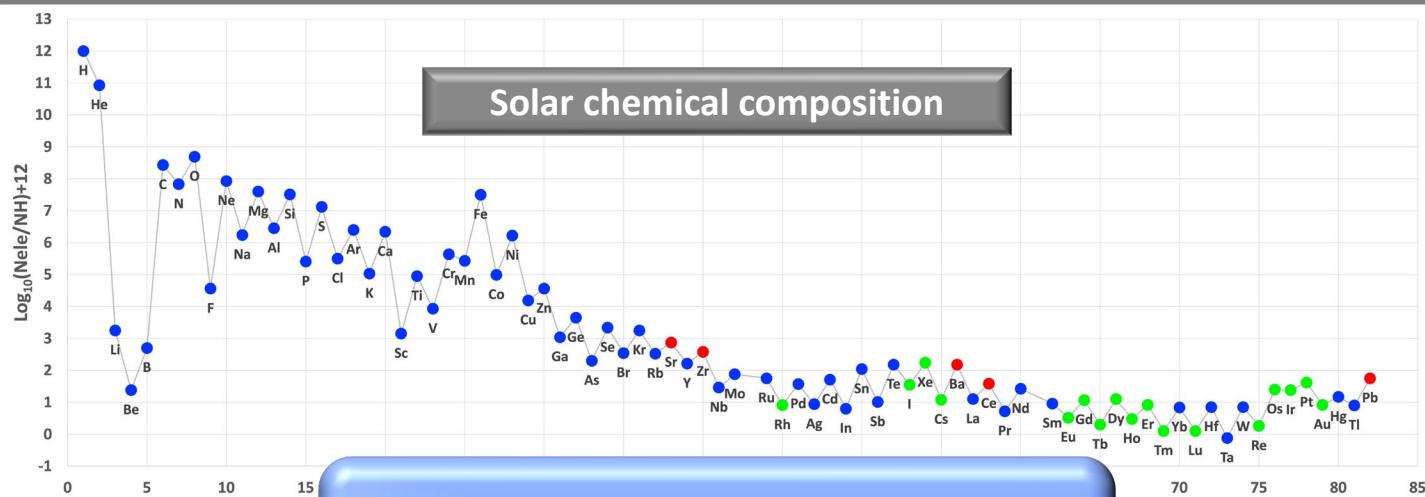




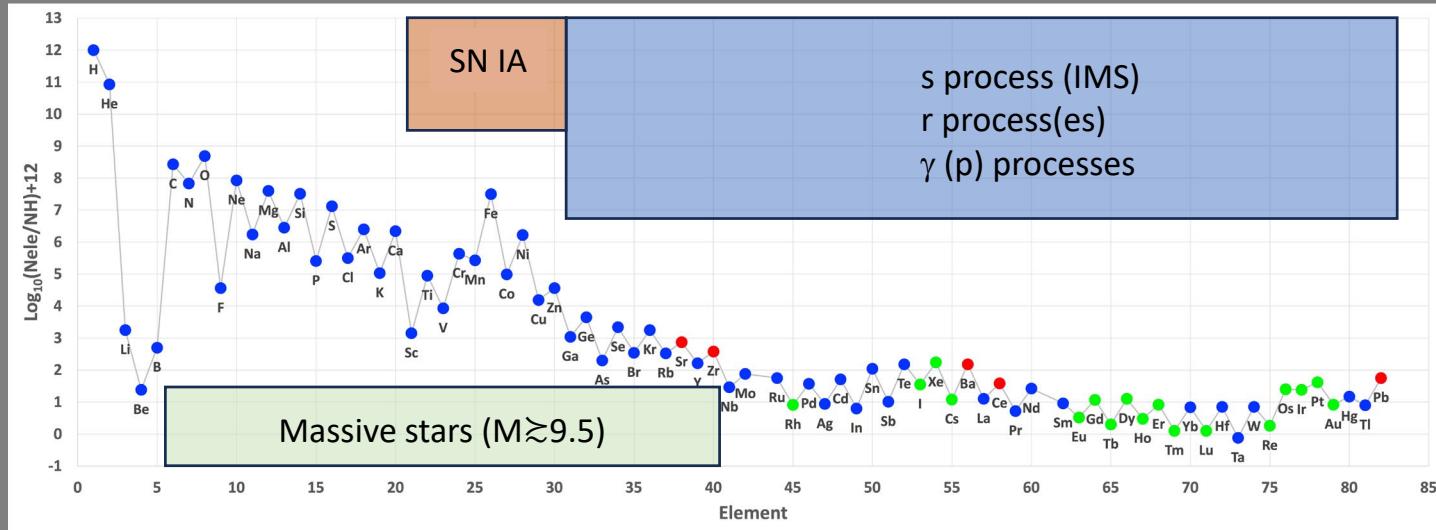
Nuclei in the Cosmos School
2025

STELLAR EVOLUTION

Alessandro Chieffi
INAF, Italy



Solar chemical composition



Golden rules

The temperature in the inner core of a star increases only as a consequence of the gravitational contraction

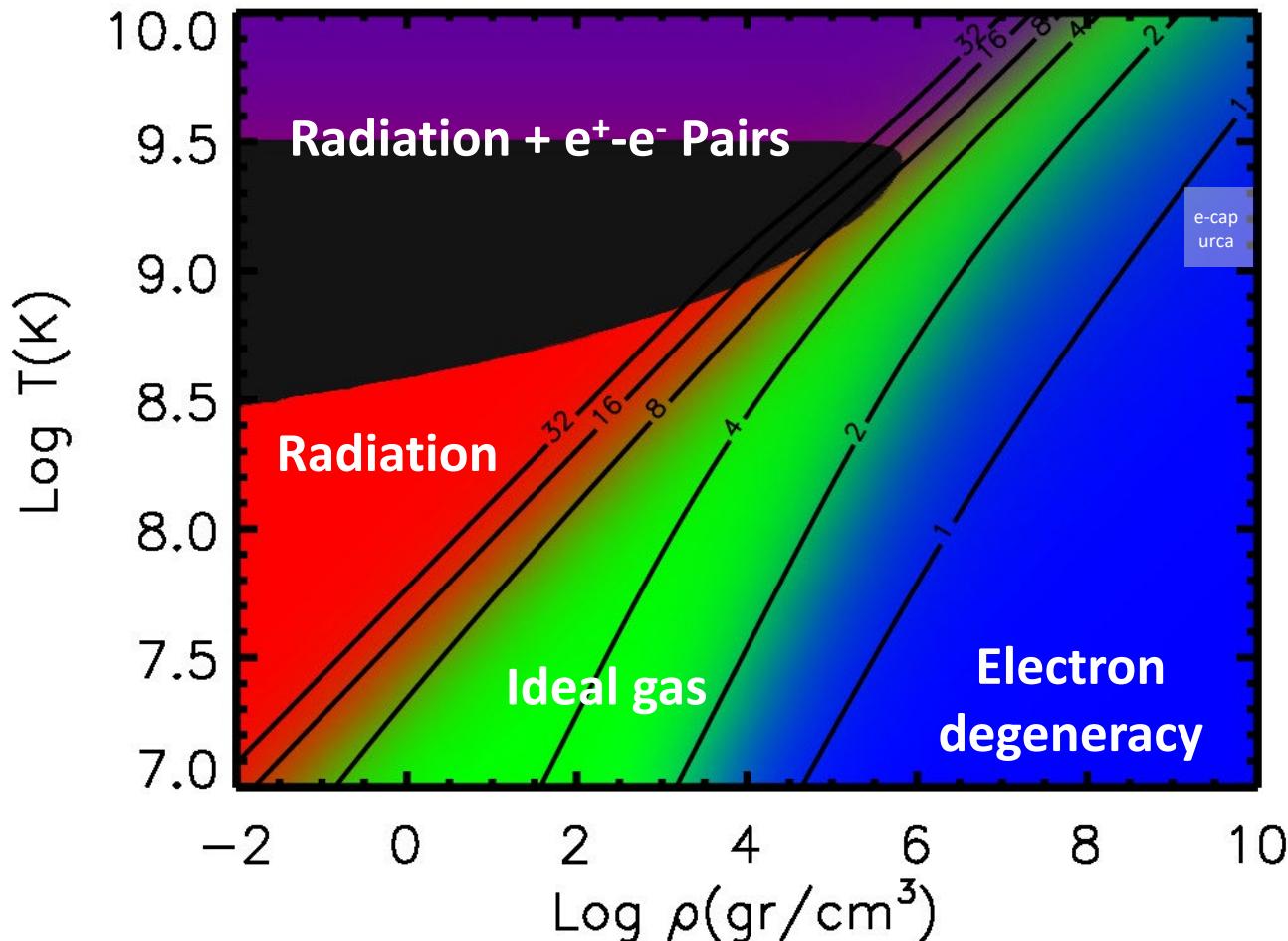
According to the Virial theorem, a fraction of the energy gained by the contraction leads to the increase of the temperature while the remaining part is lost outward (the Luminosity)

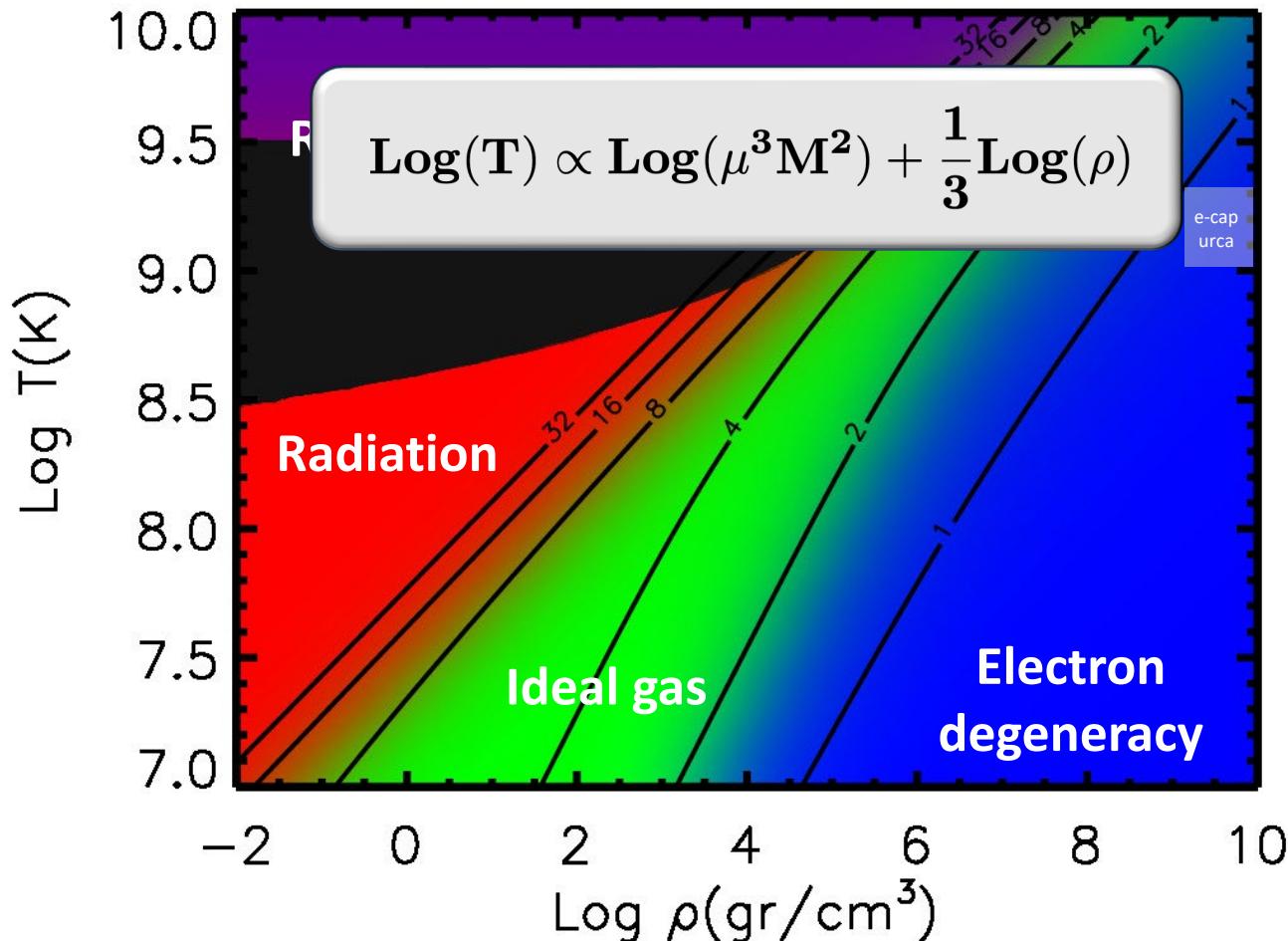
$$\text{Log}(T) \propto \text{Log}(\mu^3 M^2) + \frac{1}{3} \text{Log}(\rho)$$

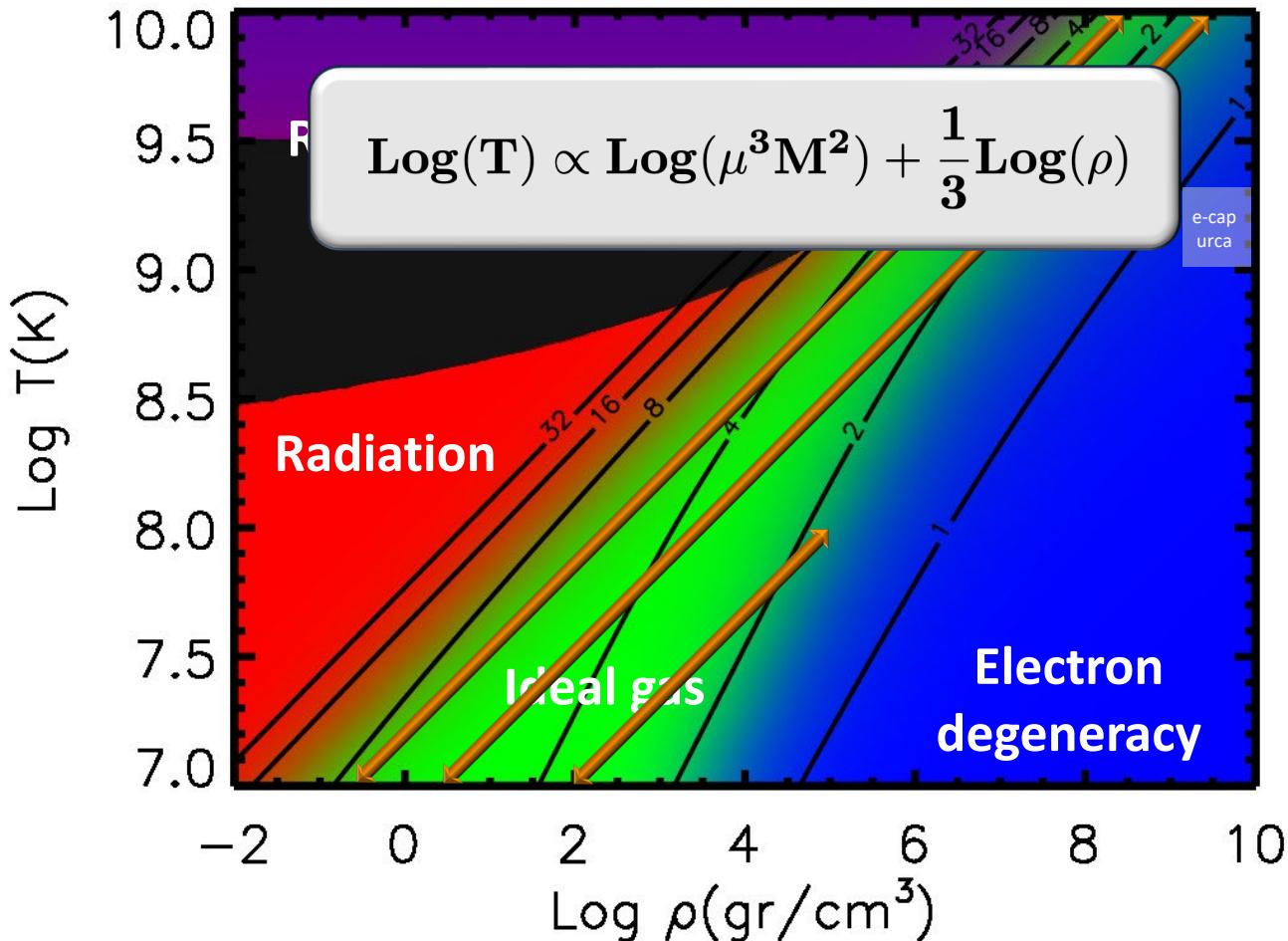
Hydrostatic equilibrium + EOS (Perfect Gas + radiation)

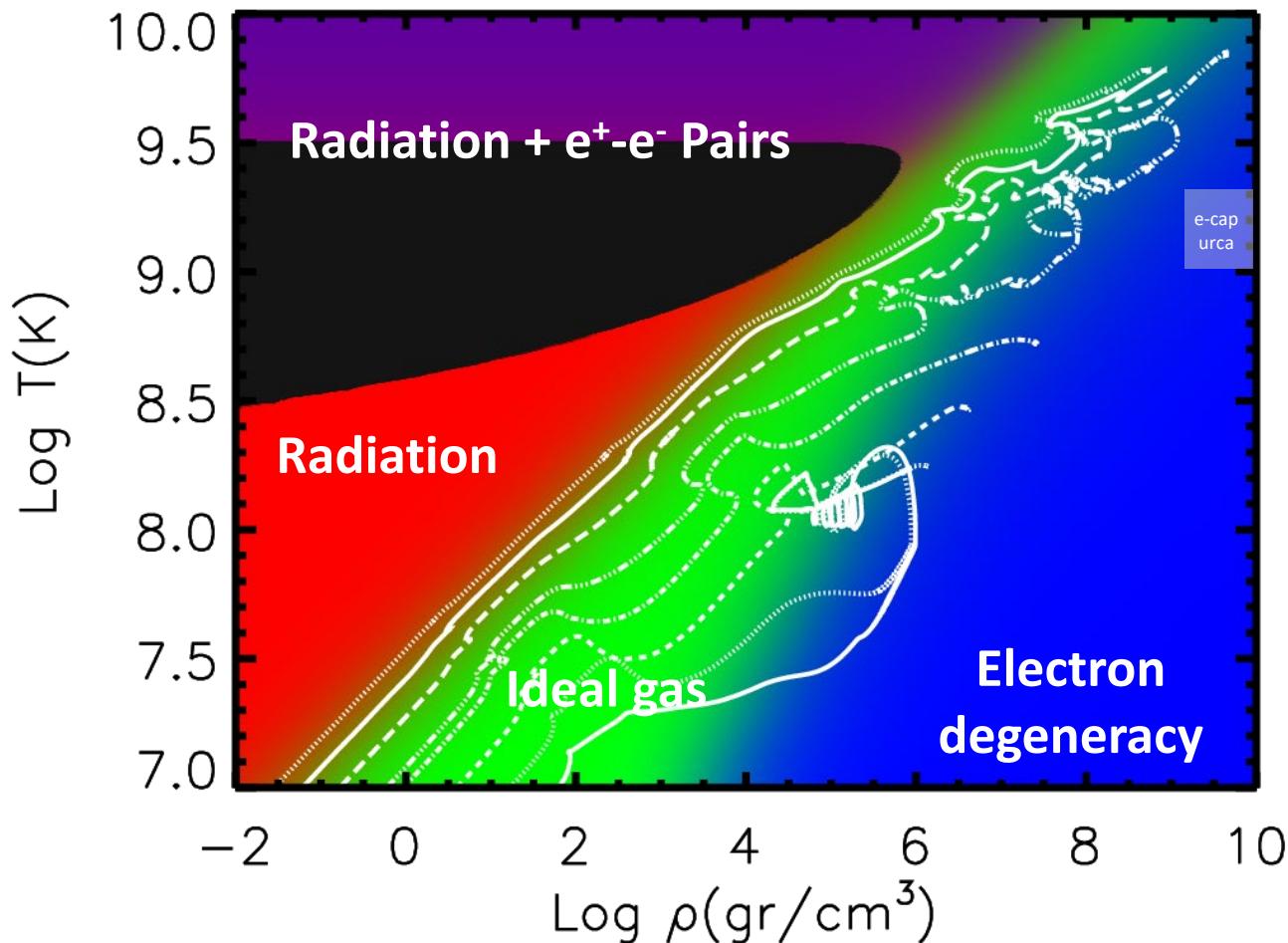
$$L \propto M^3$$

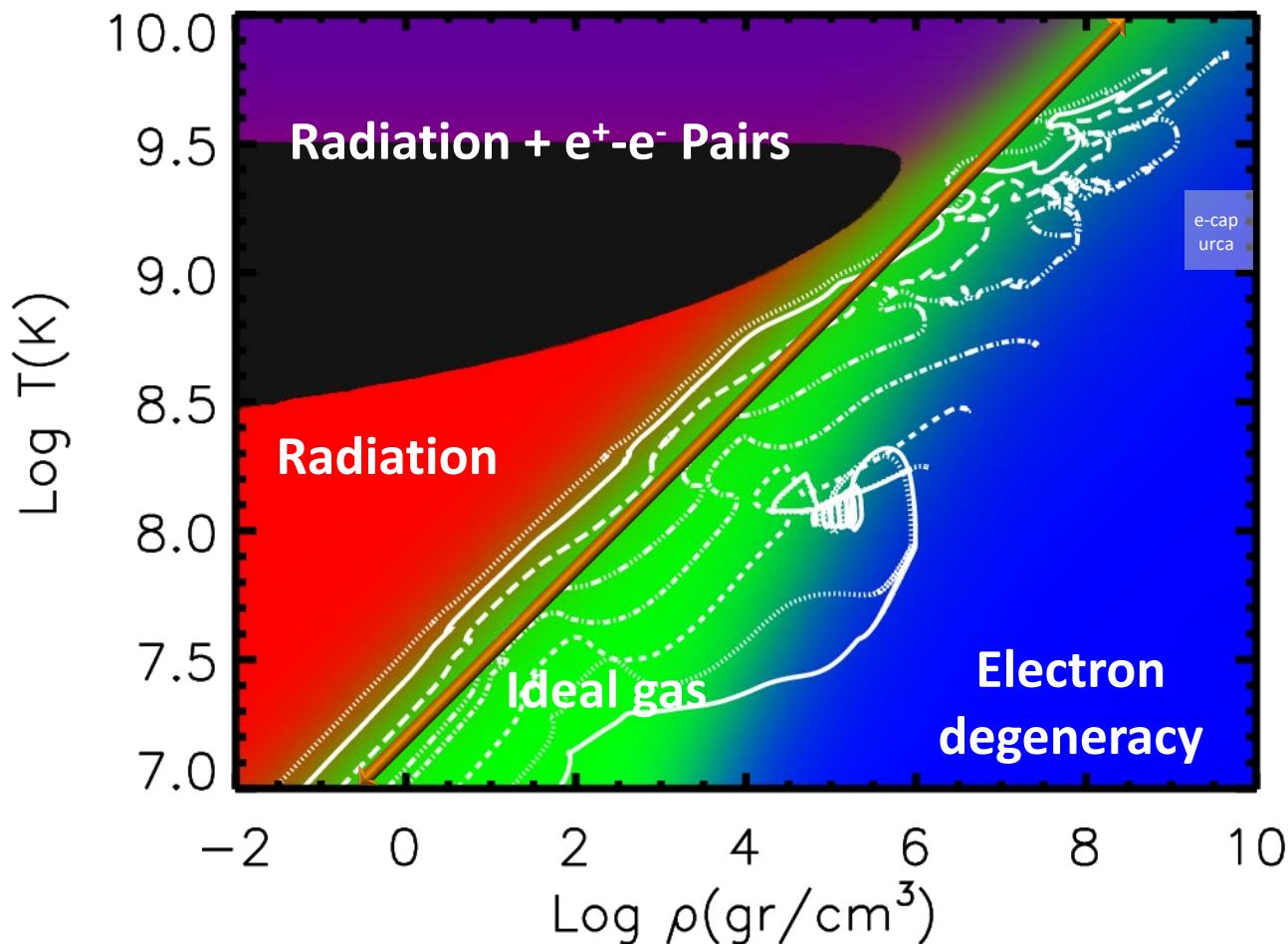
+ radiative equilibrium (valid in self regulating conditions)

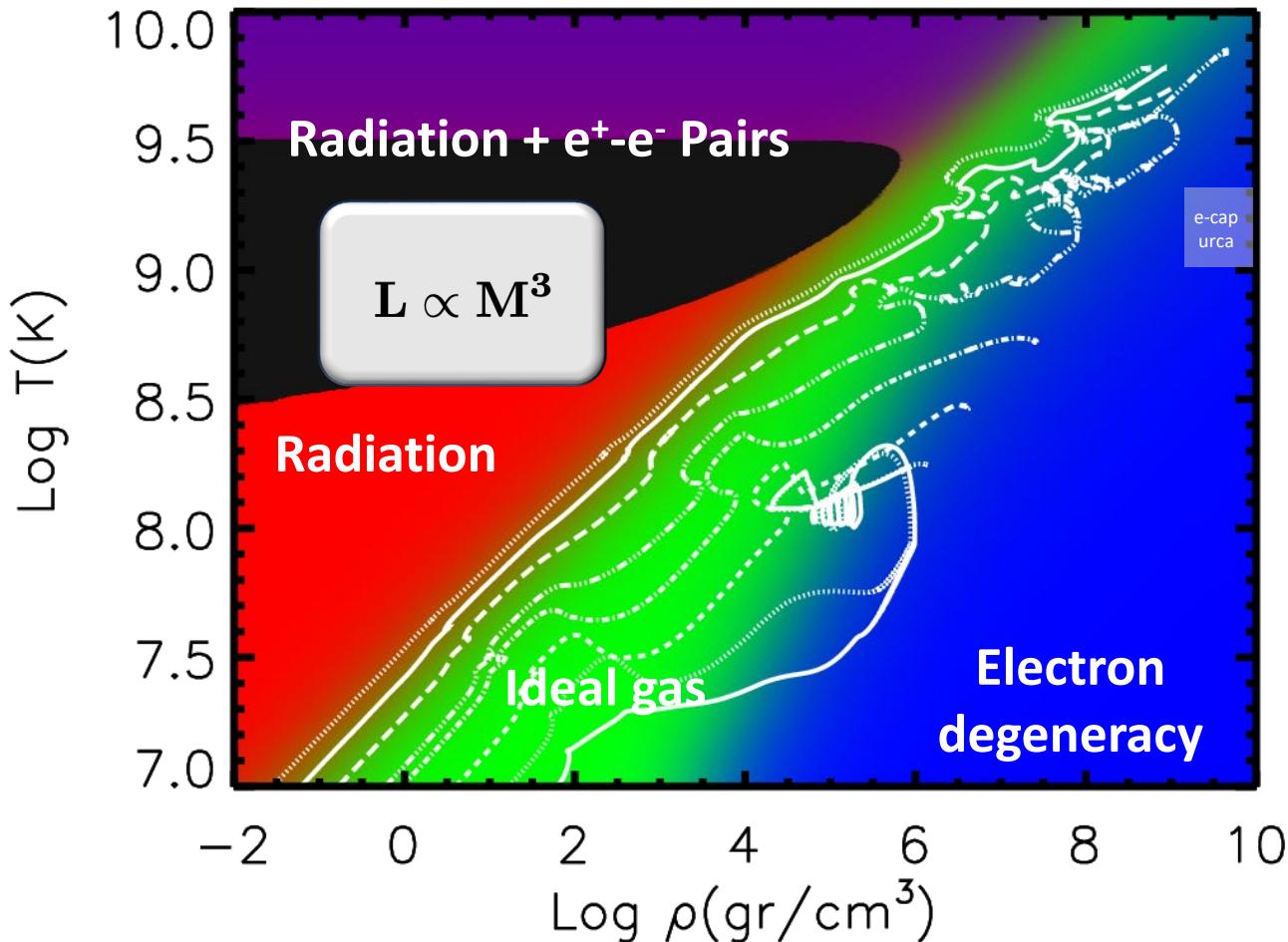


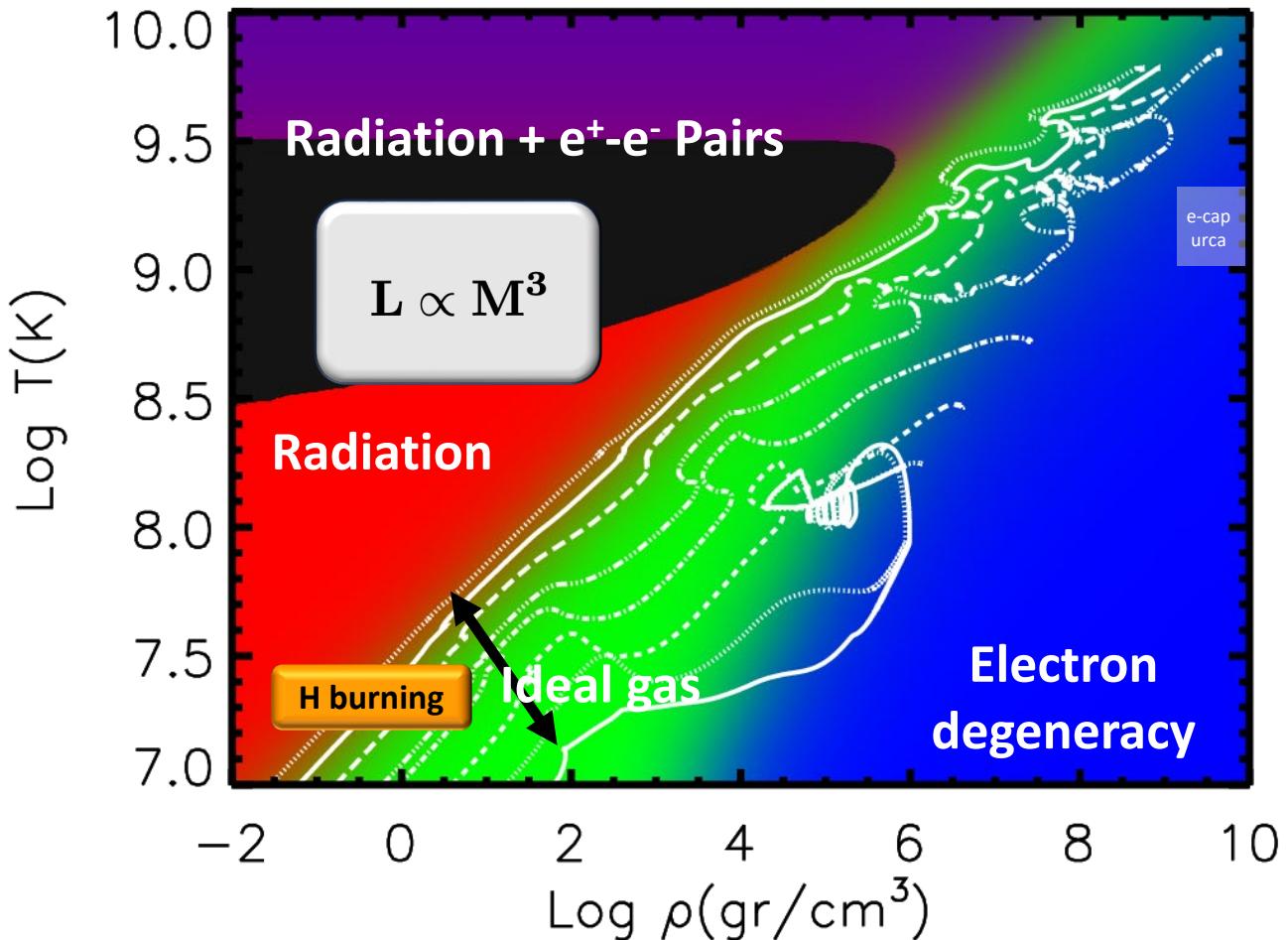


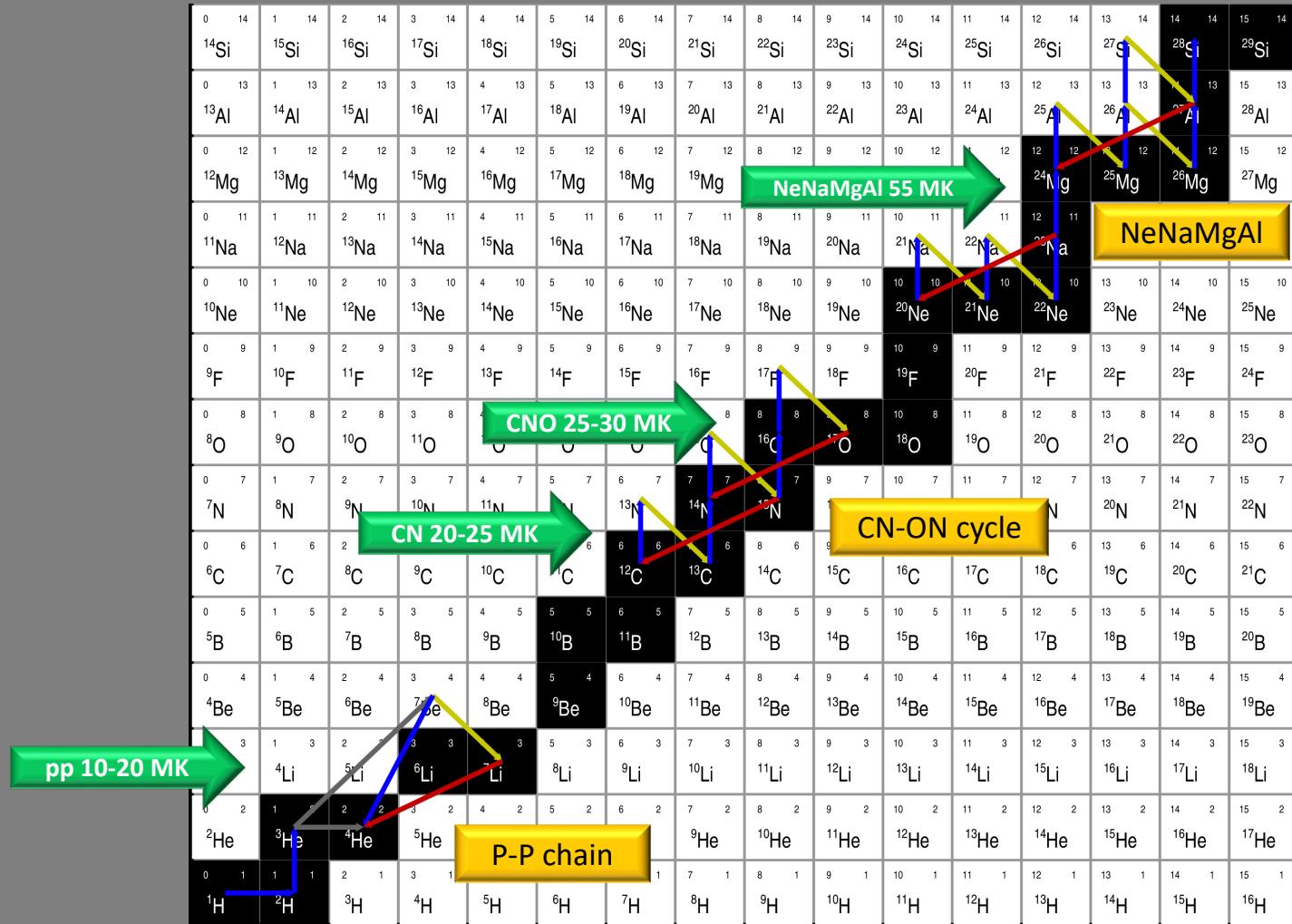


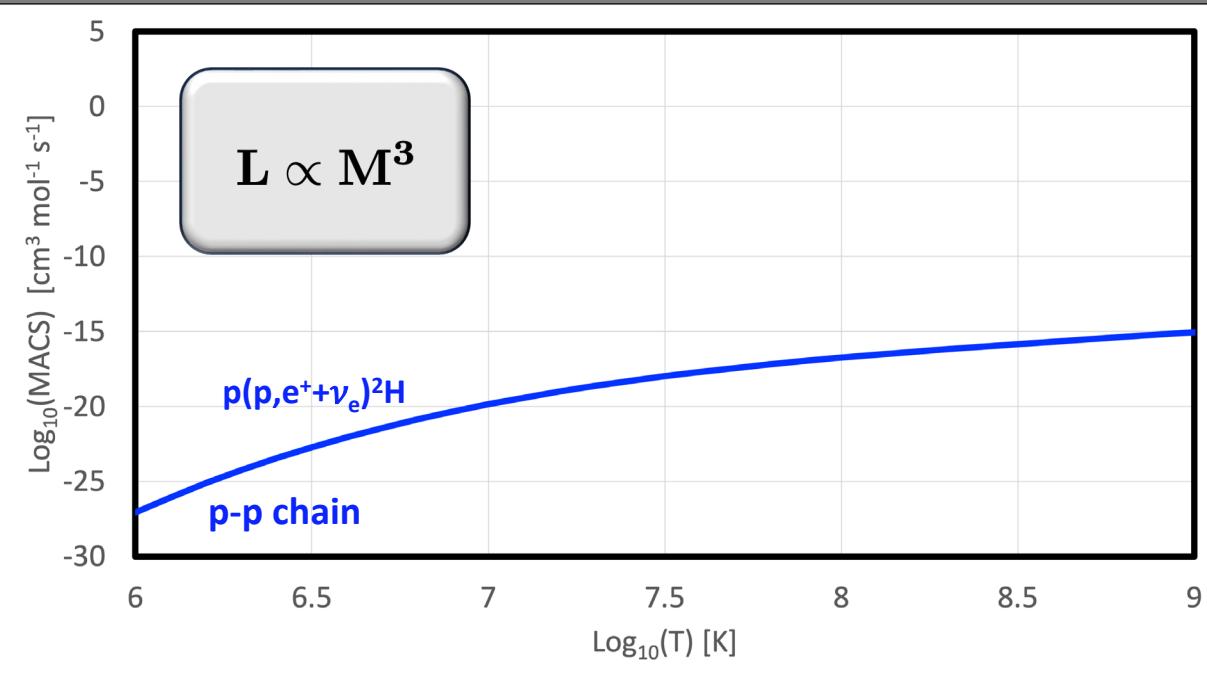




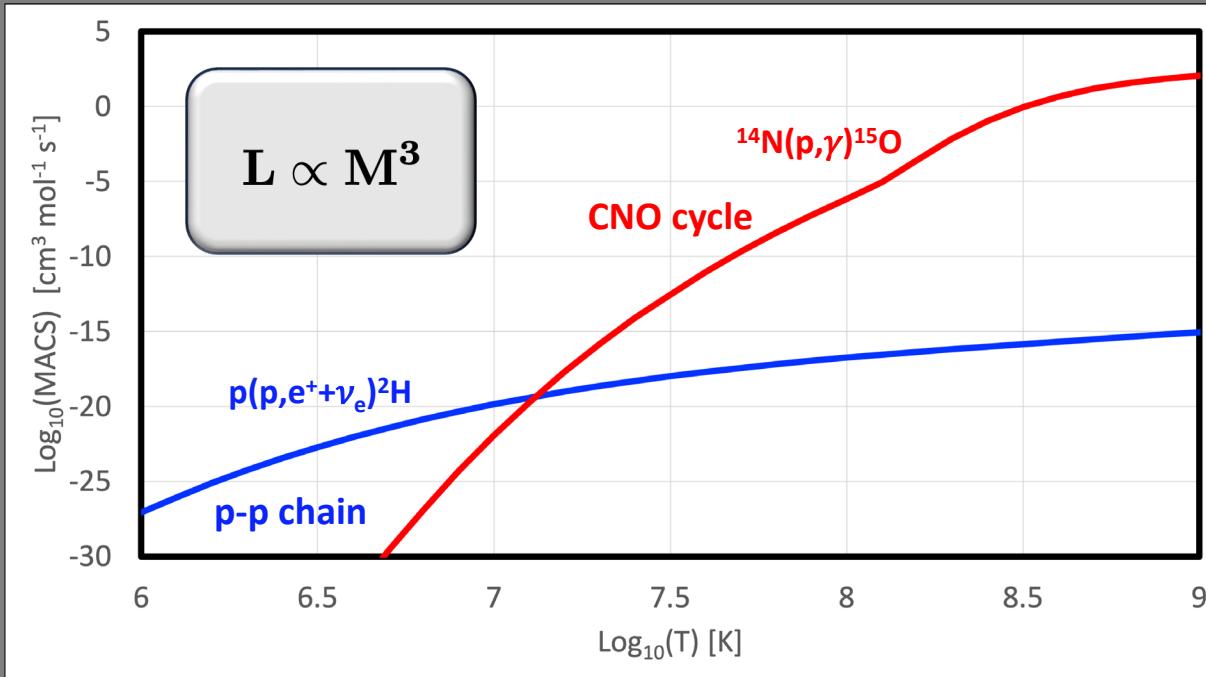




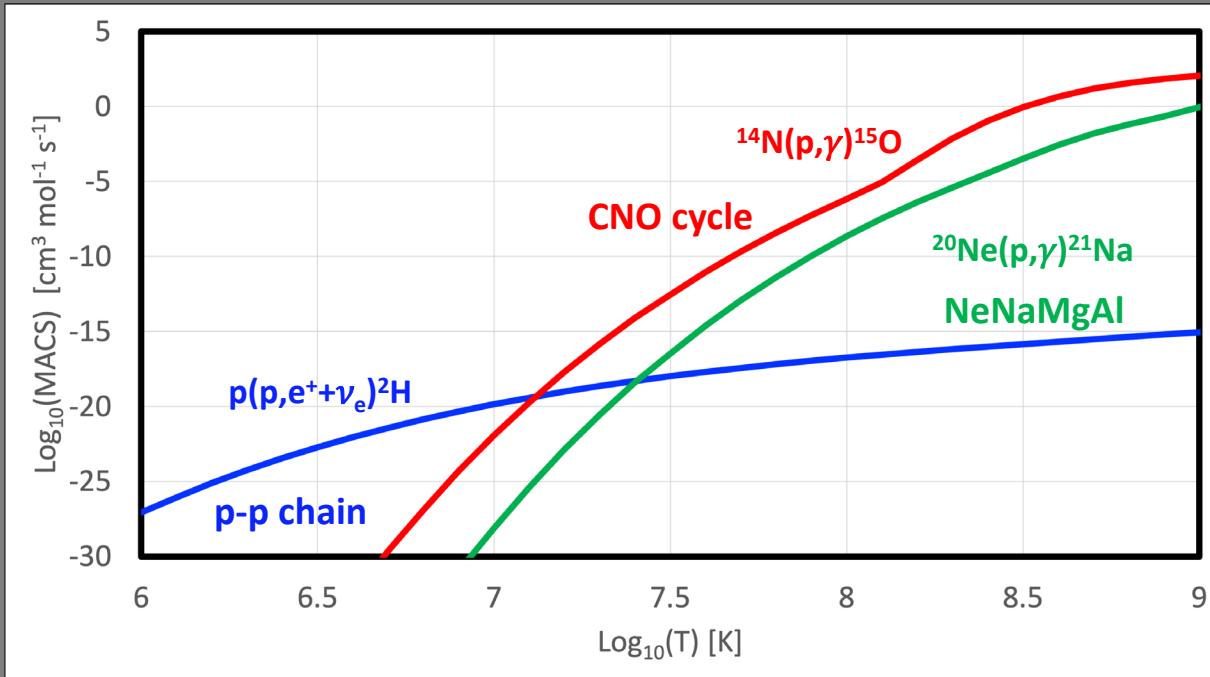




$$R_{ij} = \frac{y_i y_j}{1 + \delta_{ij}} N_A^2 \rho^2 <\sigma v>_{ij}$$



$$R_{ij} = \frac{y_i y_j}{1 + \delta_{ij}} N_A^2 \rho^2 <\sigma v>_{ij}$$

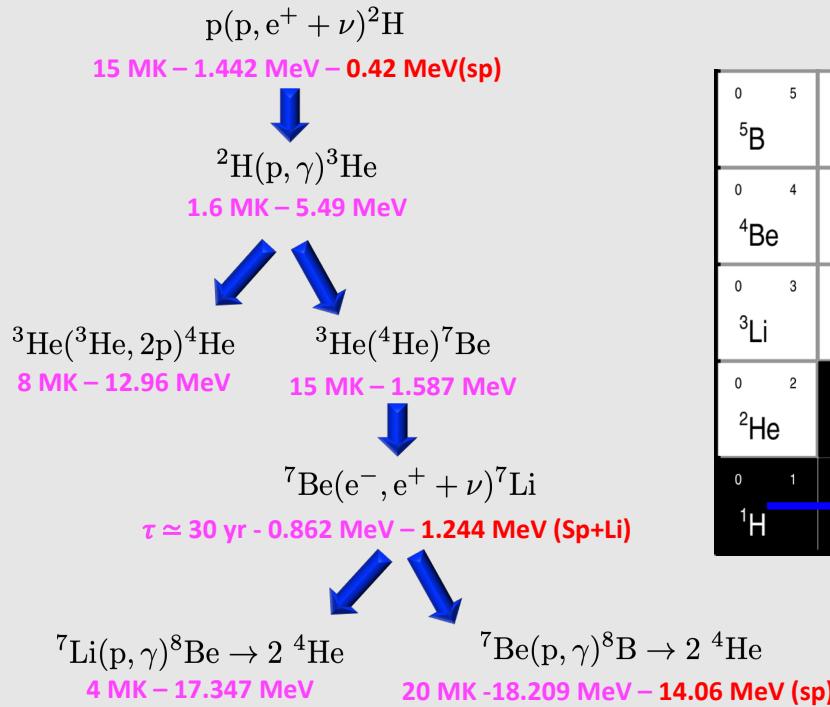


The P-P chain

$$4\text{p} \rightarrow {}^4\text{He} \quad [26.73 \text{ MeV}]$$

$$E_{\text{H} \rightarrow \text{He}} = 6.44 \times 10^{18} \text{ erg g}^{-1}$$

Energy released by an Earthquake of magnitudo 5 on the Richter scale

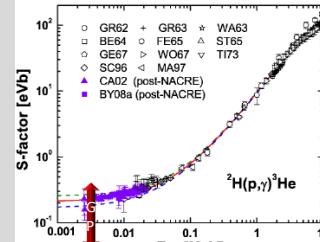
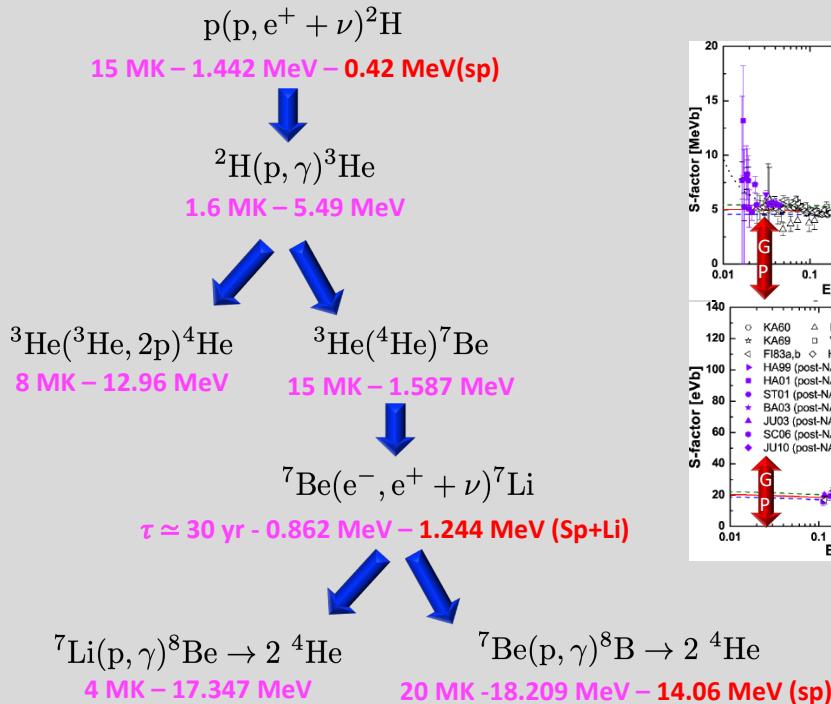


0	5	1	5	2	5	3	5	4	5	5	5
5B		6B		7B		8B		9B		10B	
0	4	1	4	2	4	3	4	4	4	5	4
4Be		5Be		6Be		7Be		8Be		9Be	
0	3	1	3	2	3	3	3	3	3	5	3
3Li		4Li		5Li		6Li		7Li		8Li	
0	2	1	2	2	2	3	2	4	2	5	2
2He		3He		4He		5He		6He		7He	
0	1	1	1	2	1	3	1	4	1	5	1
1H		2H		3H		4H		5H		6H	

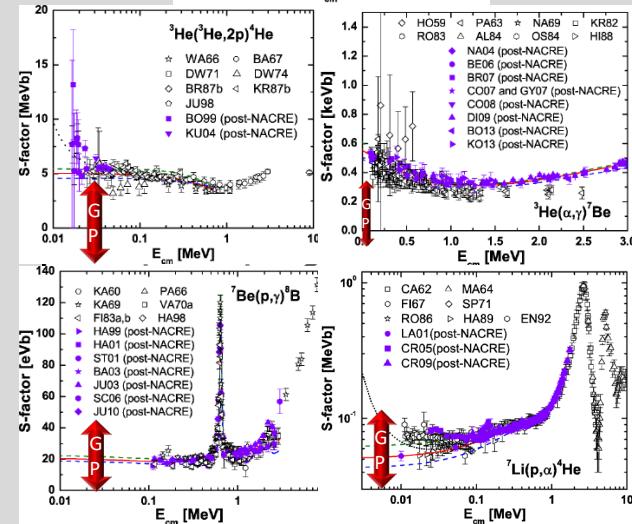
The P-P chain



$$E_{\text{H} \rightarrow \text{He}} = 6.44 \times 10^{18} \text{ erg g}^{-1}$$

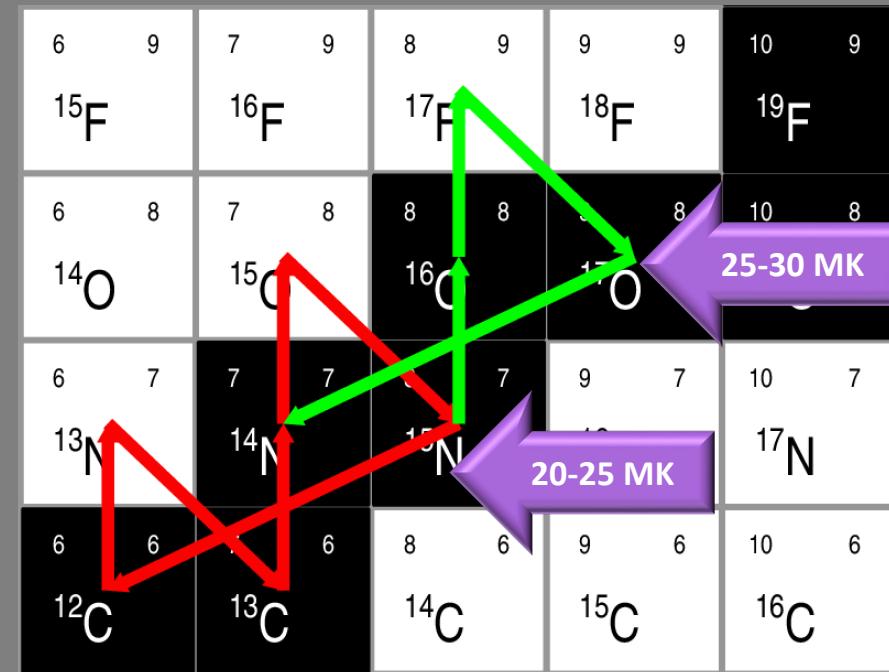
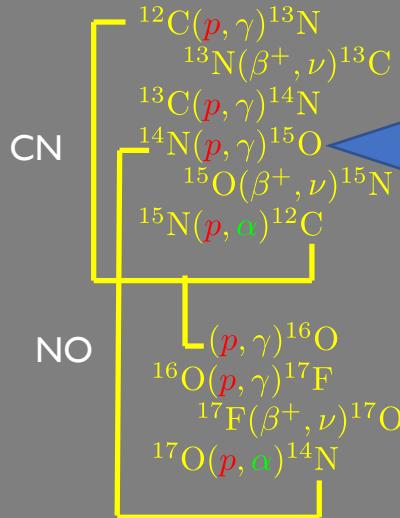


**Earthquake of
Richter scale**



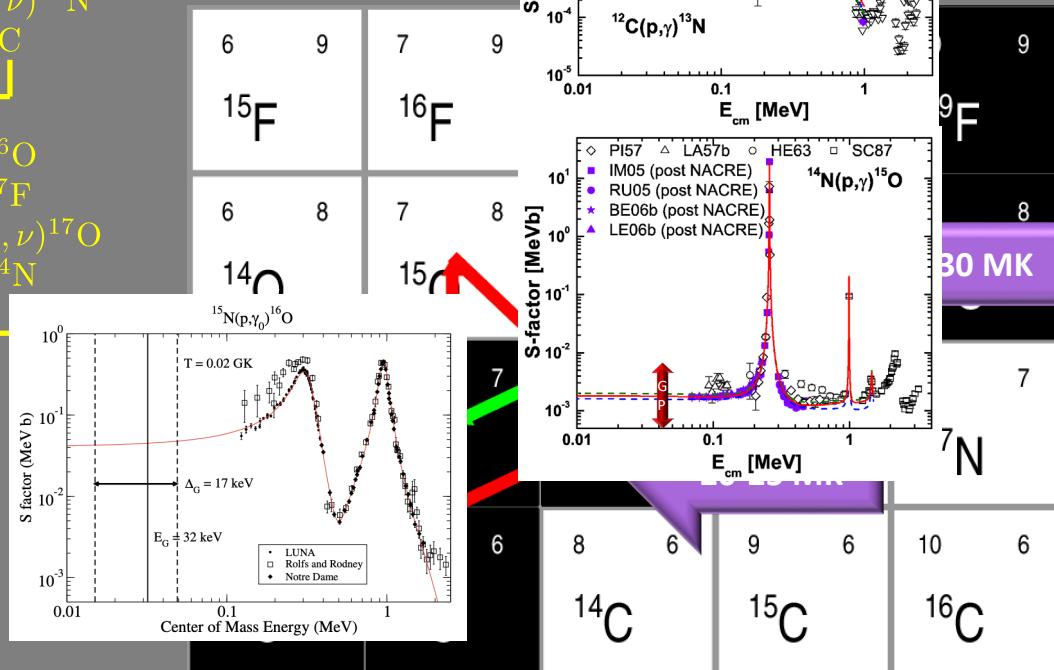
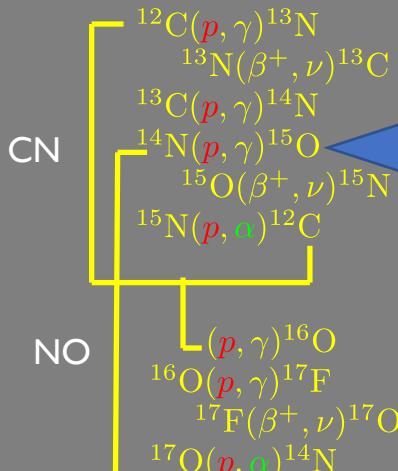
CNO cycle

$$T \sim 3 - 5 \cdot 10^7 \text{ K} \quad \rho \sim 1 - 10 \text{ gcm}^{-3}$$



CNO cycle

$$T \sim 3 - 5 \cdot 10^7 \text{ K} \quad \rho \sim 1 - 10 \text{ g cm}^{-3}$$



CNO cycle

$$T > 20 \text{ MK} \quad \rho \sim 1 - 10 \text{ [gr cm}^{-3}\text{]}$$



CNO processed material: $^{12}\text{C} \downarrow$ $^{14}\text{N} \uparrow$ $^{16}\text{O} \downarrow$

Typical equilibrium ratios:

$$\frac{^{12}\text{C}}{^{13}\text{C}} \sim 4 \quad \frac{^{14}\text{N}}{^{15}\text{N}} \sim 45000 \quad \frac{^{16}\text{O}}{^{17}\text{O}} \sim 170$$

$$\frac{N}{C} \sim 140 \quad \frac{C}{O} \sim 0.2$$

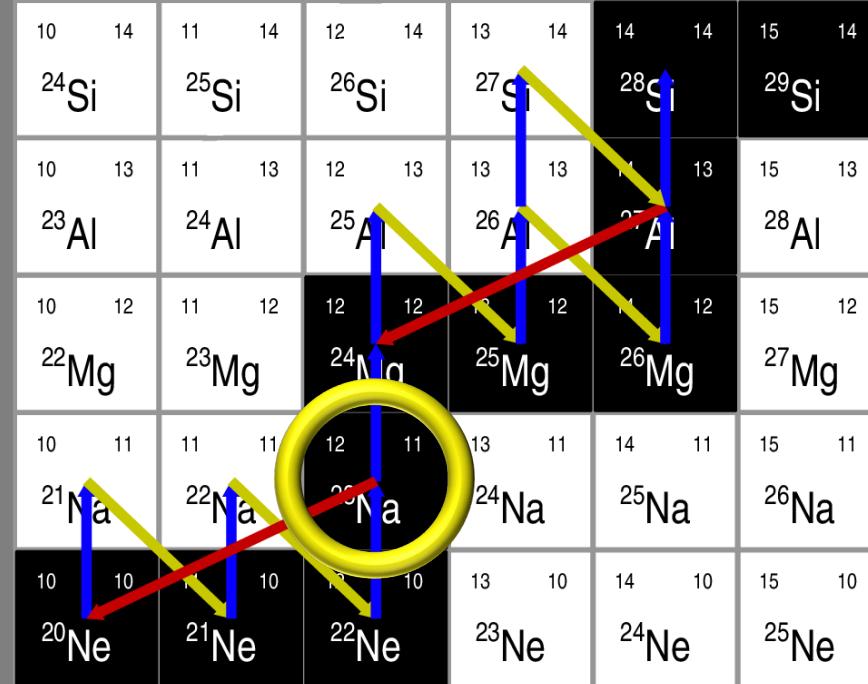
NeNa & MgAl (chains or cycles?)

T > 55 MK

$^{20}\text{Ne}(p, \gamma)^{21}\text{Na}$
 $^{21}\text{Na}(\beta^+, \nu)^{21}\text{Ne}$
 $^{21}\text{Ne}(p, \gamma)^{22}\text{Na}$
 $^{22}\text{Na}(\beta^+, \nu)^{22}\text{Ne}$
 $^{22}\text{Ne}(p, \gamma)^{23}\text{Na}$
 $^{23}\text{N}(p, \alpha)^{20}\text{Ne}$



$^{24}\text{Mg}(p, \gamma)^{25}\text{Al}$
 $^{25}\text{Al}(\beta^+, \nu)^{25}\text{Mg}$
 $^{25}\text{Mg}(p, \gamma)^{26}\text{Al}$
 $^{26}\text{Al}(\beta^+, \nu)^{26}\text{Mg}$
 $^{26}\text{Mg}(p, \gamma)^{27}\text{Al}$
 $^{27}\text{Al}(p, \alpha)^{24}\text{Mg}$



NeNa & MgAl (chains or cycles?)

NeNa

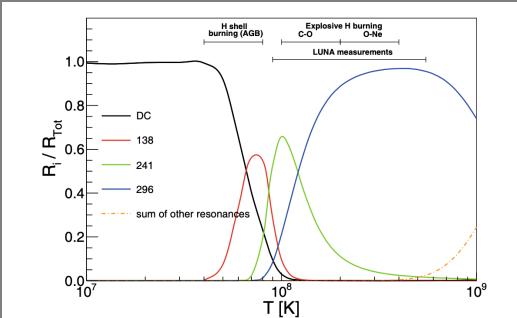
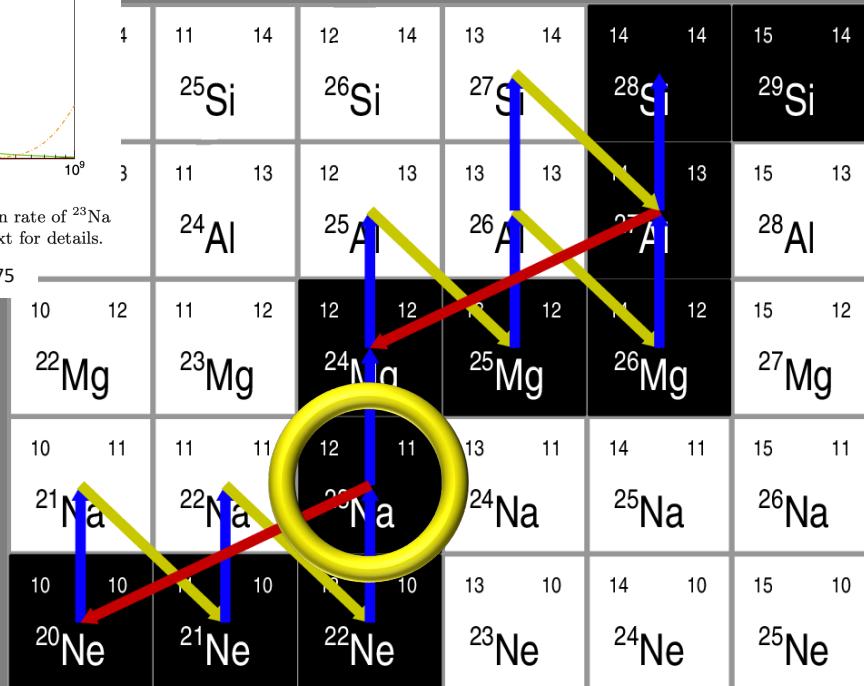


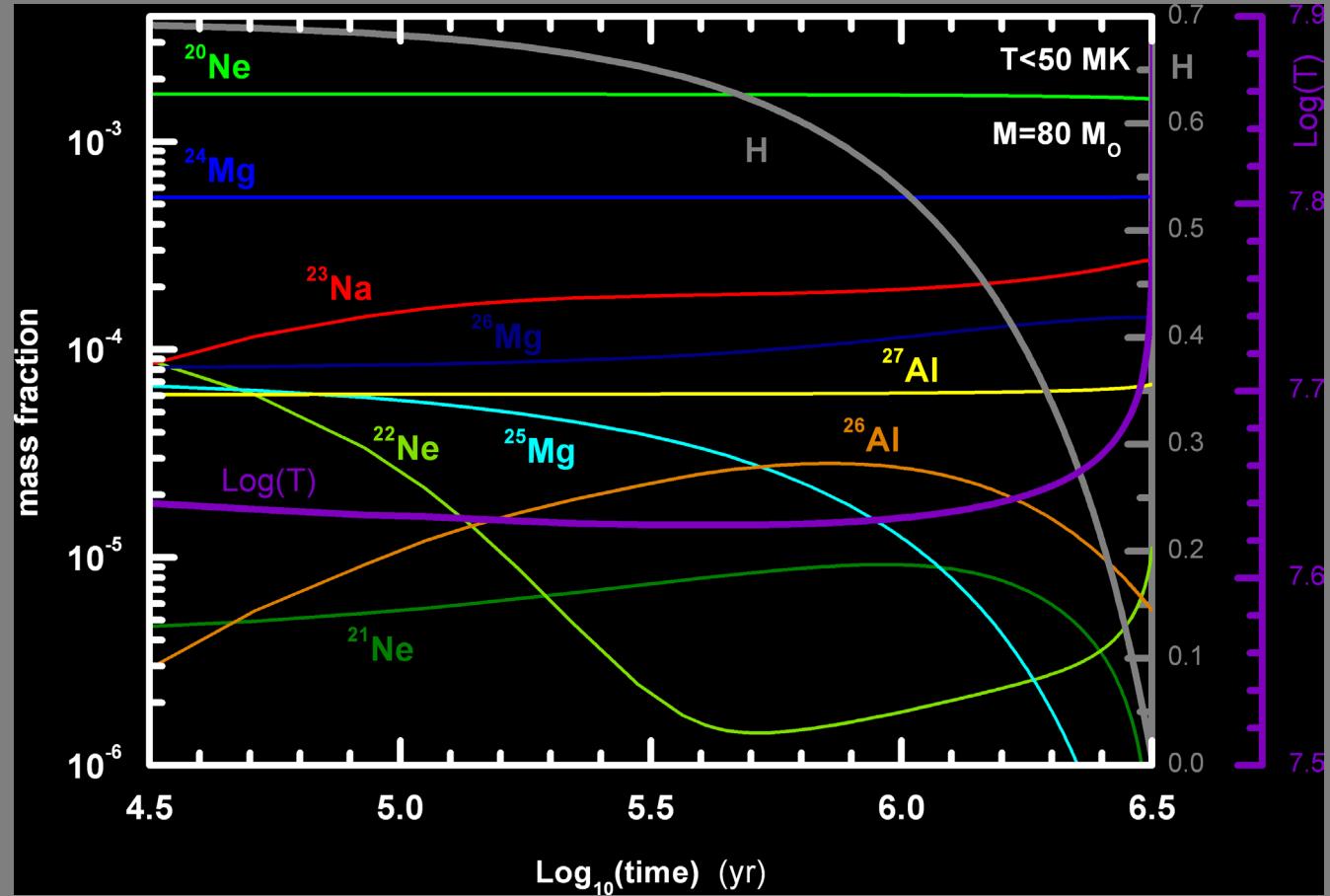
Fig. 12. Fractional contributions to the reaction rate of $^{23}\text{Na}(p, \gamma)^{24}\text{Mg}$, as a function of temperature, see text for details.

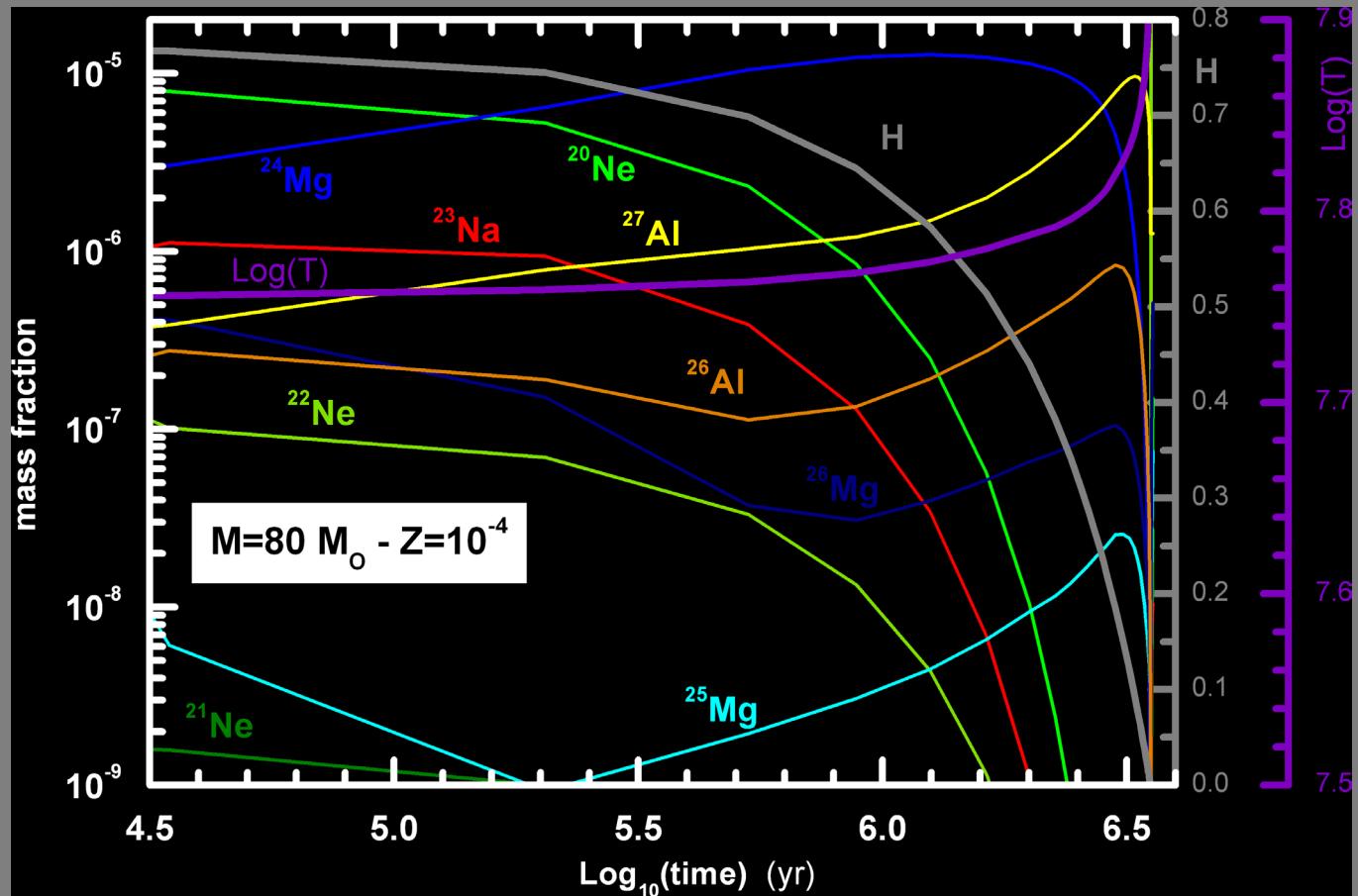
> 55 MK



MgAl

- $^{24}\text{Mg}(p, \gamma)^{25}\text{Al}$
- $^{25}\text{Al}(\beta^+, \nu)^{25}\text{Mg}$
- $^{25}\text{Mg}(p, \gamma)^{26}\text{Al}$
- $^{26}\text{Al}(\beta^+, \nu)^{26}\text{Mg}$
- $^{26}\text{Mg}(p, \gamma)^{27}\text{Al}$
- $^{27}\text{Al}(p, \alpha)^{24}\text{Mg}$





NeNa & MgAl (chains or cycles?)

$T > 55 \text{ MK}$

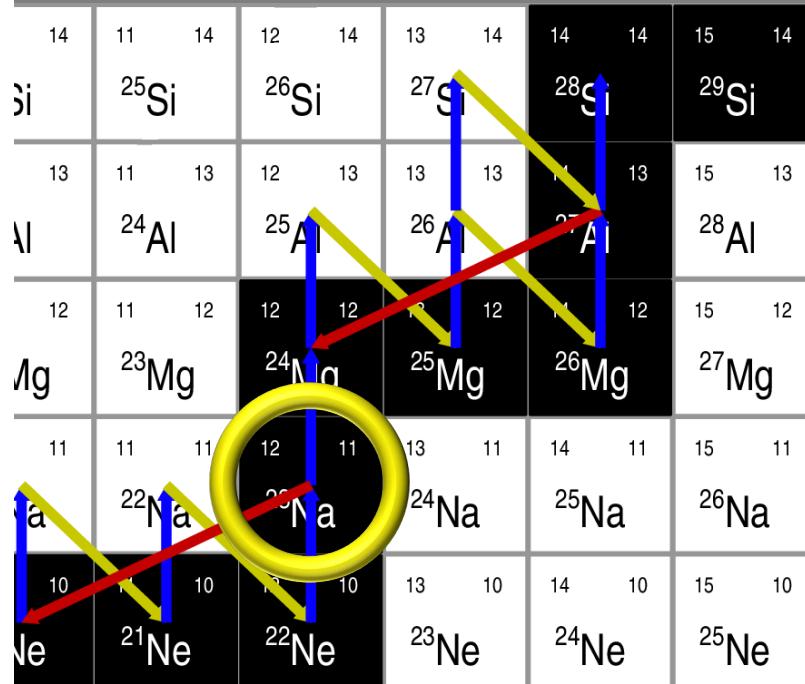
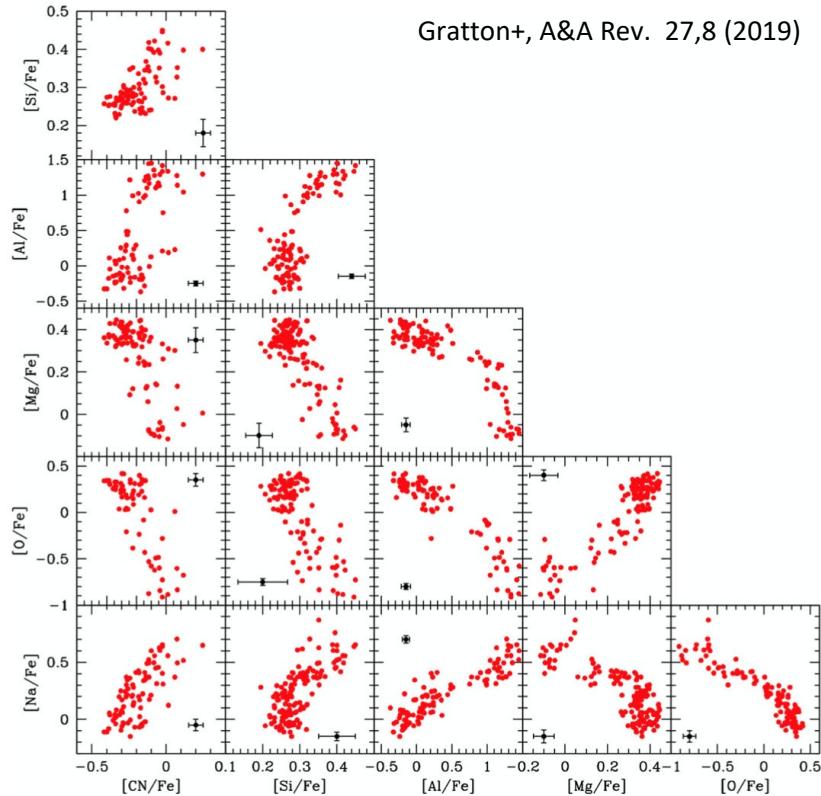
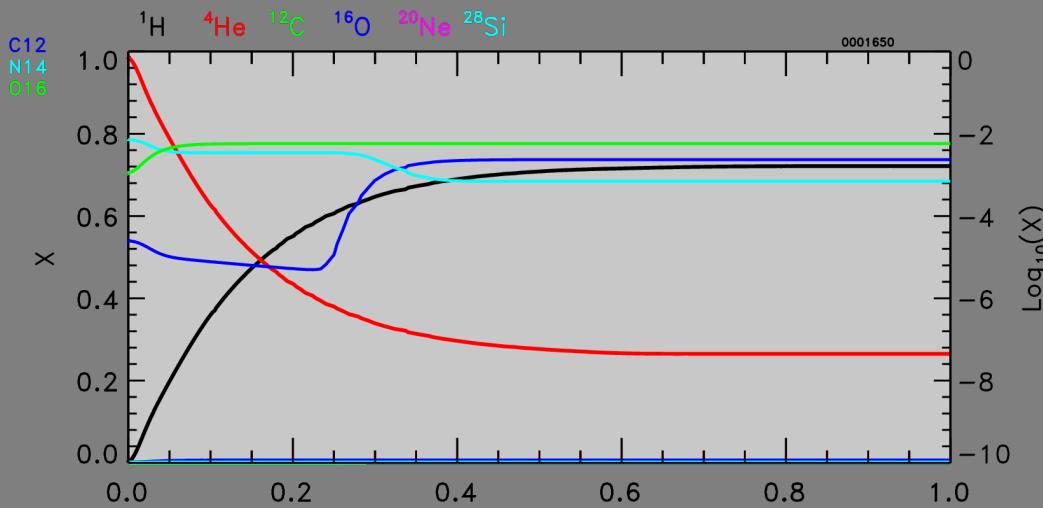


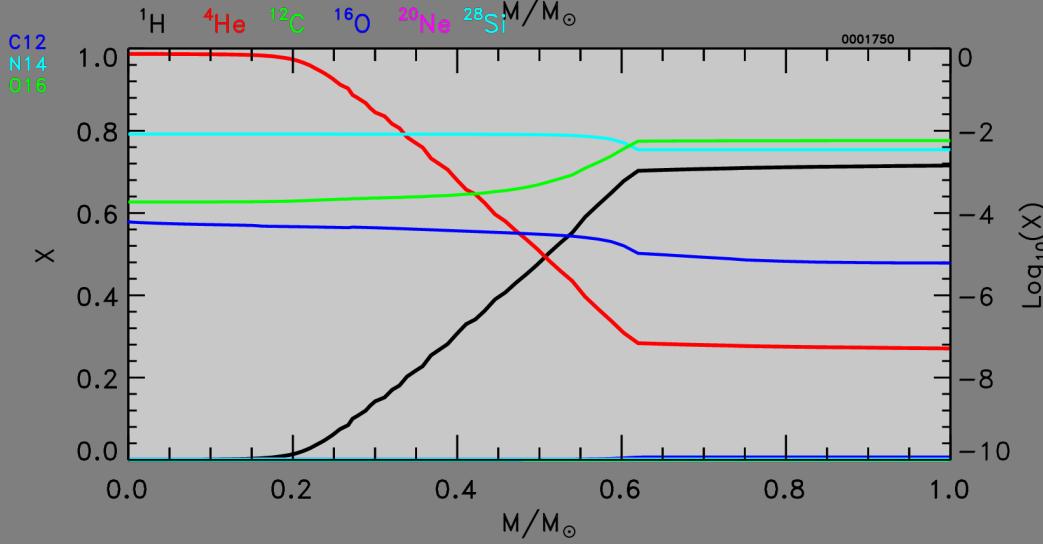
Fig. 2 Run of abundance ratios for light elements in RGB stars in NGC 2808. O, Na, Si, and Mg are from Carretta (2015), Al and CN abundances from Carretta et al. (2018). The figure is adapted from the invited review by Carretta (2016)

0	14	1	14	2	14	3	14	4	14	5	14	6	14	7	14	8	14	9	14	10	14	11	14	12	14	13	14	14	14					
¹⁴ Si		¹⁵ Si		¹⁶ Si		¹⁷ Si		¹⁸ Si		¹⁹ Si		²⁰ Si		²¹ Si		²² Si		²³ Si		²⁴ Si		²⁵ Si		²⁶ Si		²⁷ Si		²⁸ Si		²⁹ Si				
0	13	1	13	2	13	3	13	4	13	5	13	6	13	7	13	8	13	9	13	10	13	11	13	12	13	13	13	13	13					
¹³ Al		¹⁴ Al		¹⁵ Al		¹⁶ Al		¹⁷ Al		¹⁸ Al		¹⁹ Al		²⁰ Al		²¹ Al		²² Al		²³ Al		²⁴ Al		²⁵ Al		²⁶ Al		²⁷ Al		²⁸ Al				
0	12	1	12	2	12	3	12	4	12	5	12	6	12	7	12	8	12	9	12	10	12	11	12	12	12	12	12	12	12					
¹² Mg		¹³ Mg		¹⁴ Mg		¹⁵ Mg		¹⁶ Mg		¹⁷ Mg		¹⁸ Mg		¹⁹ Mg										²⁴ Mg	²⁵ Mg	²⁶ Mg	²⁷ Mg							
0	11	1	11	2	11	3	11	4	11	5	11	6	11	7	11	8	11	9	11	10	11	11	12	11	13	11	14	11	15	12				
¹¹ Na		¹² Na		¹³ Na		¹⁴ Na		¹⁵ Na		¹⁶ Na		¹⁷ Na		¹⁸ Na		¹⁹ Na		²⁰ Na		²¹ Na		²² Na		²³ Na		²⁴ Na		²⁵ Na		²⁶ Na				
0	10	1	10	2	10	3	10	4	10	5	10	6	10	7	10	8	10	9	10	10	10	10	11	12	11	13	10	14	10	15	10			
¹⁰ Ne		¹¹ Ne		¹² Ne		¹³ Ne		¹⁴ Ne		¹⁵ Ne		¹⁶ Ne		¹⁷ Ne		¹⁸ Ne		¹⁹ Ne		²⁰ Ne		²¹ Ne		²² Ne		²³ Ne		²⁴ Ne		²⁵ Ne				
0	9	1	9	2	9	3	9	4	9	5	9	6	9	7	9	8	9	9	9	10	9	11	9	12	9	13	9	14	9	15	9			
⁹ F		¹⁰ F		¹¹ F		¹² F		¹³ F		¹⁴ F		¹⁵ F		¹⁶ F		¹⁷ F		¹⁸ F		¹⁹ F		²⁰ F		²¹ F		²² F		²³ F		²⁴ F				
0	8	1	8	2	8	3	8	4																										
⁸ O		⁹ O		¹⁰ O		¹¹ O		¹² O		¹³ O		¹⁴ O		¹⁵ O		¹⁶ O		¹⁷ O		¹⁸ O		¹⁹ O		²⁰ O		²¹ O		²² O		²³ O				
0	7	1	7	2	7	3	7	4	7	5	7	6	7	7	7	8	8	8	8	10	8	11	8	12	8	13	8	14	8	15	8			
⁷ N		⁸ N		⁹ N		¹⁰ N		¹¹ N								¹³ N	¹⁴ N	¹⁵ N	¹⁶ N	¹⁷ N	¹⁸ N	¹⁹ N	²⁰ N	²¹ N	²² N									
0	6	1	6	2													¹² C	¹³ C	¹⁴ C	¹⁵ C	¹⁶ C	¹⁷ C	¹⁸ C	¹⁹ C	²⁰ C	²¹ C								
⁶ C		⁷ C		⁸ C		⁹ C		¹⁰ C		¹¹ C		¹² C		¹³ C		¹⁴ C		¹⁵ C		¹⁶ C		¹⁷ C		¹⁸ C		¹⁹ C		²⁰ C		²¹ C				
0	5	1	5	2	5	3	5	4	5	5	5	6	5	7	5	8	5	9	5	10	5	11	5	12	5	13	5	14	5	15	5			
⁵ B		⁶ B		⁷ B		⁸ B		⁹ B		¹⁰ B		¹¹ B		¹² B		¹³ B		¹⁴ B		¹⁵ B		¹⁶ B		¹⁷ B		¹⁸ B		¹⁹ B		²⁰ B		²¹ B		
0	4	1	4	2	4	3	4	4	4	5	4	6	4	7	4	8	4	9	4	10	4	11	4	12	4	13	4	14	4	15	4			
⁴ Be		⁵ Be		⁶ Be		⁷ Be		⁸ Be		⁹ Be		¹⁰ Be		¹¹ Be		¹² Be		¹³ Be		¹⁴ Be		¹⁵ Be		¹⁶ Be		¹⁷ Be		¹⁸ Be		¹⁹ Be				
0	3	1	3	2	3	2	3	3	3	3	5	3	6	3	7	3	8	3	9	3	10	3	11	3	12	3	13	3	14	3	15	3		
² He		³ He		⁴ He		⁵ He		⁶ He		⁷ He		⁸ He		⁹ He		¹⁰ He		¹¹ He		¹² He		¹³ He		¹⁴ He		¹⁵ He		¹⁶ He		¹⁷ He		¹⁸ He		¹⁹ He
0	1	1	1	1	1	2	1	3	1	4	1	5	1	6	1	7	1	8	1	9	1	10	1	11	1	12	1	13	1	14	1	15	1	
¹ H		² H		³ H		⁴ H		⁵ H		⁶ H		⁷ H		⁸ H		⁹ H		¹⁰ H		¹¹ H		¹² H		¹³ H		¹⁴ H		¹⁵ H		¹⁶ H				

$M < 1.3 M_{\odot}$
Radiative core
PP chain



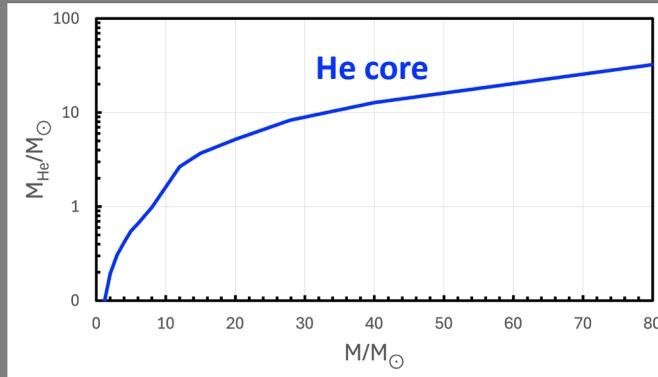
$M > 1.3 M_{\odot}$
convective core
CNO cycle



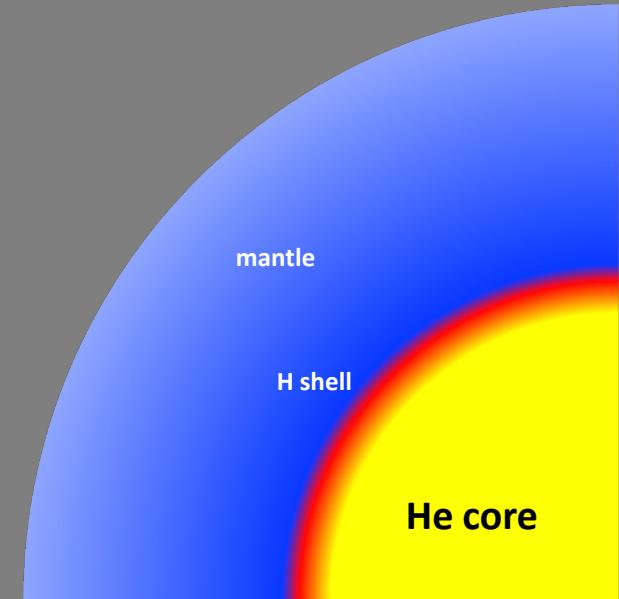
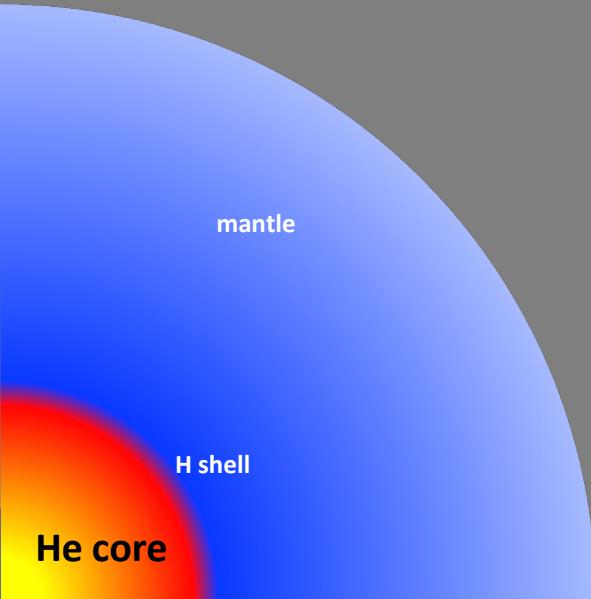
Central H exhaustion

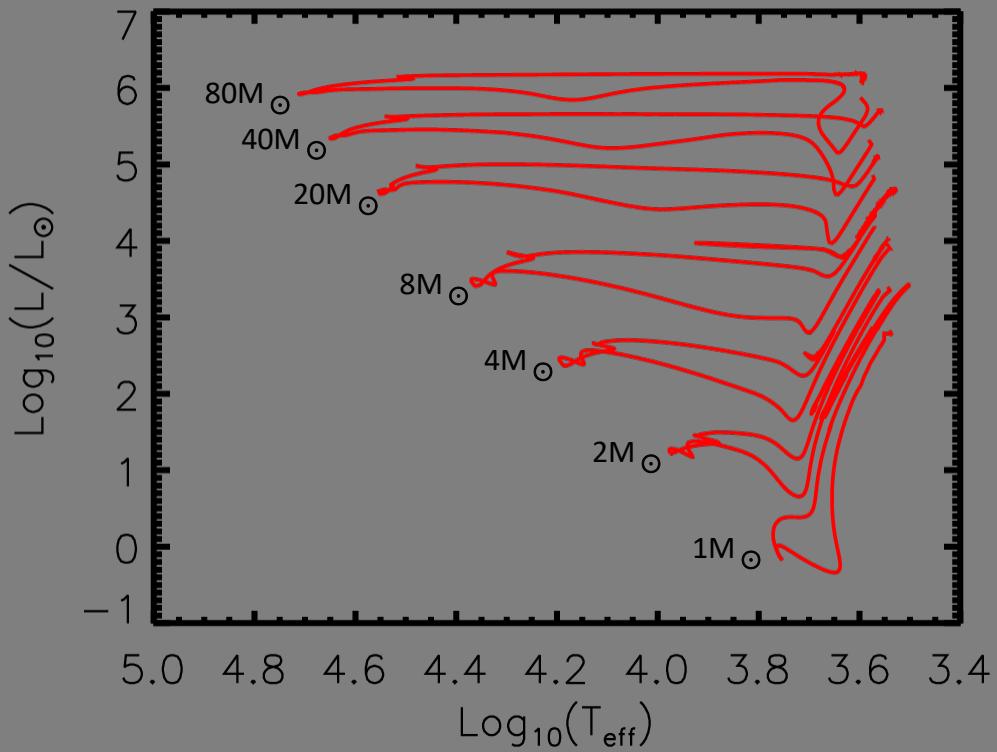
monotonic family of stars

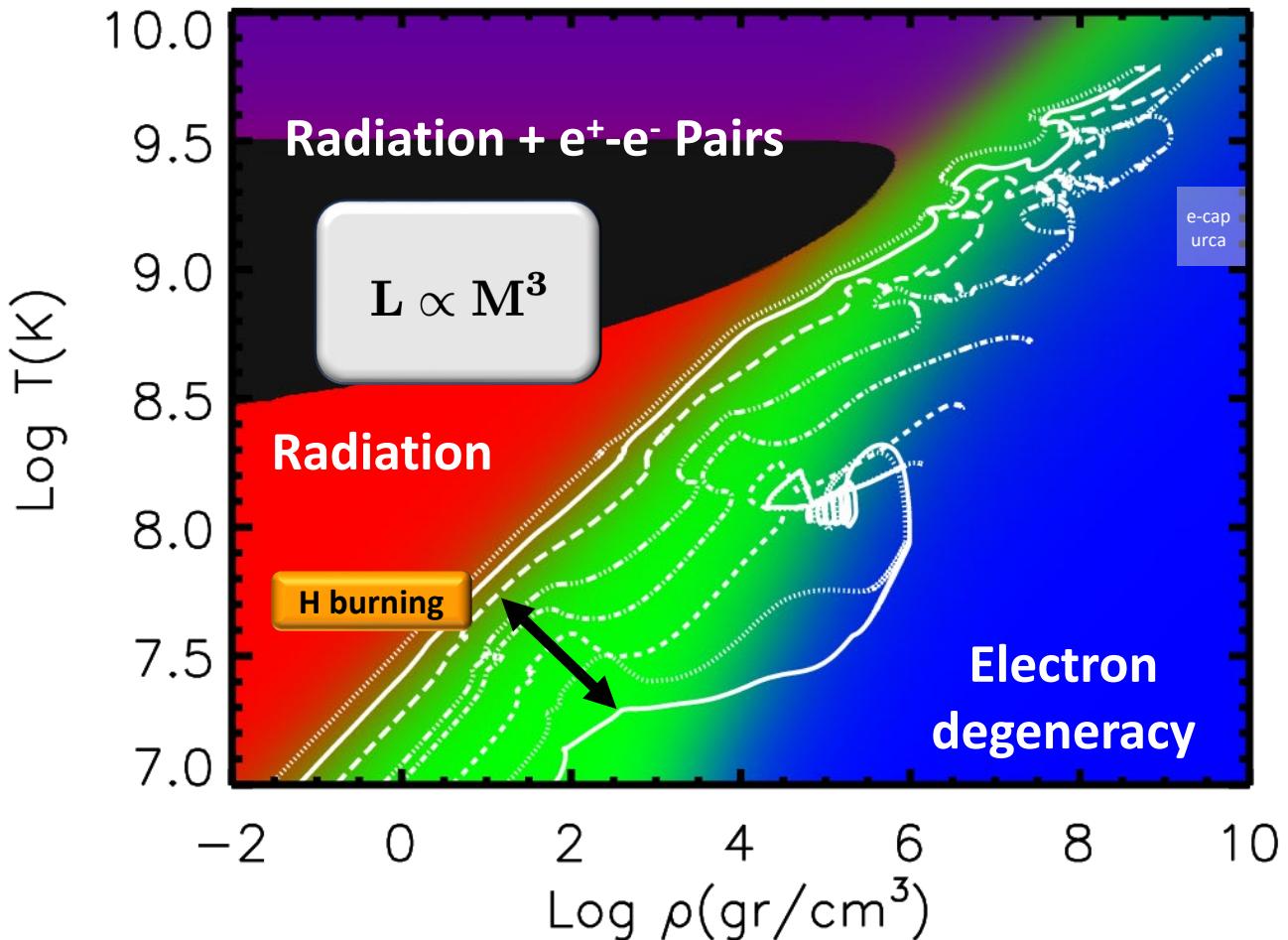
$M < 1.3 M_{\odot}$
Radiative core
PP chain

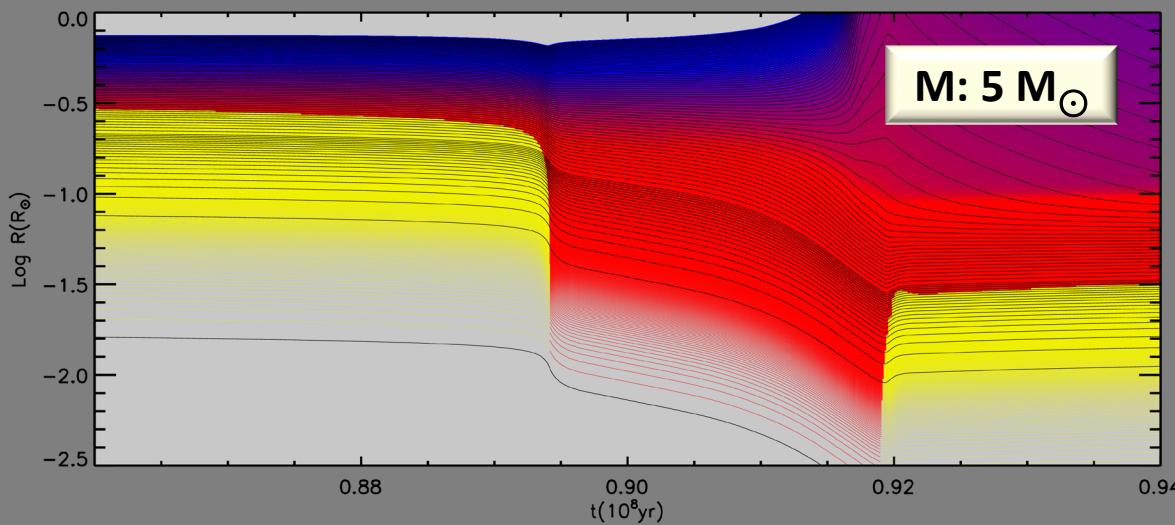
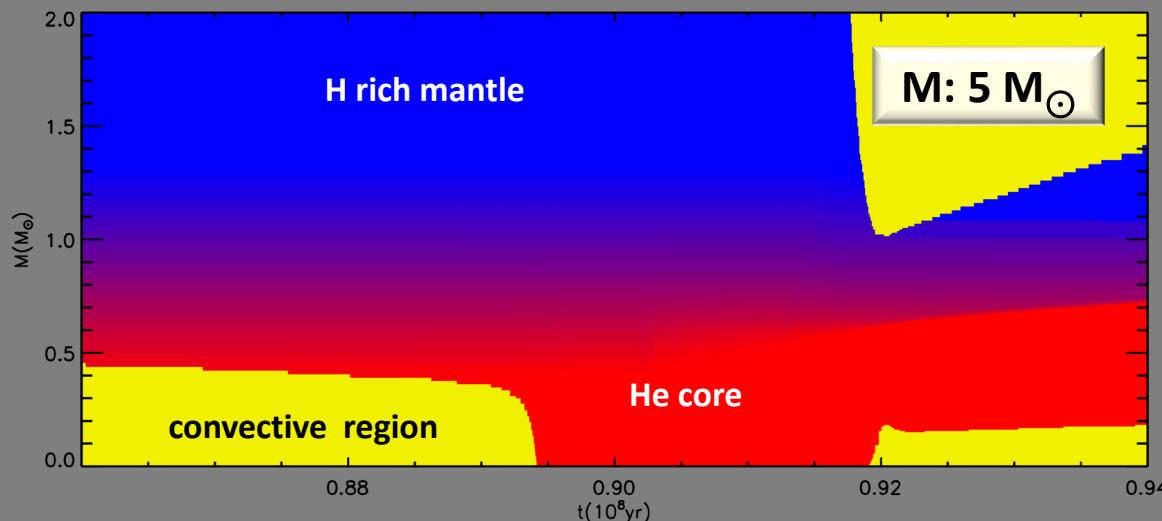


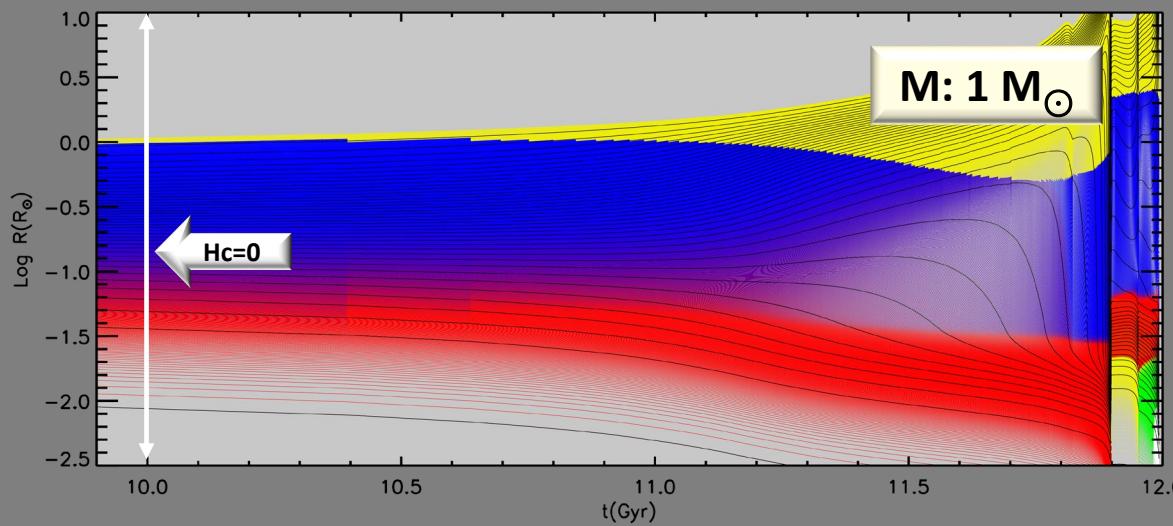
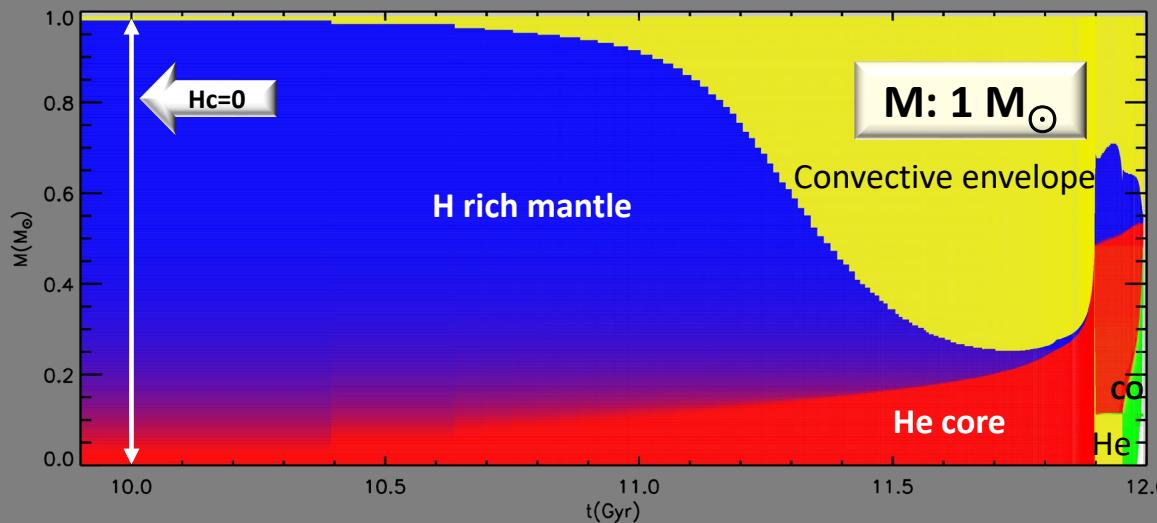
$M > 1.3 M_{\odot}$
convective core
CNO cycle

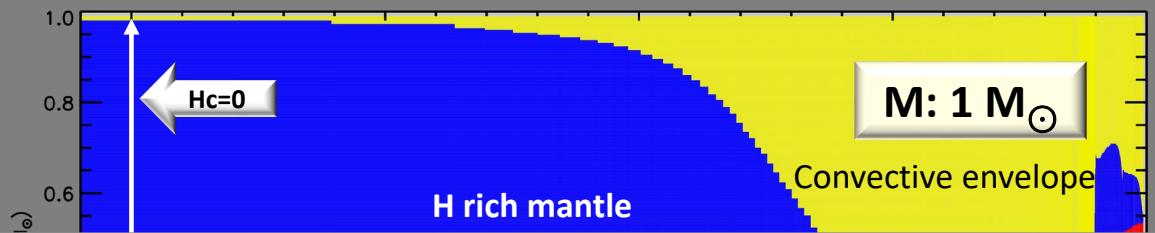




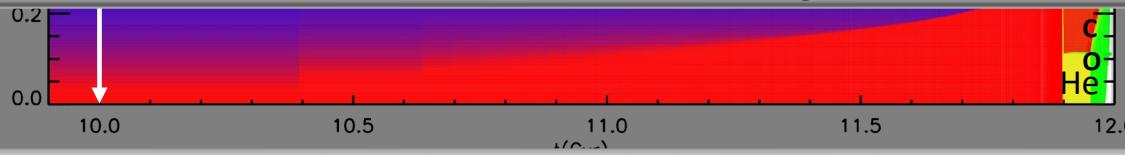








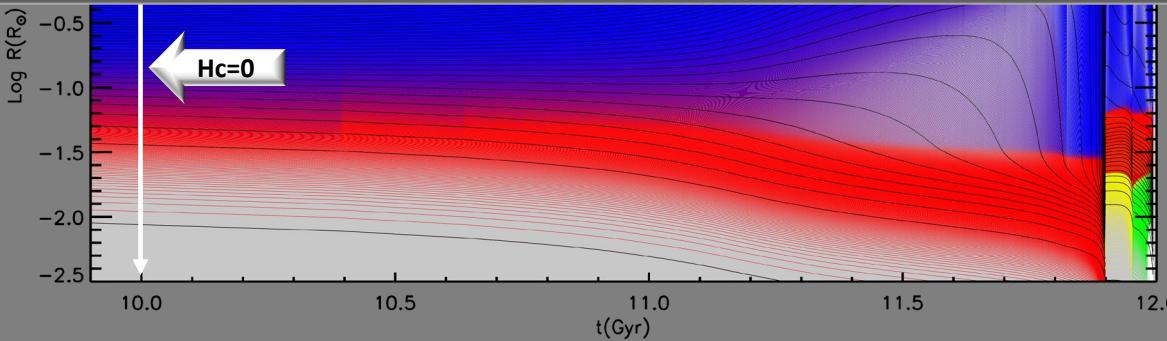
Stars less massive than, roughly, $2.3M_{\odot}$ ignite He off-center when the He core reaches a value $\sim 0.5M_{\odot}$

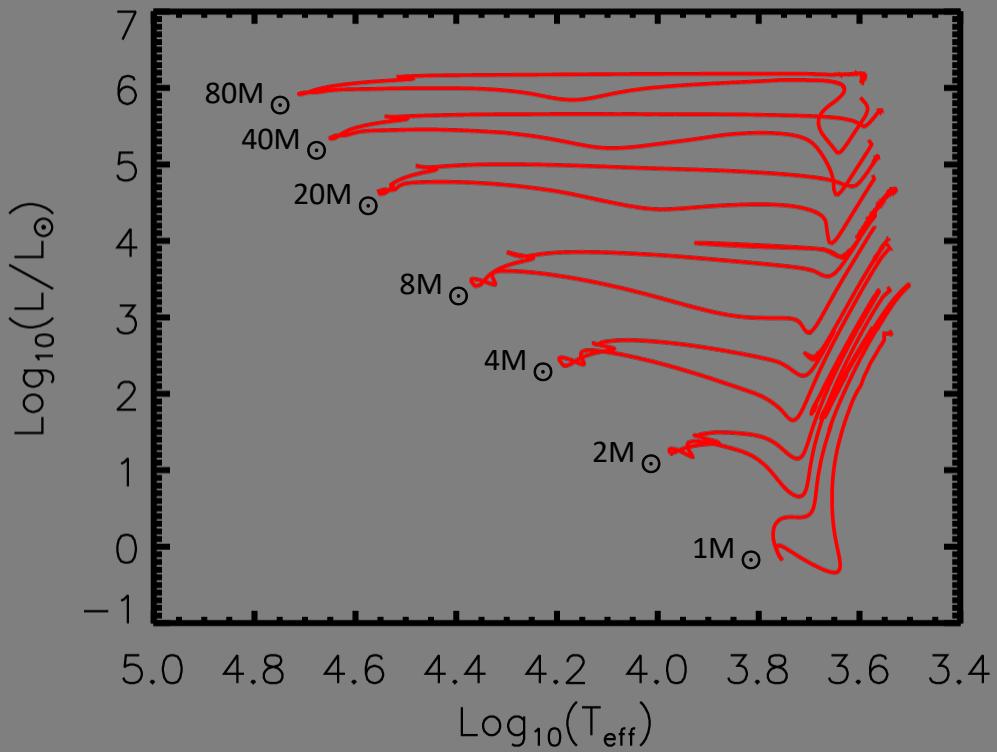


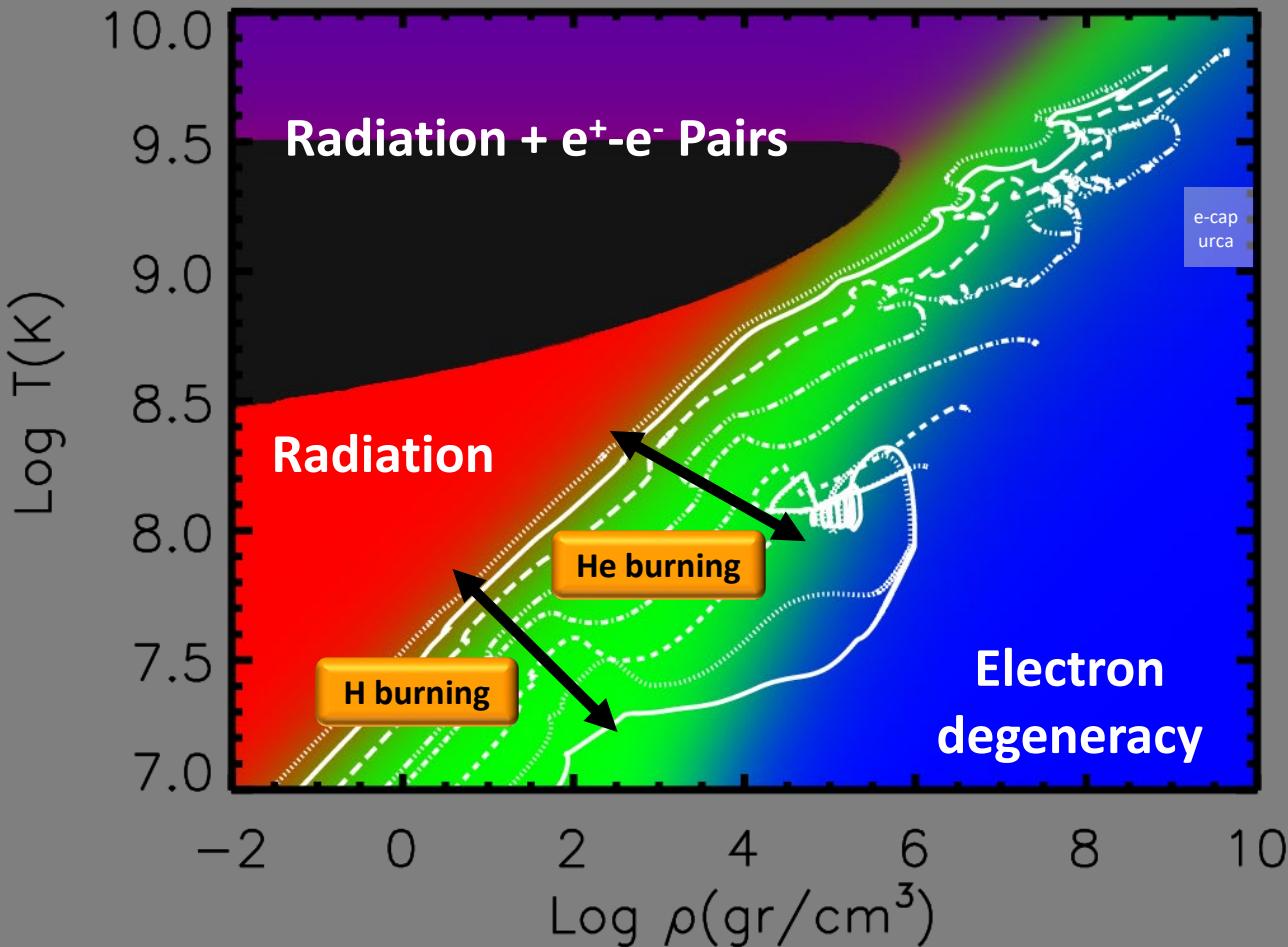
All stars are able to lift the electron degeneracy because:

Energy required to lift the degeneracy at $\text{Log}_{10}(\rho)=6 \Rightarrow 5 \cdot 10^{22} \text{ [erg/cm}^3]$

Energy provided by the 3α at $\text{Log}_{10}(\rho)=6 \Rightarrow 8 \cdot 10^{23} \text{ [erg/cm}^3]$



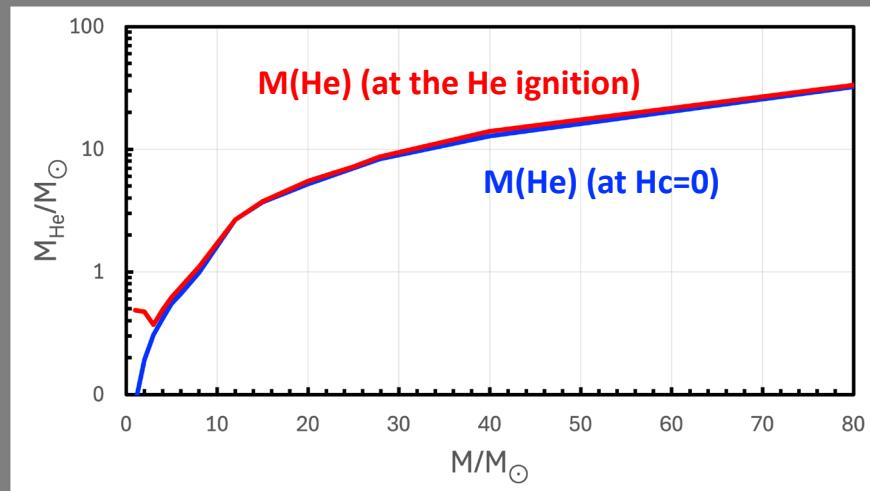
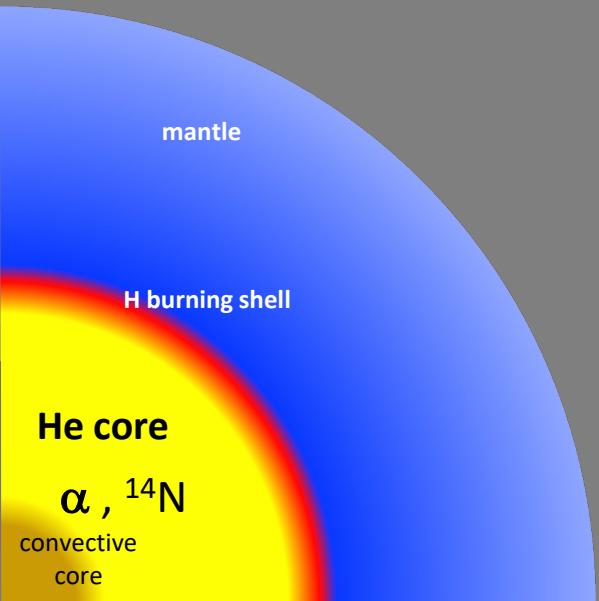




He burning

The key parameter that drives the evolution in He burning is the He core mass and NOT the total mass

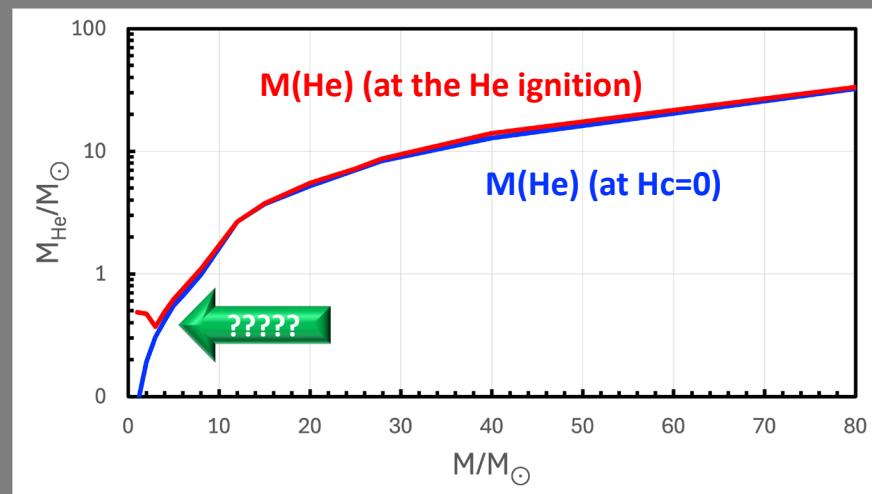
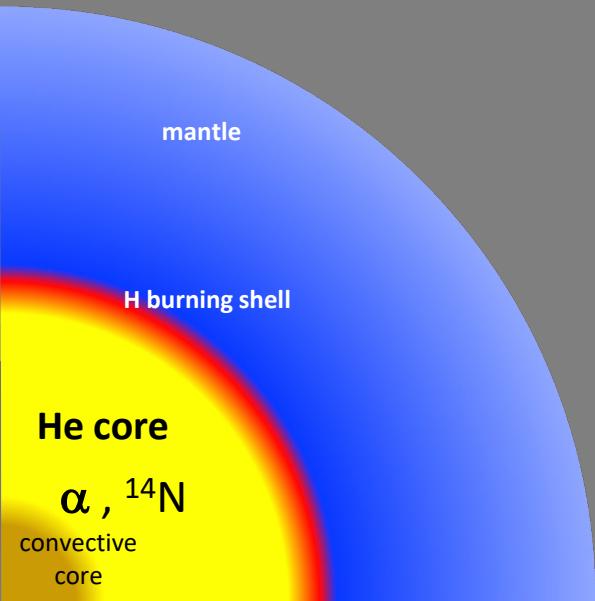
Central He burning always occurs in a convective environment



He burning

The key parameter that drives the evolution in He burning is the He core mass and NOT the total mass

Central He burning always occurs in a convective environment



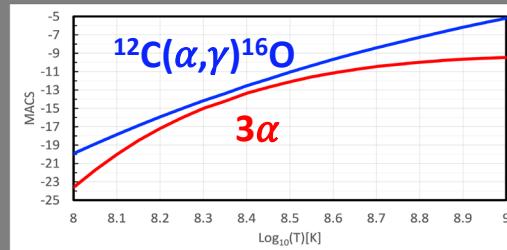
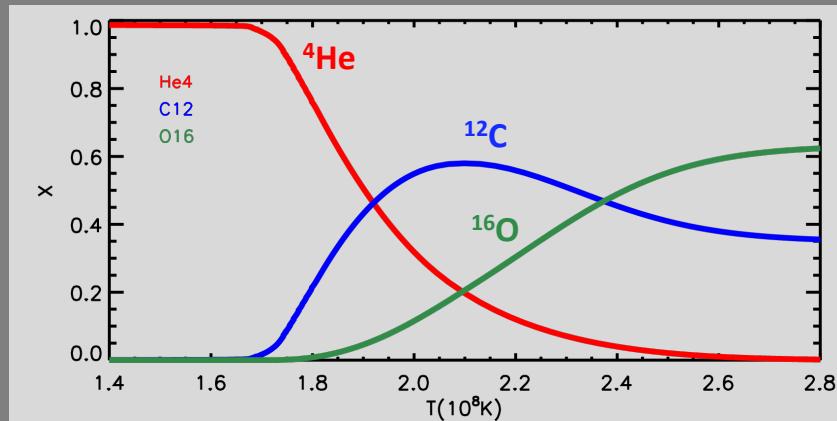
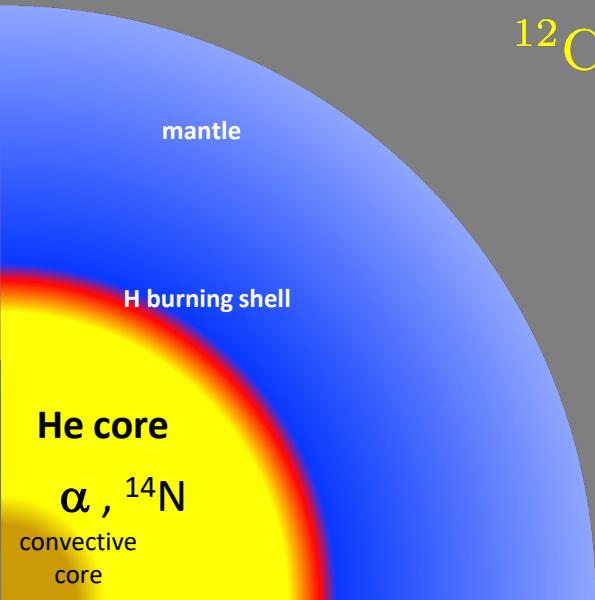
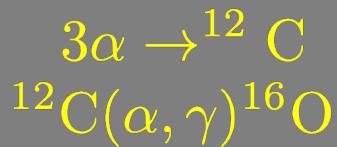
He burning

The key parameter that drives the evolution in He burning is the He core mass and NOT the total mass

Central He burning always occurs in a convective environment

$$T \sim 1.5 - 3.5 \cdot 10^8 \text{ K} \quad \rho \sim 0.2 - 4 \cdot 10^3 \text{ gcm}^{-3}$$

$$4 \alpha \rightarrow {}^{16}\text{O} \quad \Delta M = 4 \times 4.0026 - 15.9949 = 0.015 \text{ MeV} \quad E_{nuc} = 8.70 \cdot 10^{17} \text{ erg/g}$$



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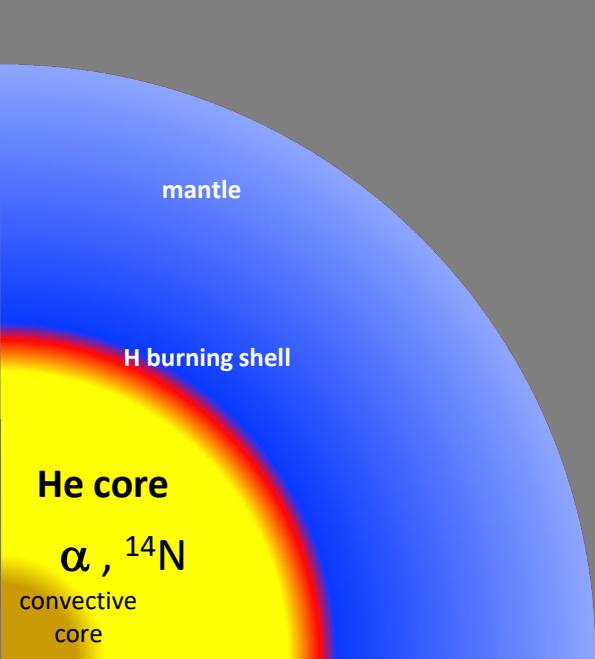
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13	Al	14	Al	15	Al	16	Al	17	Al	18	Al	19	Al	20	Al	21	Al	22	Al	23	Al	24	Al	25	Al	26	Al	27	Al	28	Al	29	Al
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12	Mg	13	Mg	14	Mg	15	Mg	16	Mg	17	Mg	18	Mg	19	Mg	20	Mg	21	Mg	22	Mg	23	Mg	24	Mg	25	Ig	26	Ig	27	Mg	28	Ig



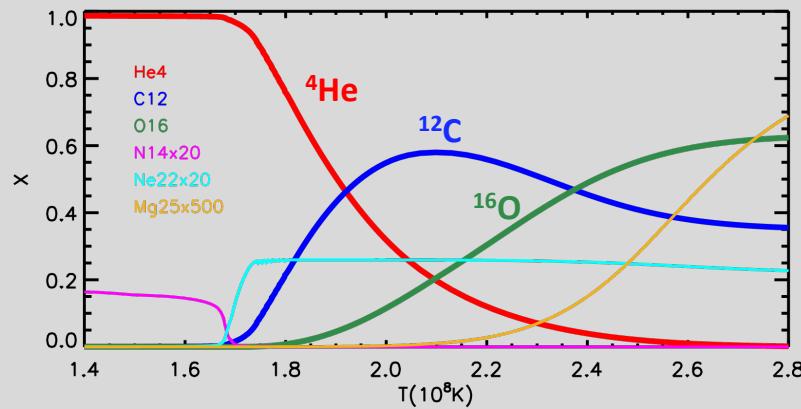
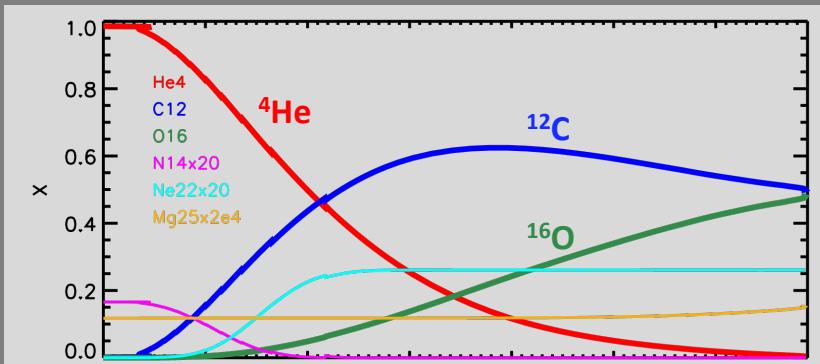
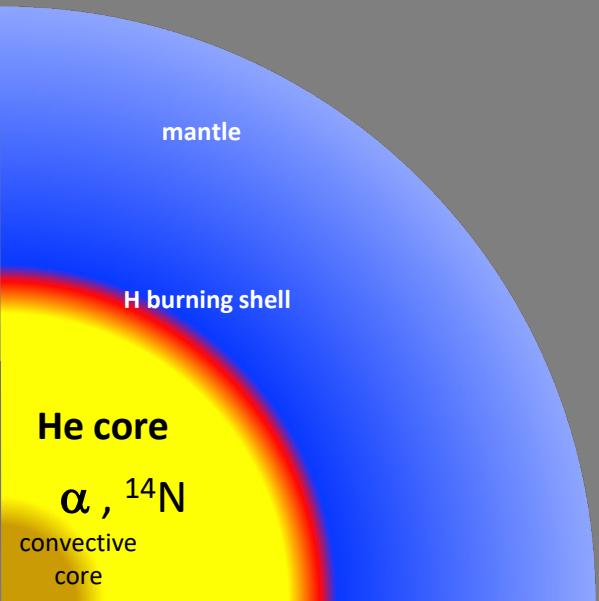
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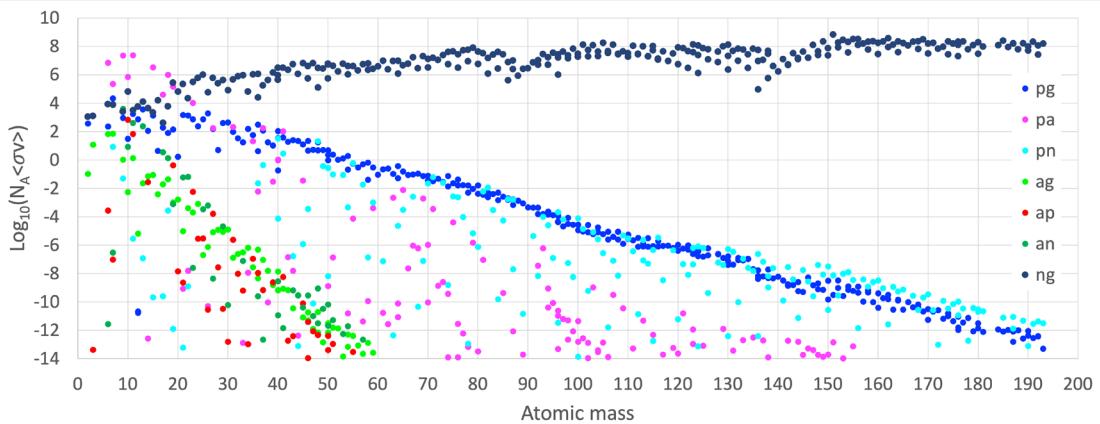


He burning

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Central He burning always occurs in a convective environment





S-process nucleosynthesis (weak component)

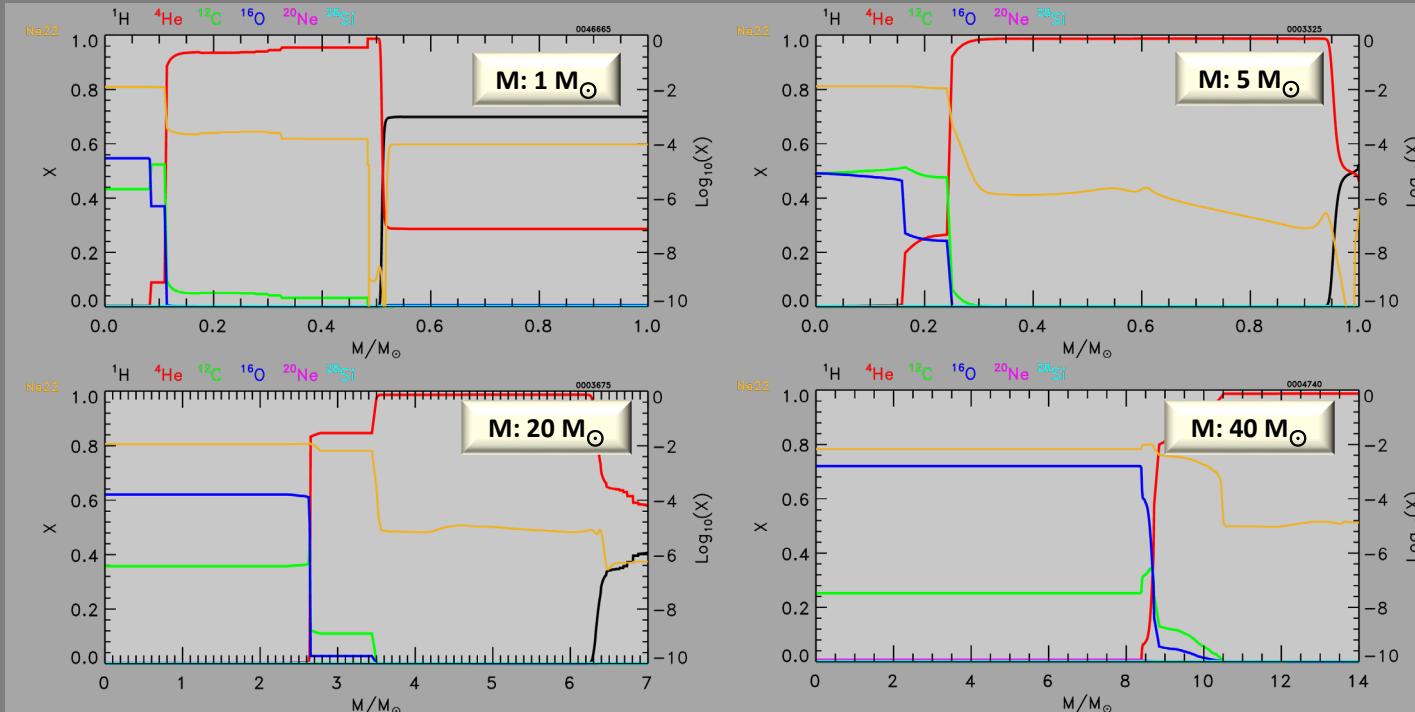


He burning

At the central He exhaustion we do not have any more a monoparametric family of stars but a Bi-parametric family.

In fact from now on there are two leading parameters that drive the further evolution of each star:

the CO core mass and the amount of C left by the He burning.

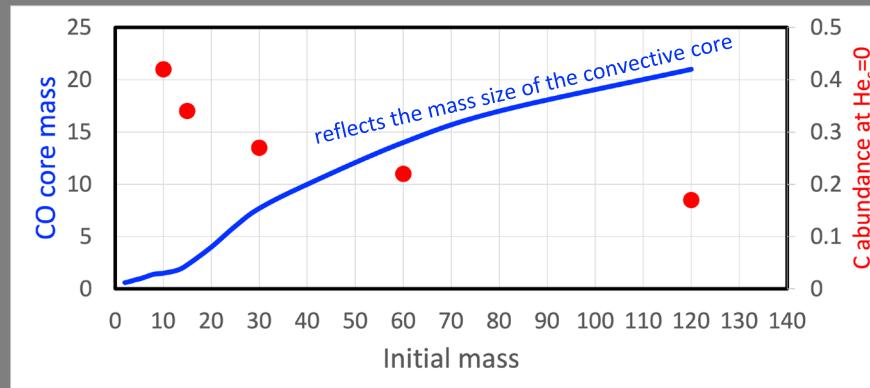


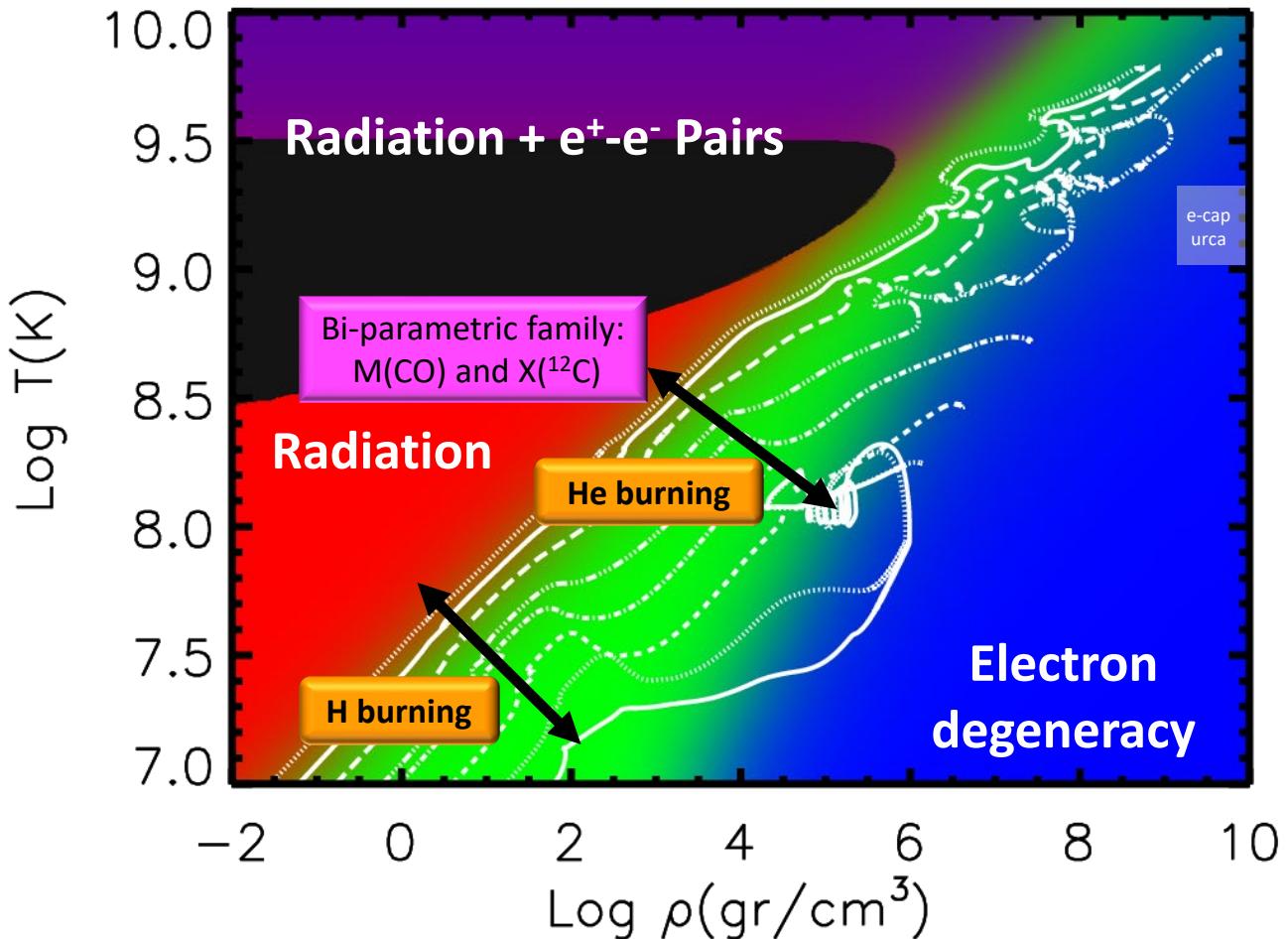
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Let us discuss the advanced phases of massive stars first

EVOLUTION OF MASSIVE STARS

R. J. TAYLER*

Princeton University Observatory

Received March 15, 1954

ABSTRACT

The evolution is considered of massive stars with opacity due to electron scattering. As the stars evolve, the convective core retreats, and a zone of continuously varying composition is set up between the core and the envelope. Ten models have been obtained with core hydrogen content changing from 100 to 6 per cent. The evolutionary tracks of the models have been plotted in the H-R diagram. It is found that, although the individual tracks are very different from those found by other authors, the H-R diagram for stars of different masses but of the same age exhibits the well-known "knee," which occurs at the point of 11 per cent over-all reduction of hydrogen content.

I. INTRODUCTION

The course of stellar evolution is largely determined by the existence or nonexistence of general mixing currents in stellar interiors. If there are efficient mixing currents, a star



Alice crosses the mirror

neutrino energy losses

Plasma neutrinos

$$\gamma \rightarrow \nu_e \bar{\nu}_e$$

Photo neutrinos

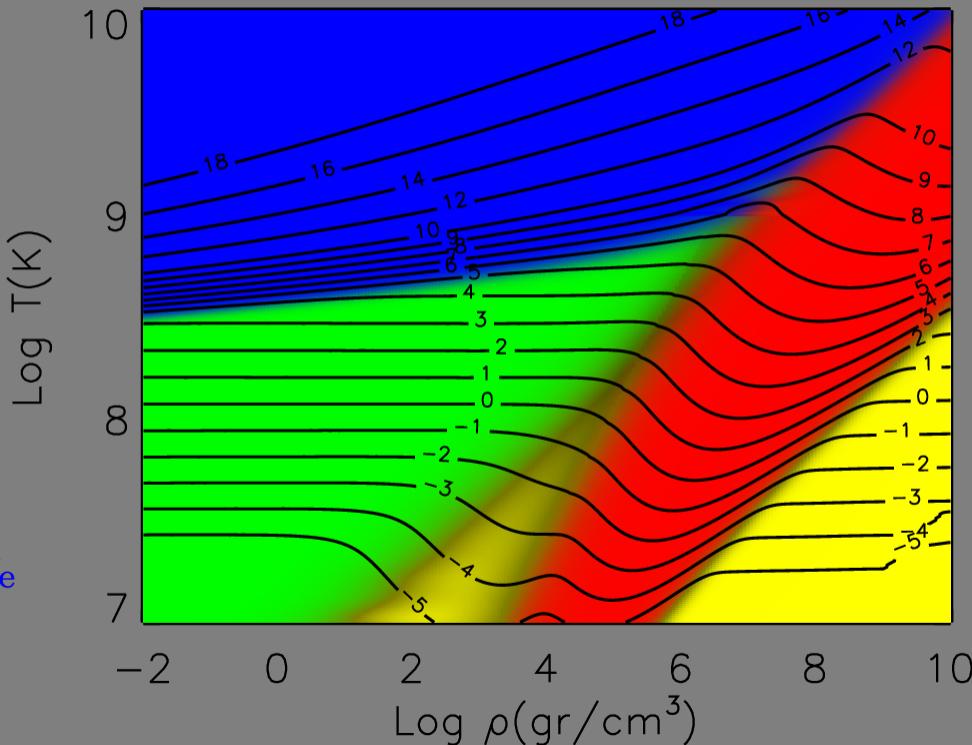
$$\gamma + e^- \rightarrow e^- \nu_e \bar{\nu}_e$$

Bremstrahlung

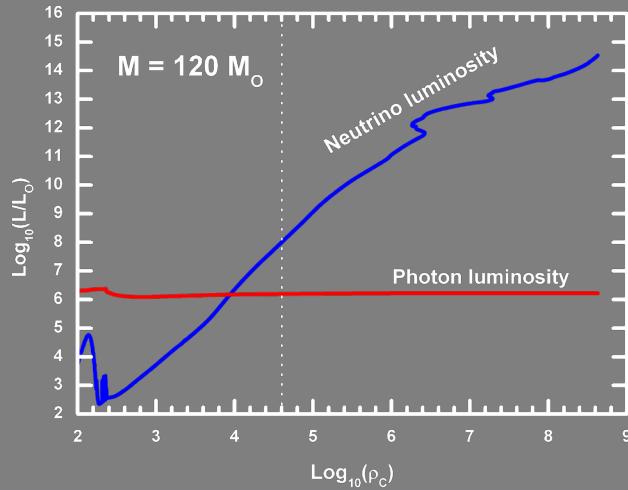
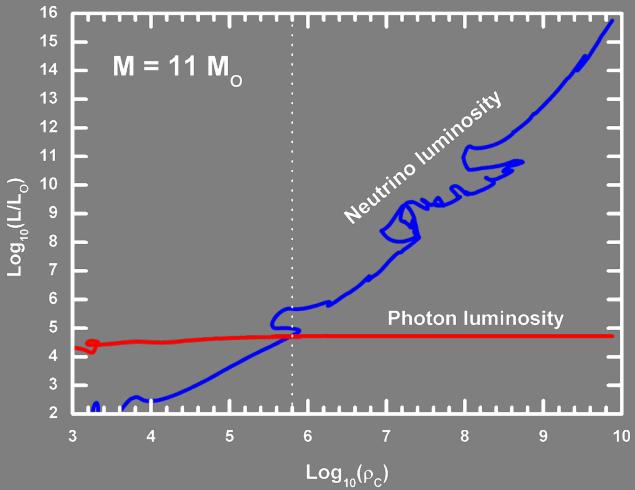
$$\gamma e^- \rightarrow \gamma e^- \nu_e \bar{\nu}_e$$

Pair neutrinos

$$\gamma\gamma \rightarrow e^+ e^- \rightarrow \nu_e \bar{\nu}_e$$



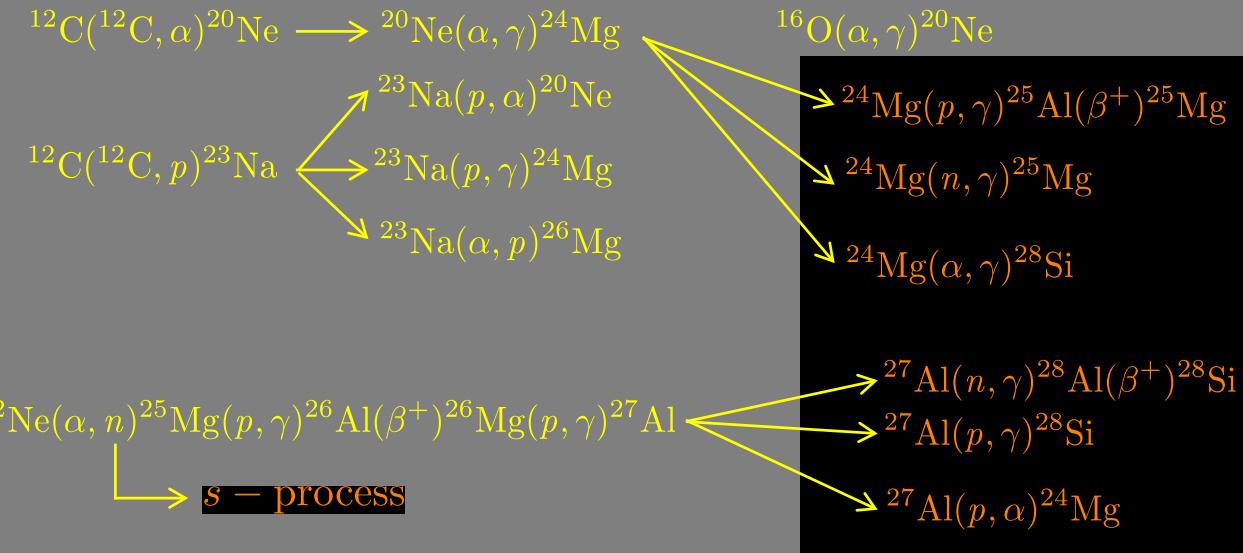
Pair neutrinos



Massive Stars: Carbon Burning

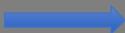
Following He burning the most abundant isotopes are $^{12}\text{C}, ^{16}\text{O}$

$$T \simeq M^{2/3} \rho^{1/3} \quad \longrightarrow \quad T \sim 8 \cdot 10^8 \text{ K} \quad \rho \sim 8 \cdot 10^5 \text{ g/cm}^3$$



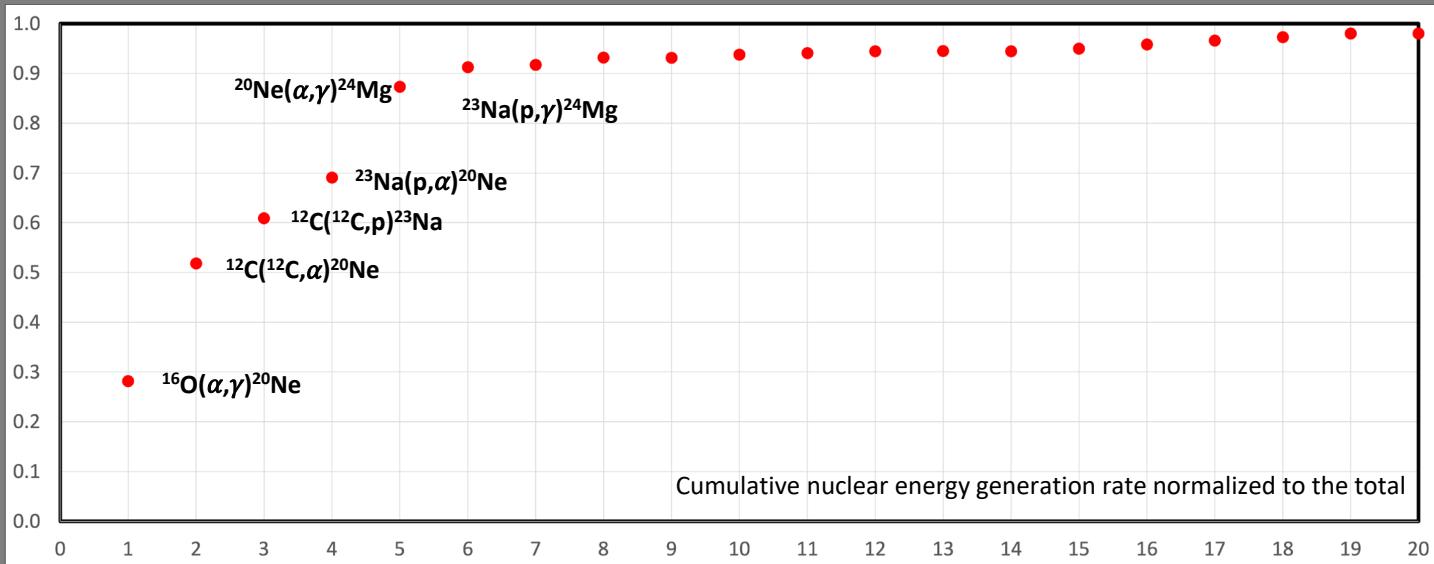
Main Products of C burning $^{20}\text{Ne}, ^{23}\text{Na}, ^{24}\text{Mg}, ^{27}\text{Al}$

Secondary Products of C burning $^{25}\text{Mg}, ^{26}\text{Mg}, s\text{-process}$



$$E_{\text{nuc}} = 1.85 \cdot 10^{17} \text{ erg/g}$$

C burning



$^{12}\text{C} + ^{12}\text{C}$

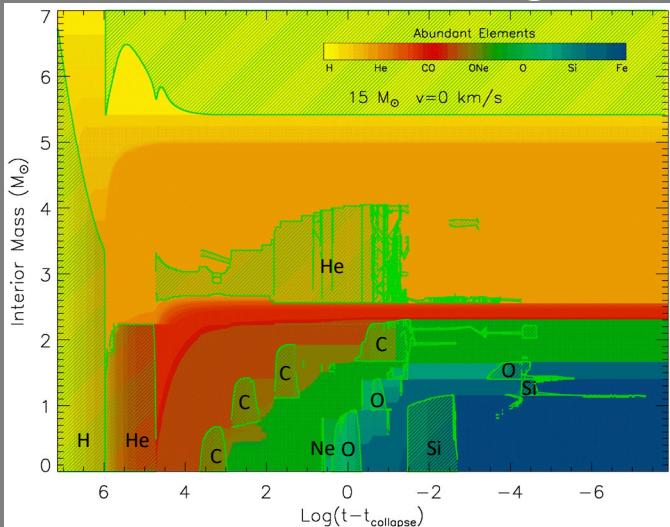
A photograph of a small, weathered building in a desert landscape, likely a former gas station. The building has a blue door and a blue garage door. Two blue vintage-style gas pumps stand in front. A yellow sign above the pumps displays the nuclear reaction $^{12}\text{C} + ^{12}\text{C}$. A black tire lies on the ground near the pumps. In the background, numerous tall saguaro cacti stand against a bright blue sky with wispy white clouds. A small windmill is mounted on top of the building's roofline.

α

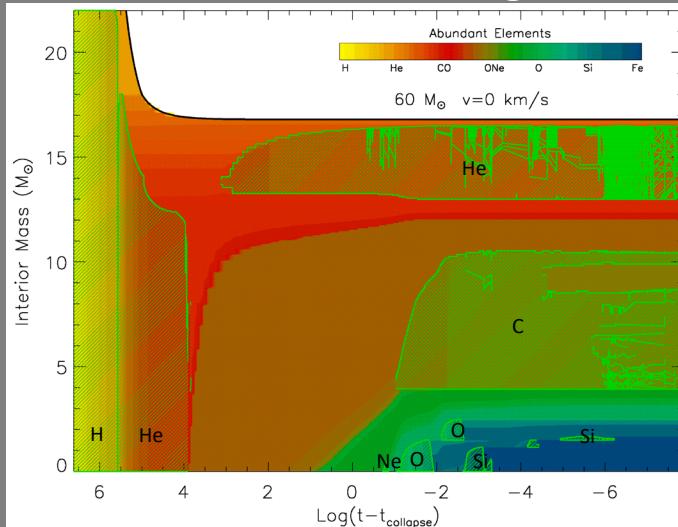
p

Massive Stars: Carbon Burning

$M < 25/30 M_{\odot}$

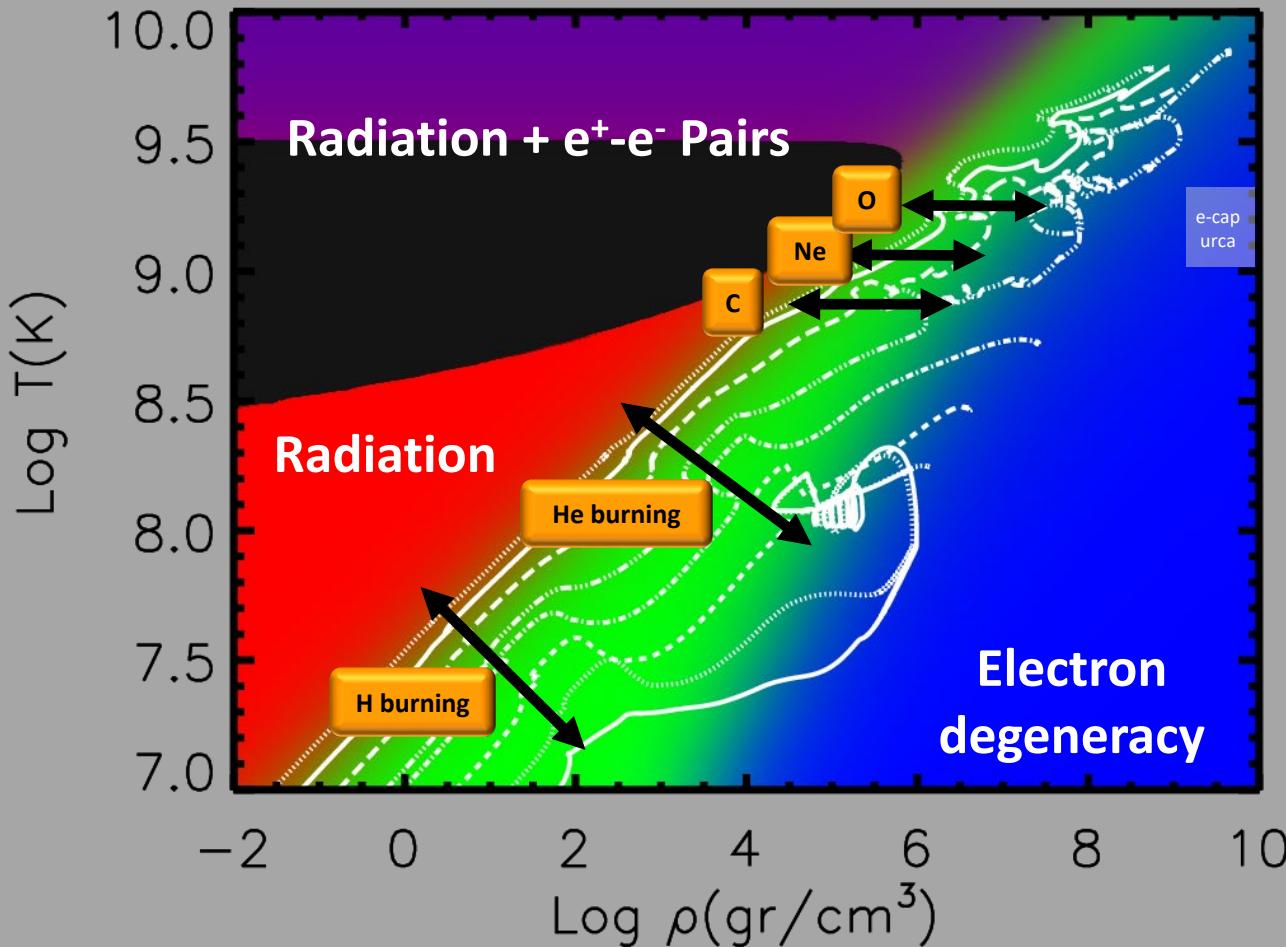


$M > 25/30 M_{\odot}$



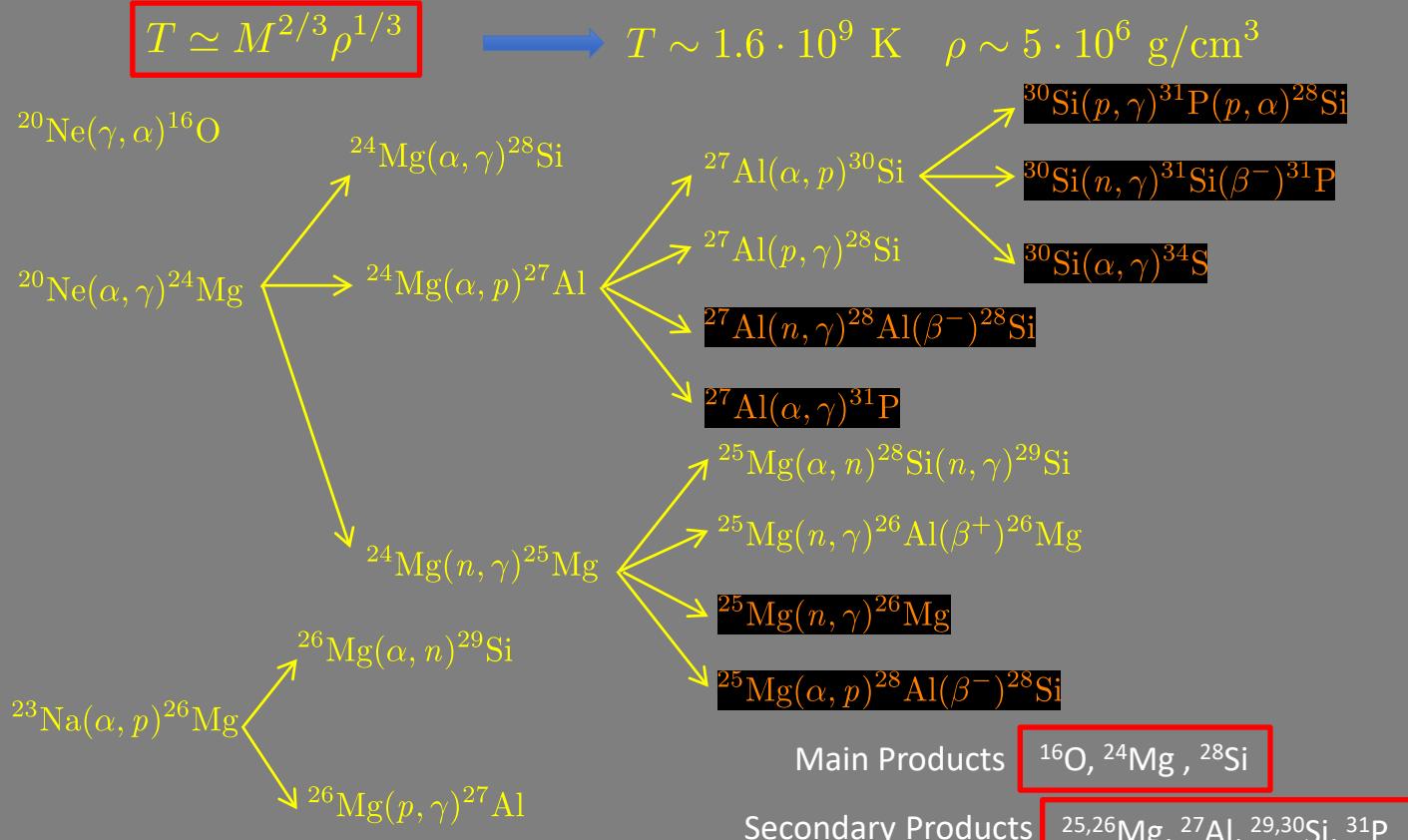
$L \Rightarrow$ total luminosity: $L_{\gamma} + L_v$

$M=80 M_{\odot}$ $t = E / L(10^6)$		M_{CC}	Estimated lifetime	Real lifetime	Revised lifetime	L_{TOT}
$H \Rightarrow He$	$6.44 \cdot 10^{18} \text{ erg gr}^{-1}$	60	$6 \cdot 10^6$	$3.2 \cdot 10^6$		10^6
$He \Rightarrow C$	$5.84 \cdot 10^{17} \text{ erg gr}^{-1}$	20	$2 \cdot 10^5$	$3.3 \cdot 10^5$		10^6
$C \Rightarrow Ne$	$1.85 \cdot 10^{17} \text{ erg gr}^{-1}$	1.5	$4.5 \cdot 10^3$	$4.7 \cdot 10^2$	$4.5 \cdot 10^2$	10^7
$O \Rightarrow Si$	$2.89 \cdot 10^{17} \text{ erg gr}^{-1}$	1	$4.8 \cdot 10^3$	$4.6 \cdot 10^{-2}$	$4.8 \cdot 10^{-2}$	10^{11}
$Si \Rightarrow Ni$	$1.88 \cdot 10^{17} \text{ erg gr}^{-1}$	1	$3.1 \cdot 10^3$	$4.3 \cdot 10^{-3}$	$3.1 \cdot 10^{-3}$	10^{12}



Massive Stars: Neon Burning

Following C burning the most abundant isotopes are ^{16}O (He-burn.), ^{20}Ne , ^{23}Na , ^{24}Mg

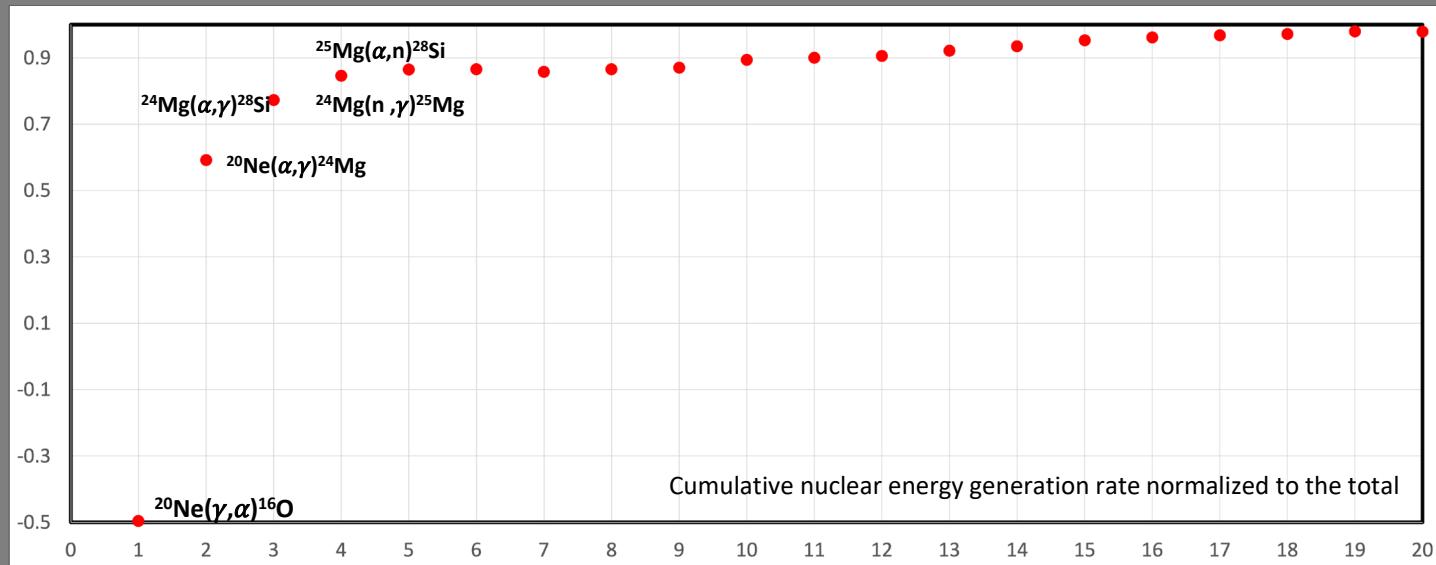


$$2 \ ^{20}\text{Ne} \rightarrow ^{16}\text{O} + ^{24}\text{Mg}$$

$$\longrightarrow$$

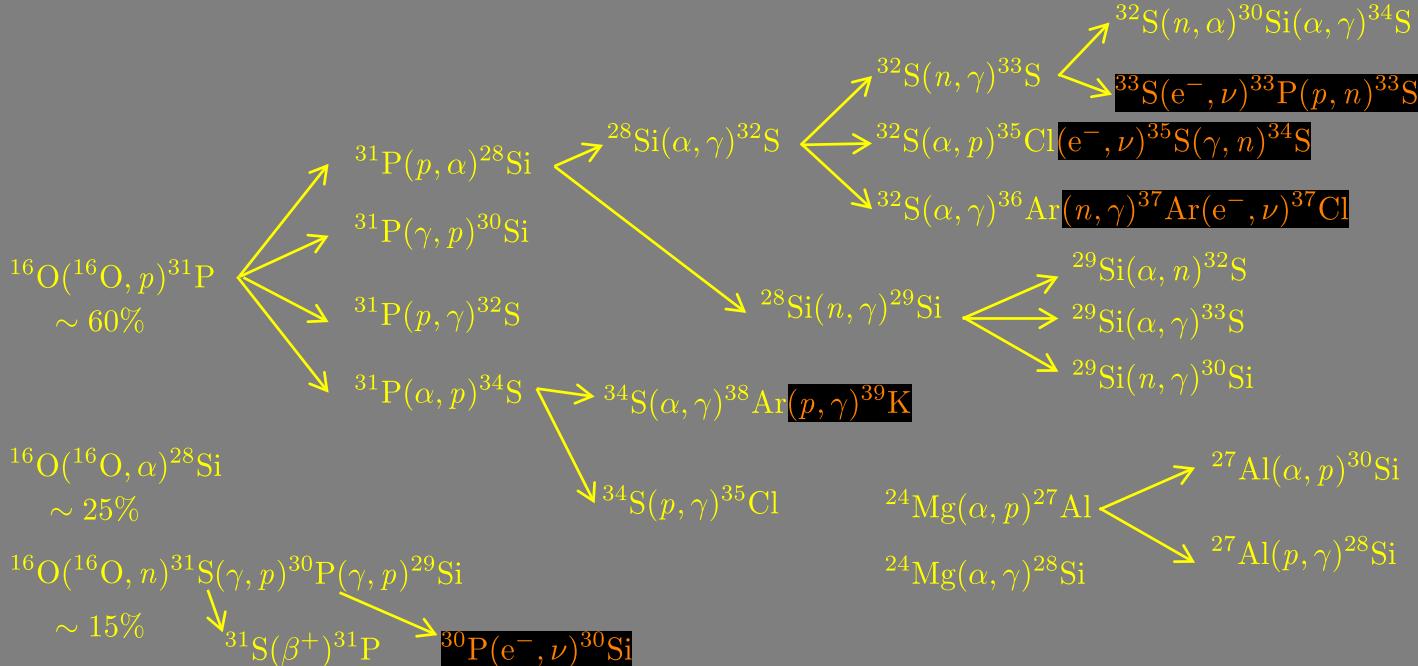
$$E_{\text{nuc}} = 1.10 \cdot 10^{17} \text{ erg/g}$$

Ne burning

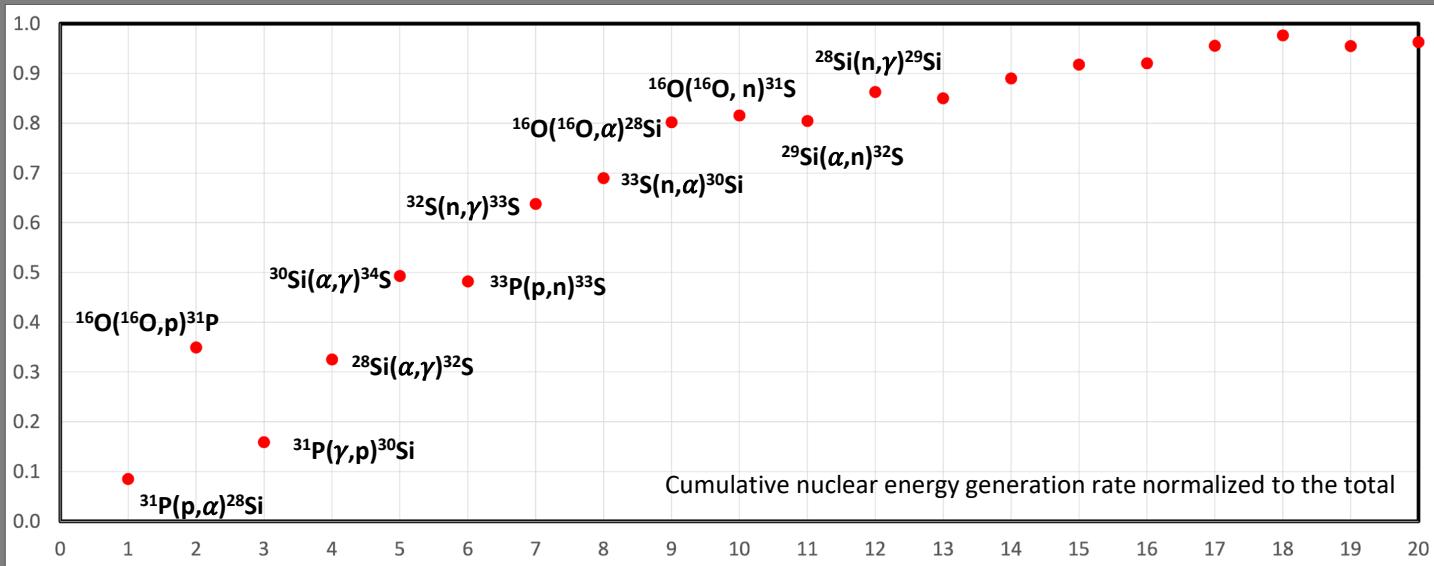


Massive Stars: Oxygen Burning

The most abundant nuclei left by the Ne burning are: ^{16}O , ^{24}Mg , ^{28}Si

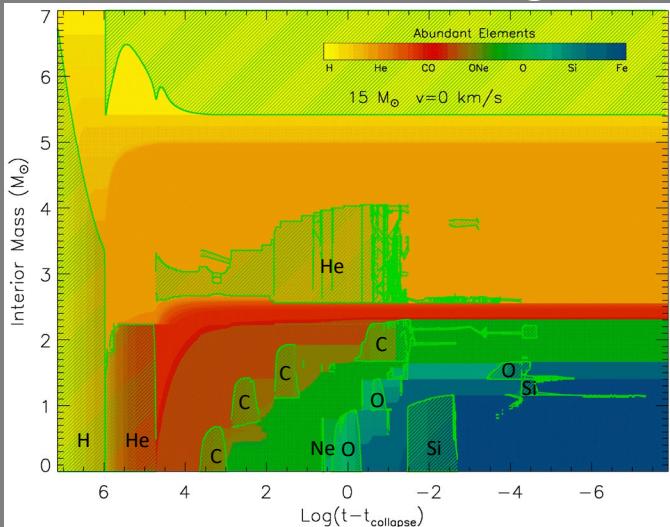


O burning

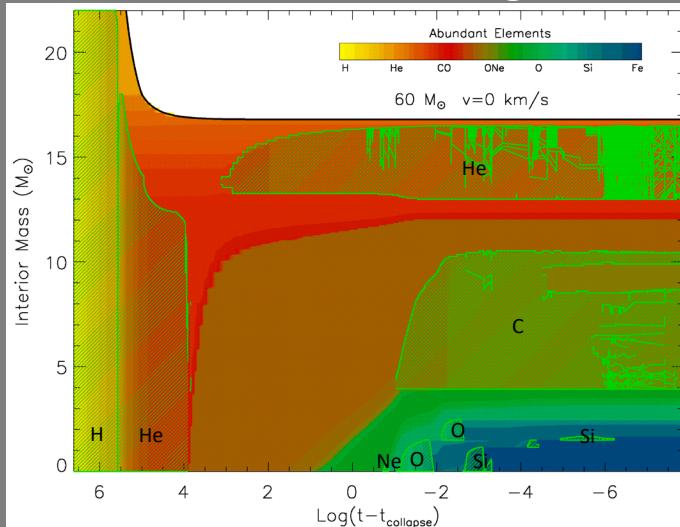


Massive Stars: Carbon Burning

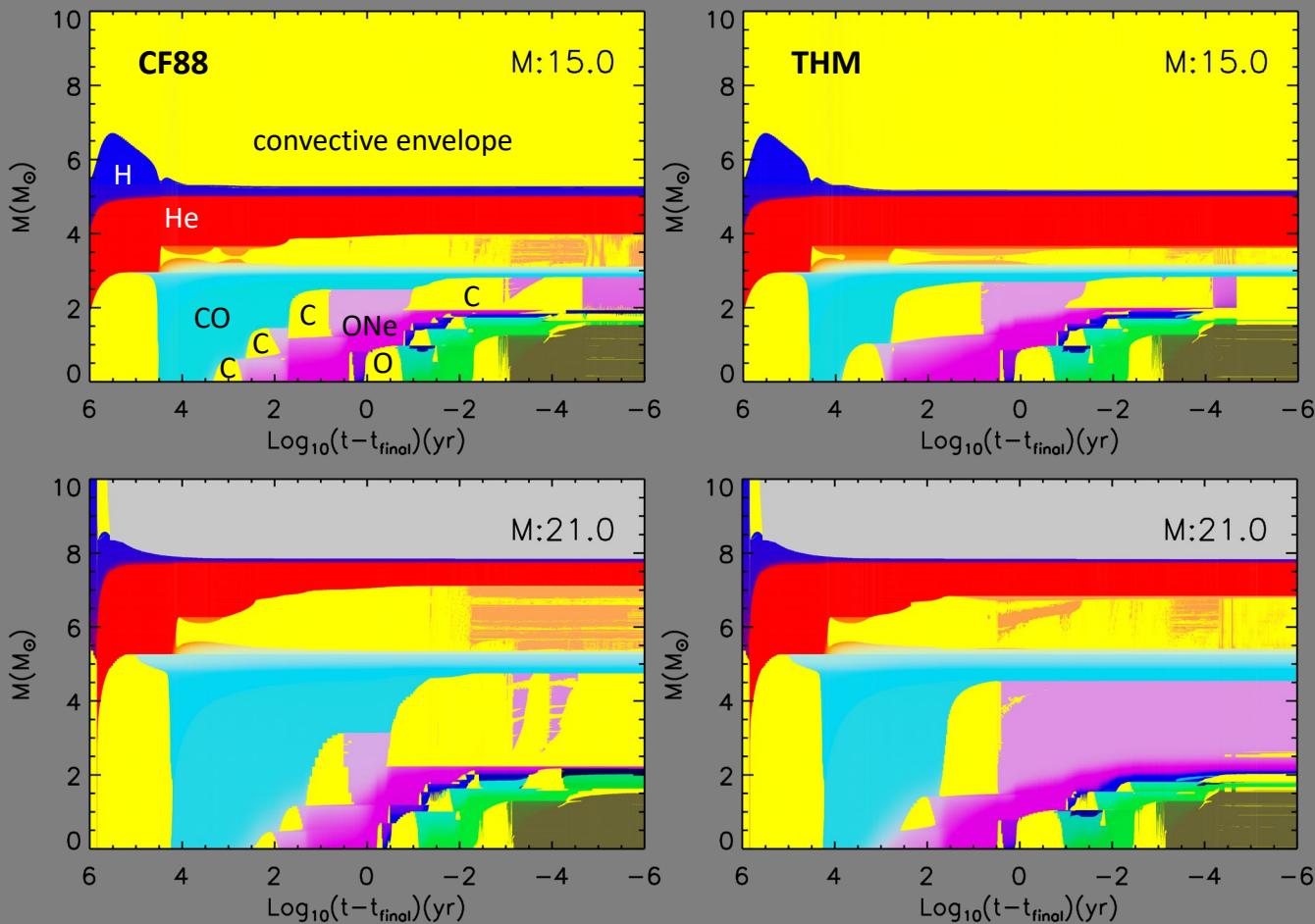
$M < 25/30 M_{\odot}$



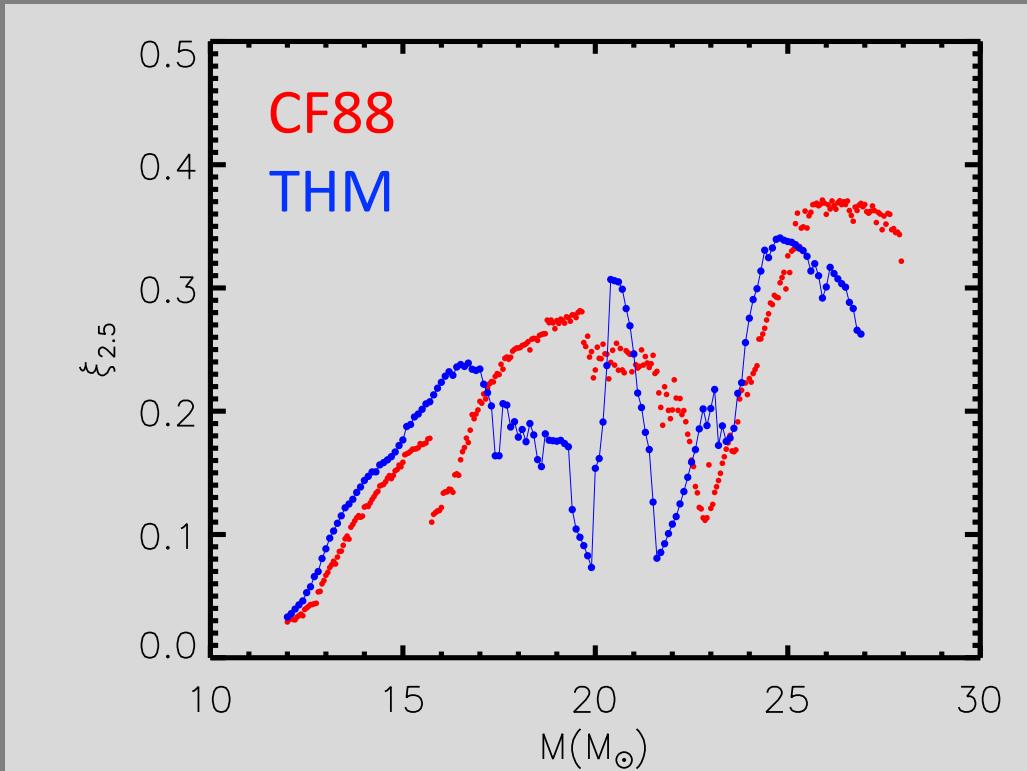
$M > 25/30 M_{\odot}$



The advanced burning

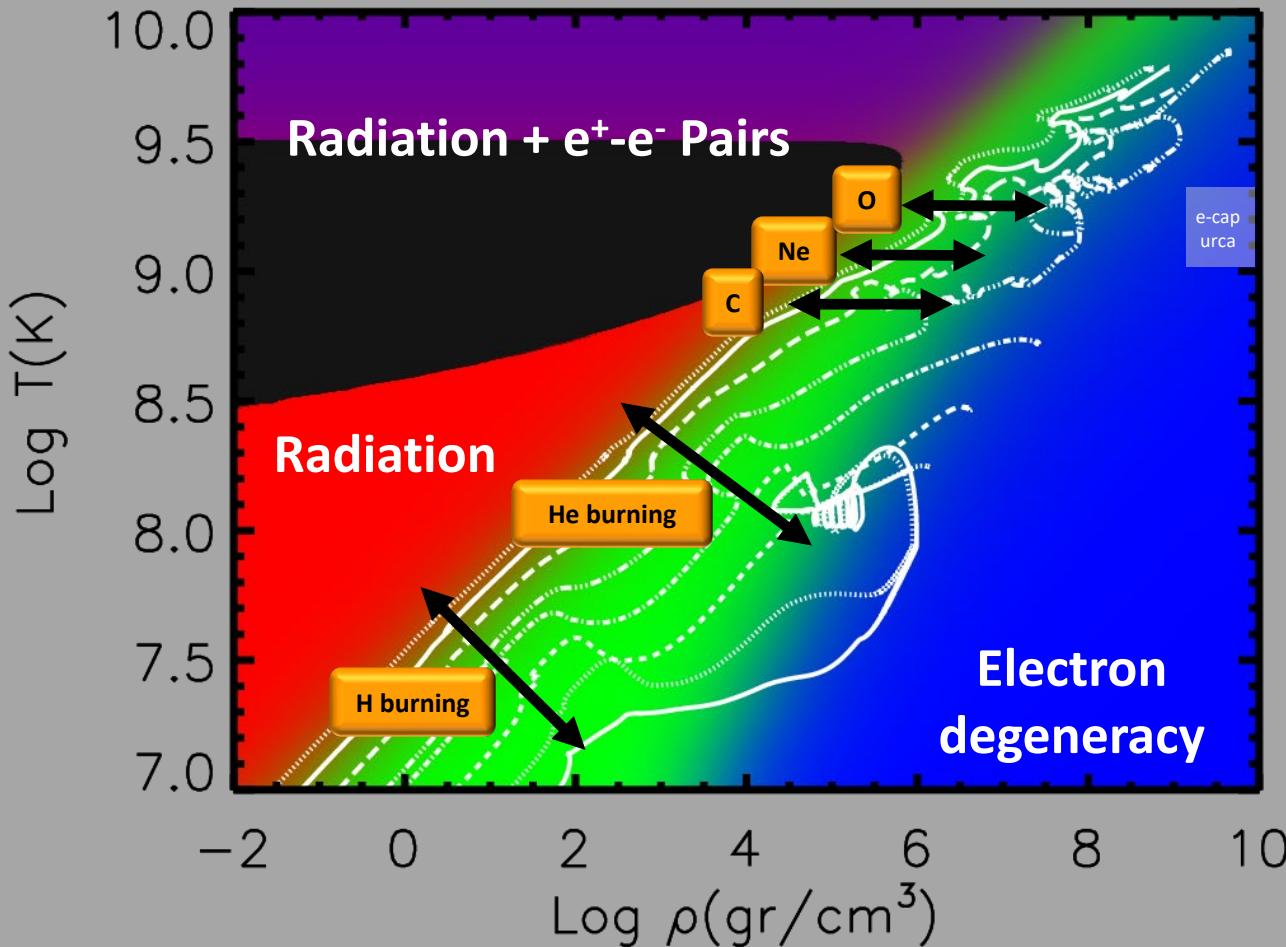


The final compactness



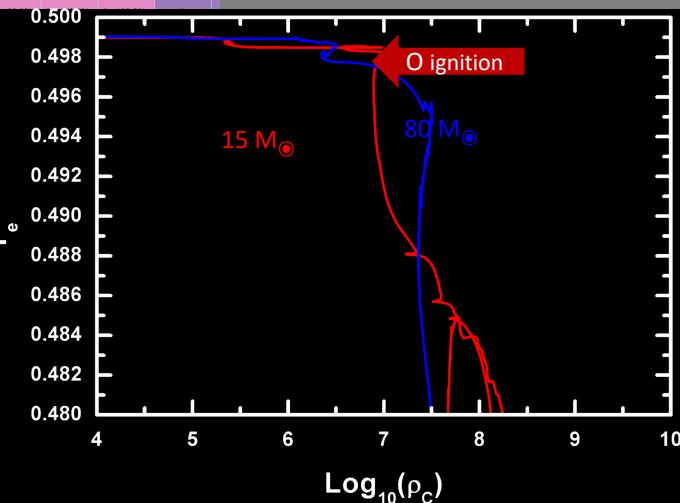
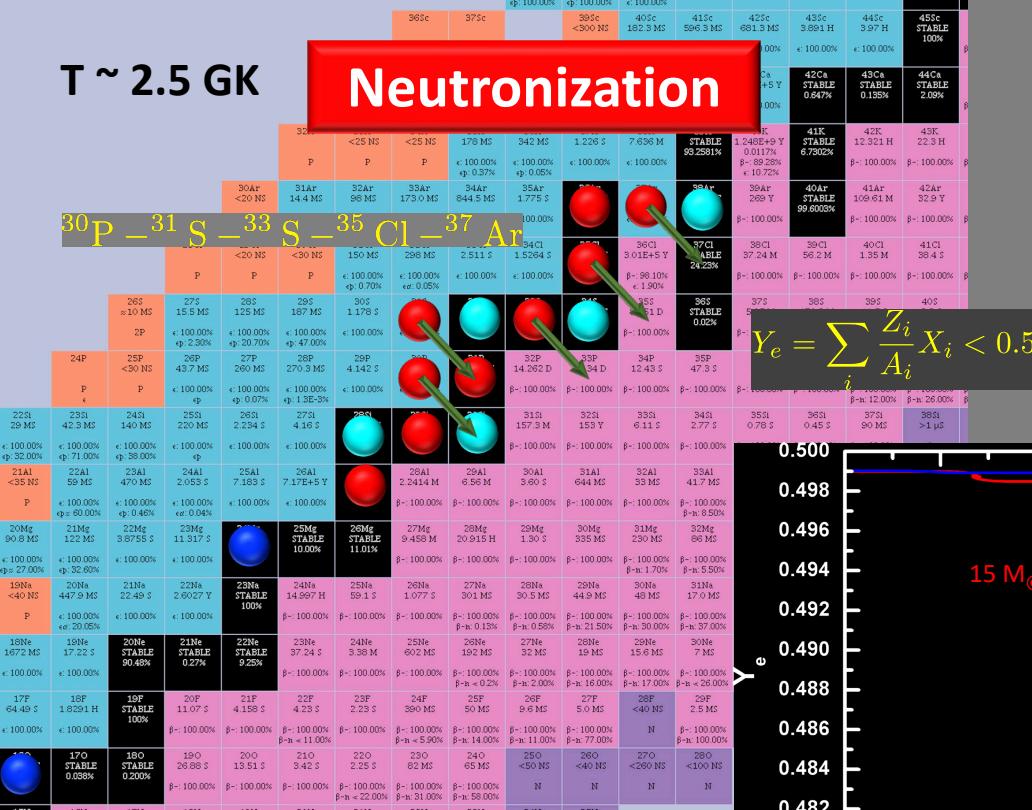
Compactness parameter ξ
(O'Connor & Ott 2011, ApJ 730, 70)

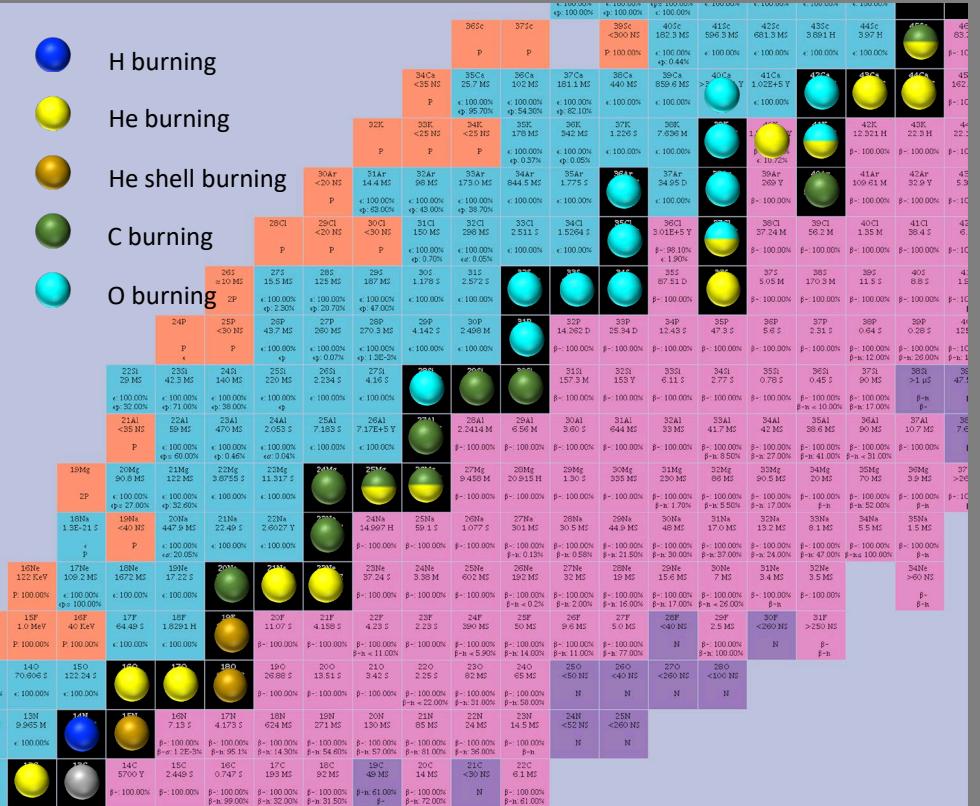
$$\xi_i = \frac{M_i(M_\odot)}{R_i(10^3 \text{ km})}$$

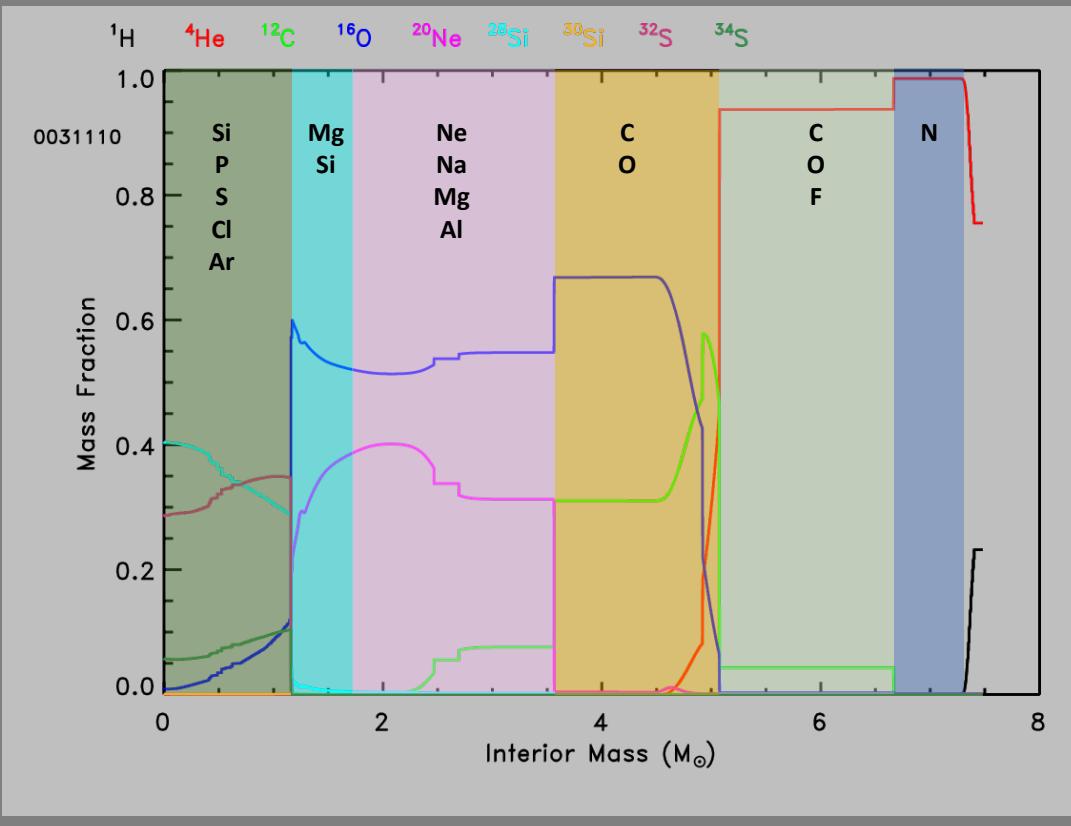


T ~ 2.5 GK

Neutronization







Beyond the O burning matter approaches progressively the Nuclear Statistical Equilibrium

At the central O exhaustion => $T \sim 2.5 \cdot 10^9$ K

direct



reverse



$$\langle \sigma v \rangle_{jl} \propto \langle \sigma v \rangle_{ik} e^{-\frac{Q_{ik}[\text{MeV}]}{0.086 T_9}}$$

$$R_{ik} \sim R_{jl}$$

To quantify how close a pair of processes is to the equilibrium let us define a parameter φ :

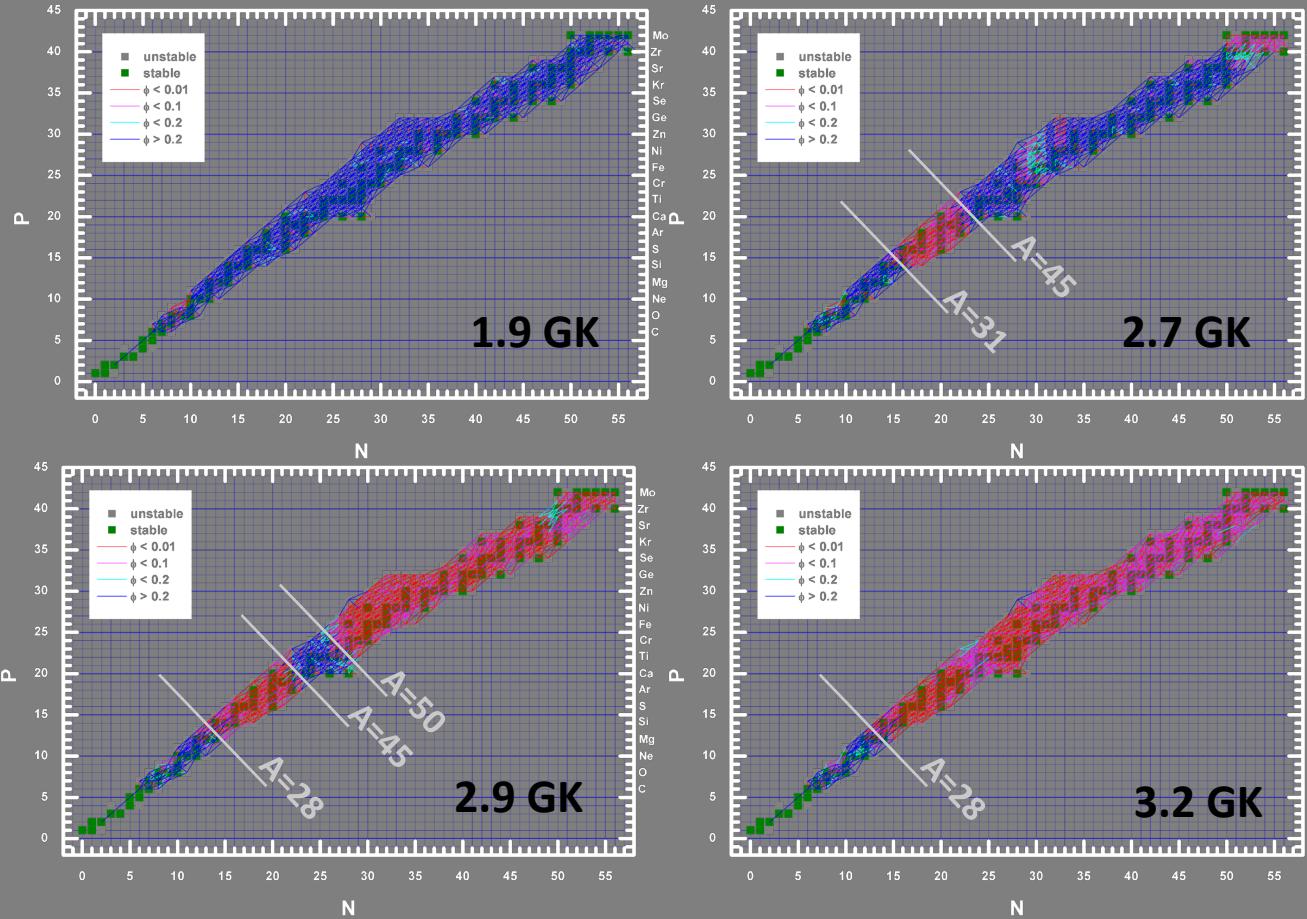
$$\varphi(i,j) = \frac{|r_{ij} - r_{ji}|}{\max(r_{ij}, r_{ji})}$$

obviously

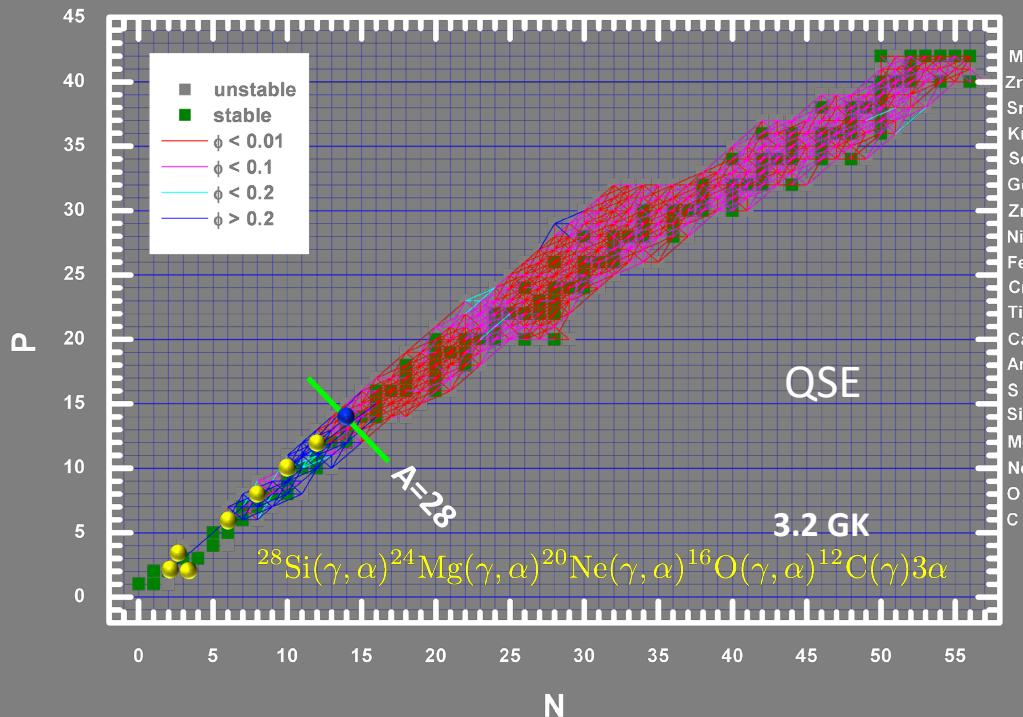
$\varphi \rightarrow 0$ full equilibrium

$\varphi \rightarrow 1$ no equilibrium

Beyond O burning...the path towards the Nuclear Statistical Equilibrium



...the approach to the Nuclear Statistical Equilibrium



Nuclear Statistical Equilibrium

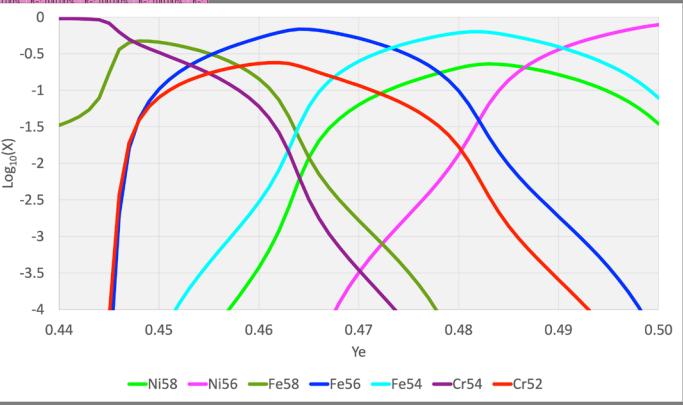
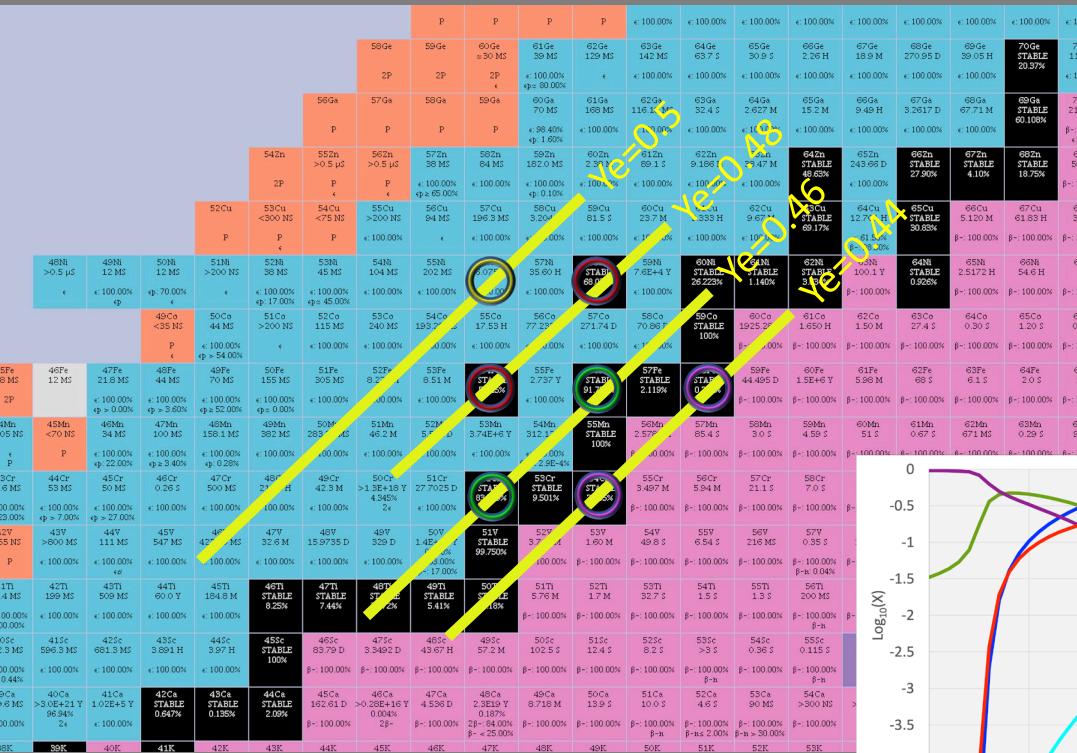
At $T > 5 \text{ GK}$ full equilibrium between direct and reverse processes

$$(N, Z) \rightleftharpoons Zp + Nn$$

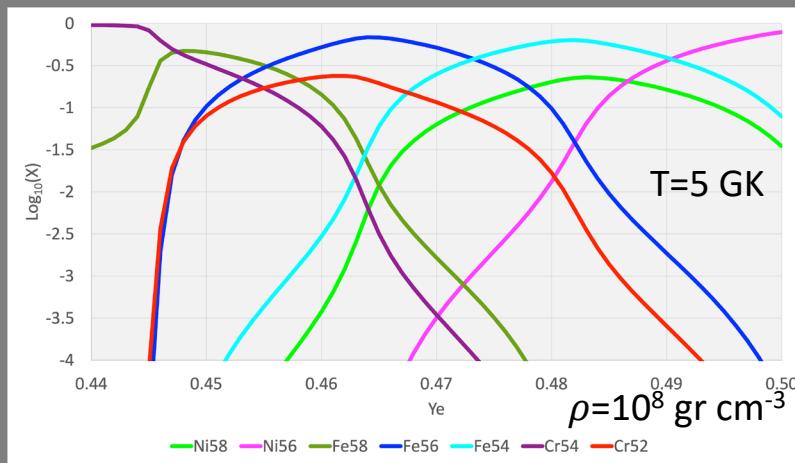
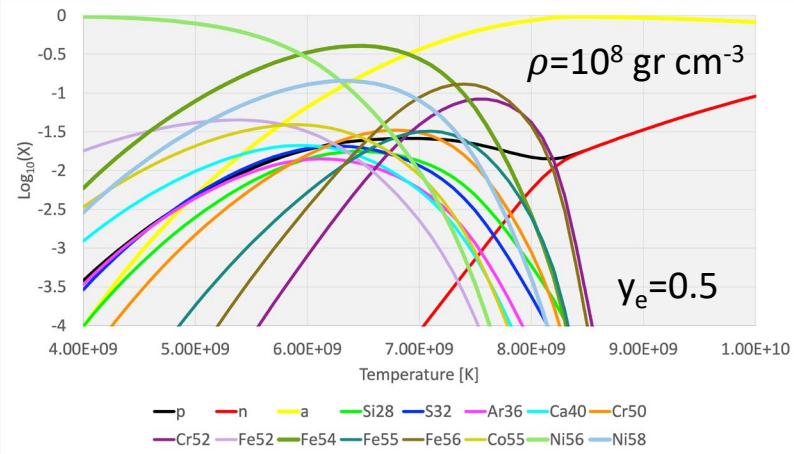
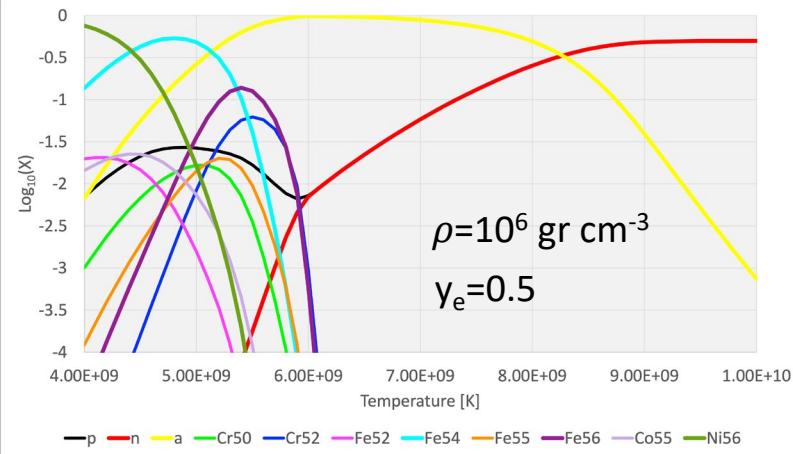
$$y_{i(Z,N)} = \omega(z, n) (\rho N_A)^{A-1} \left(\frac{A m_p k T}{2\pi \hbar^2} \right)^{3/2} y_p^z y_n^n 2^{-A} \left(\frac{2\pi \hbar^2}{m_p k T} \right)^{A/2} e^{-\frac{Q(z,n)}{kT}}$$

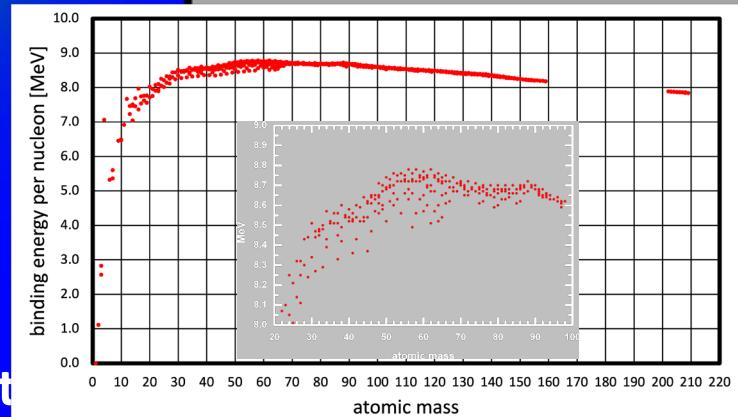
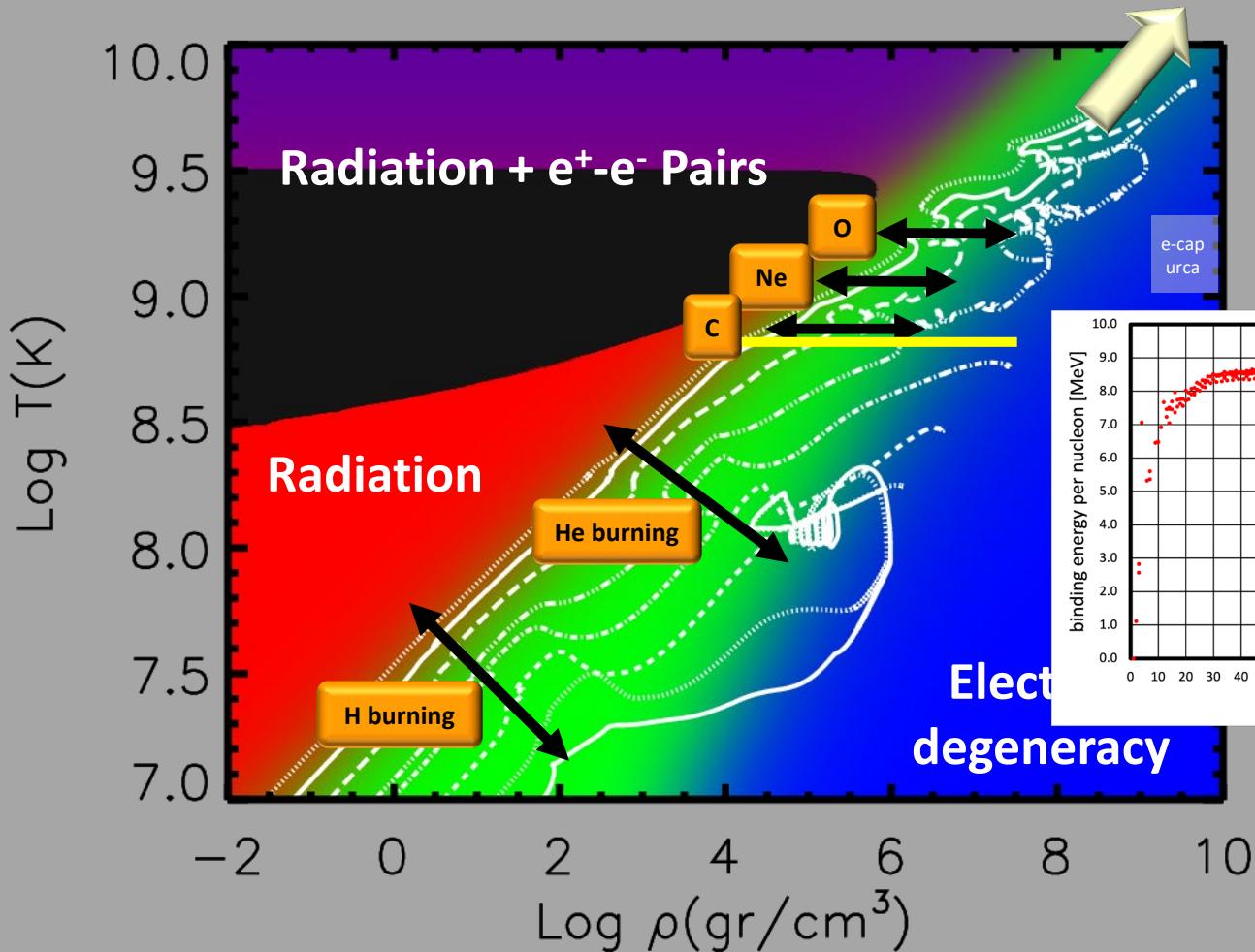
$$y_{i(Z,N)} = f(A, T, \rho) y_p^z y_n^n e^{-\frac{Q(z,n)}{kT}}$$

Nuclear Statistical Equilibrium



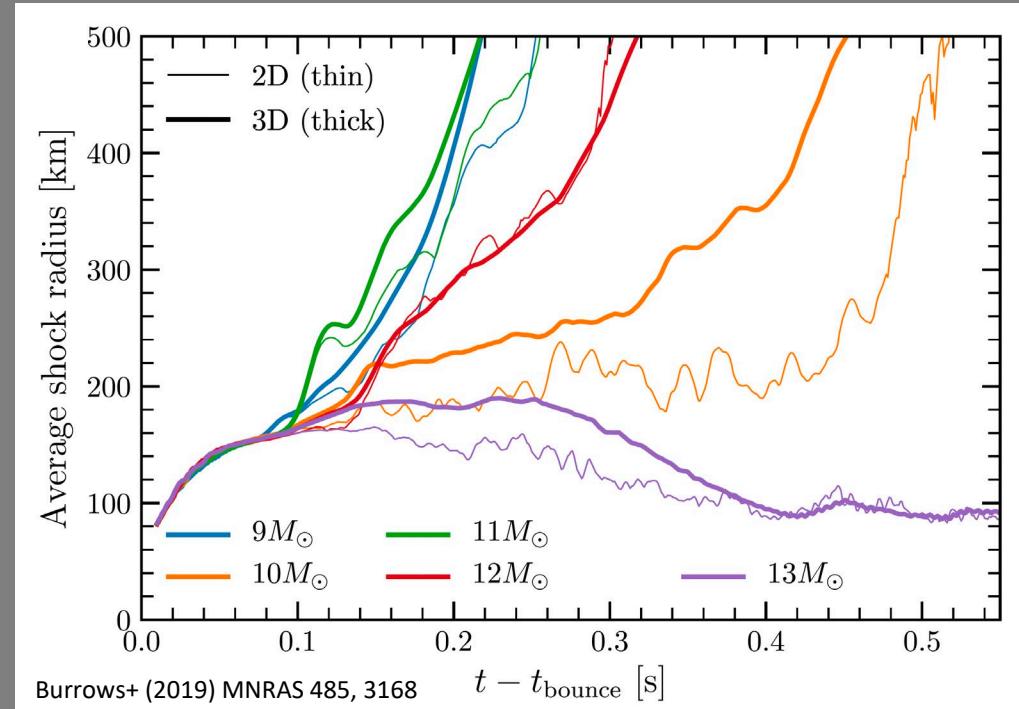
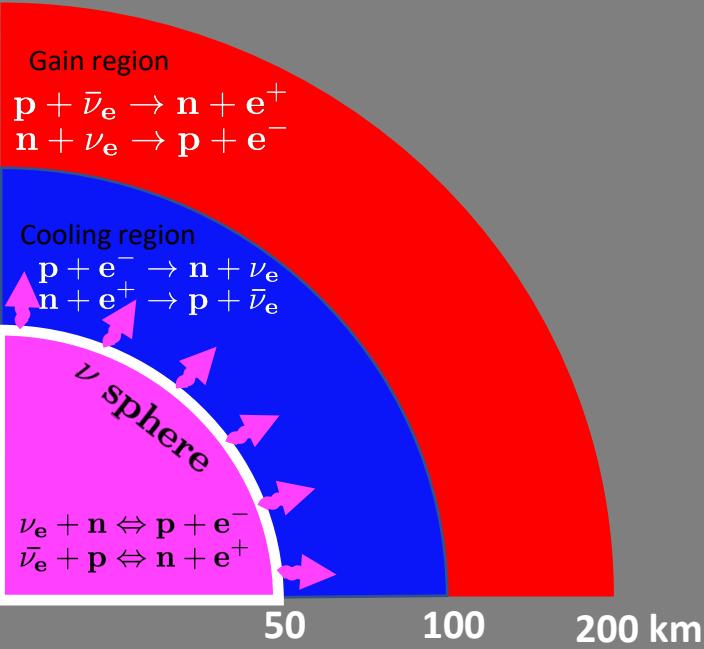
Nuclear Statistical Equilibrium





Total amount of energy released by the gravitational collapse amounts to, roughly:

$$E = 1.6 \cdot 10^{53} \text{ erg}$$



Total amount of energy released by the gravitational collapse amounts to, roughly: $E = 1.6 \cdot 10^{53}$ erg

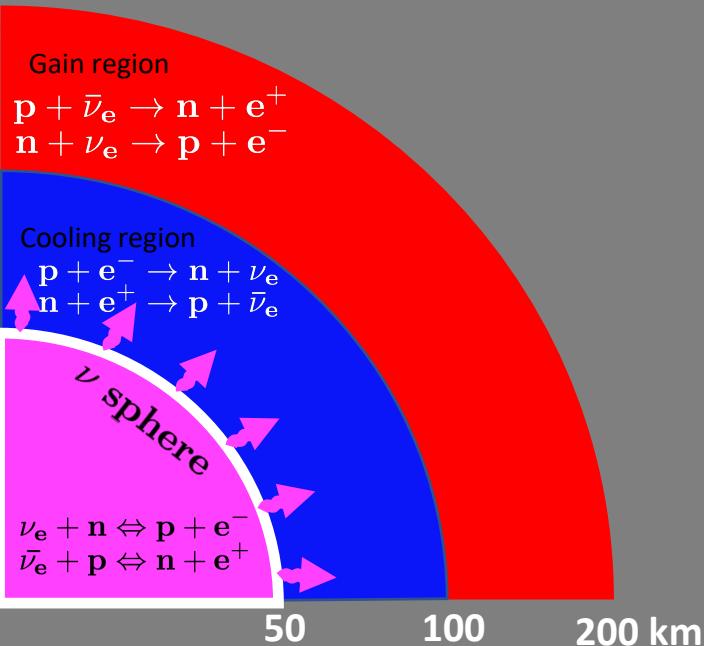
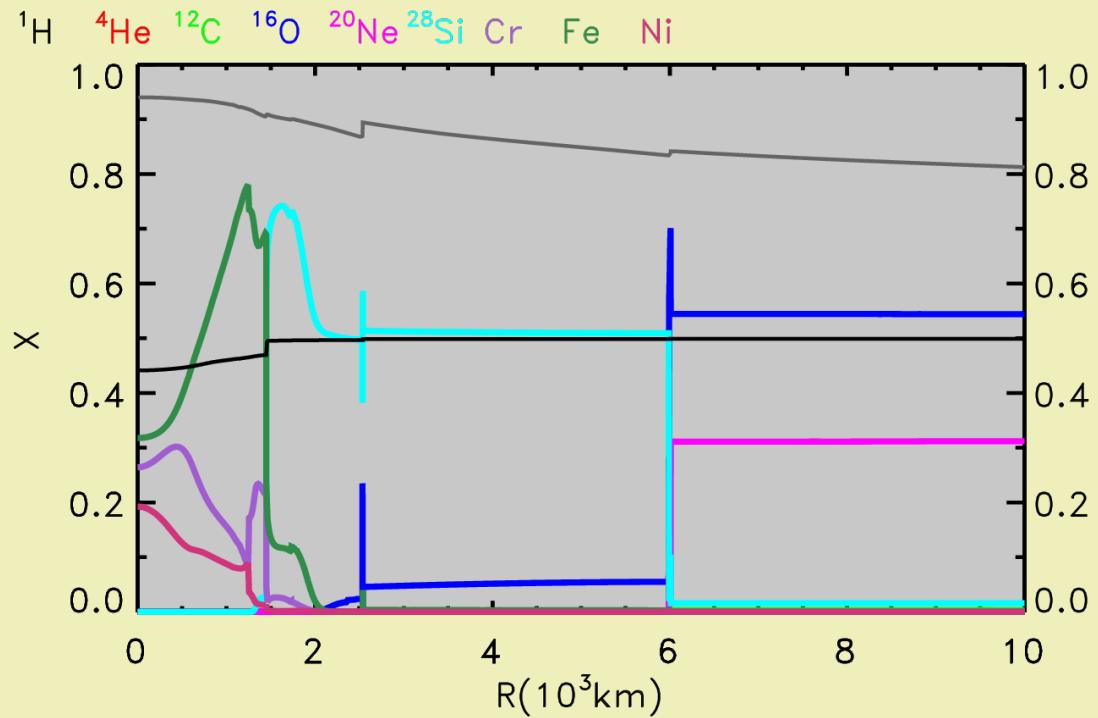
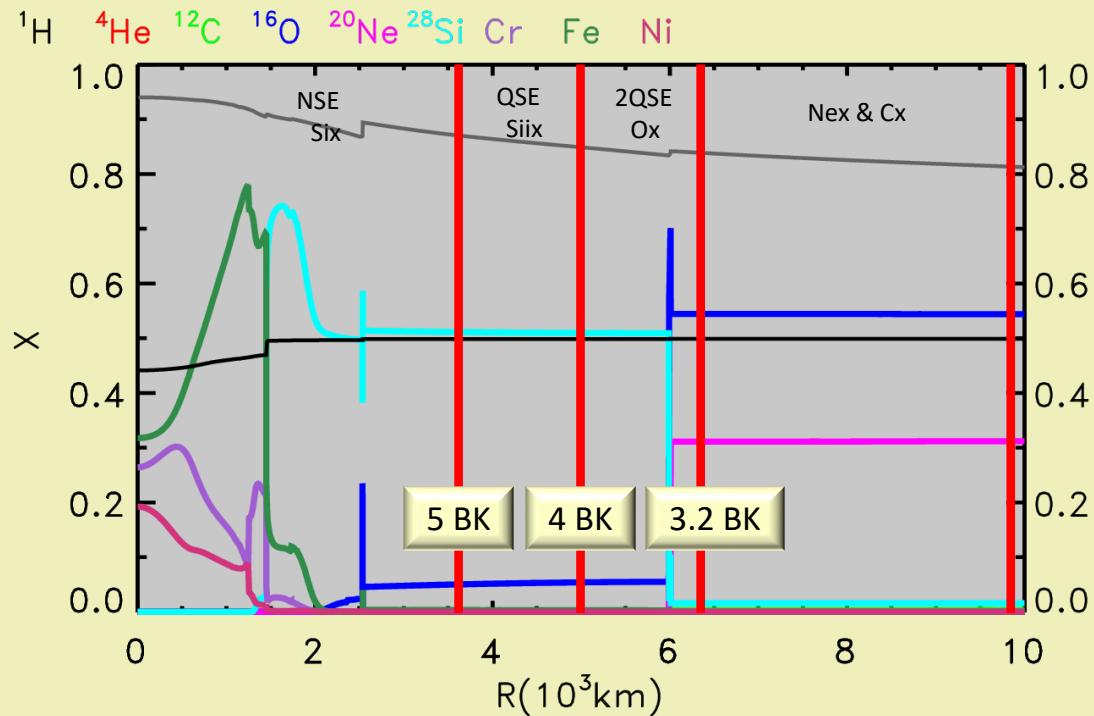
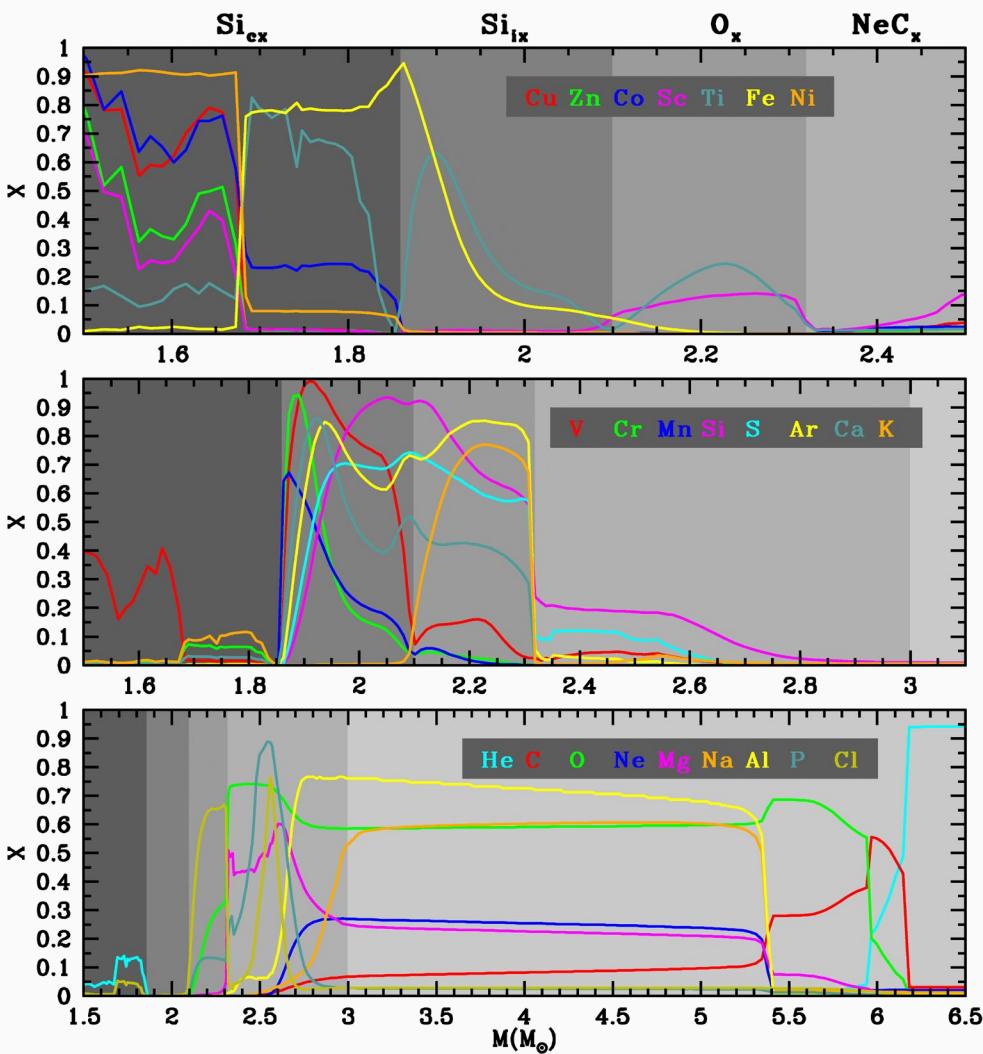


Table 2. For the runs presented in this paper, the mean shock radius (in units of 1000 kilometres) and mean shock speed (in units of 1000 km s^{-1}) at the end of each simulation. Note that the shock is still stalled at the end of the simulation only for the 2D and 3D 13- M_\odot models.

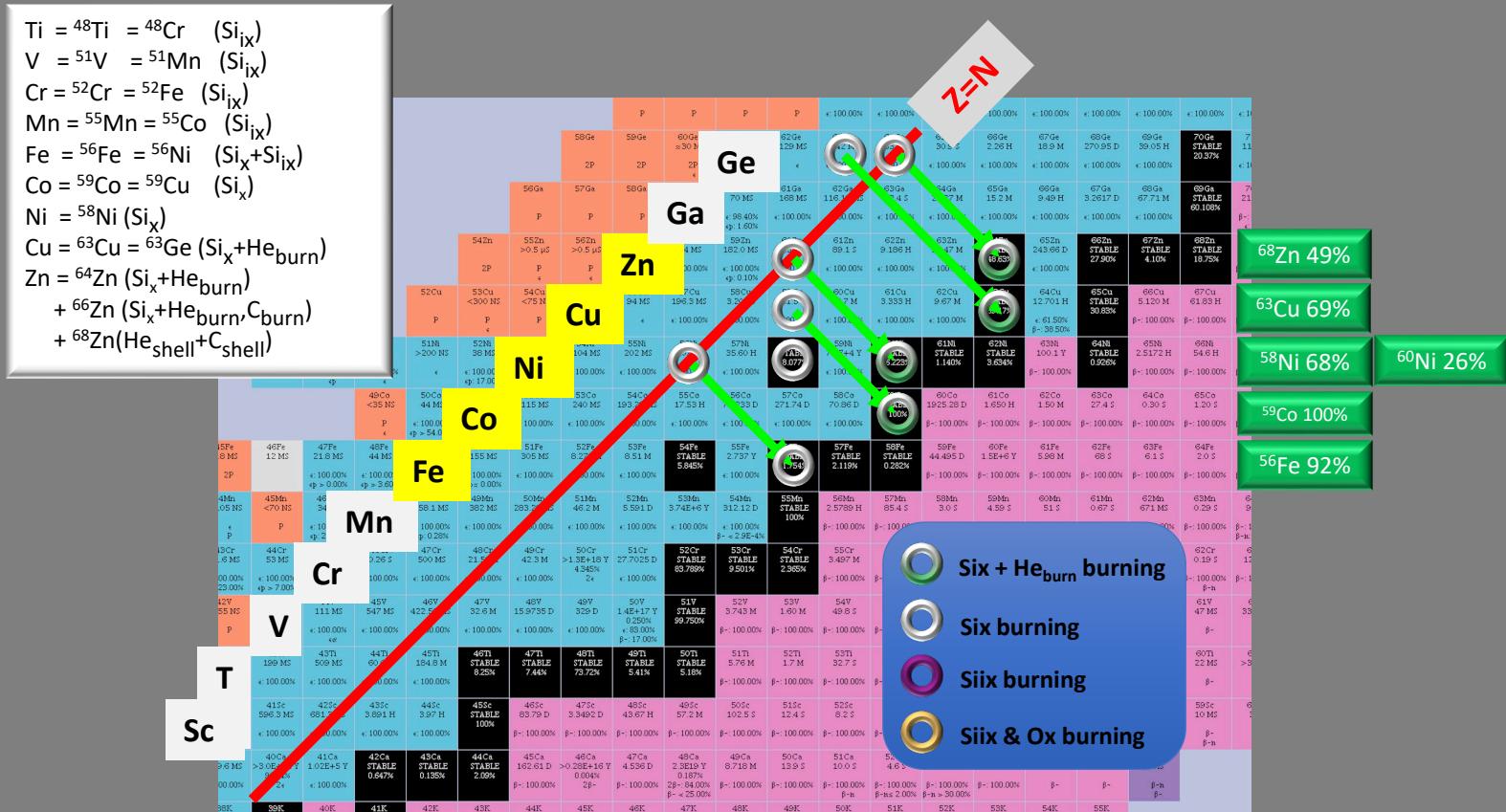
	t (final) (s)	Mean shock radius (1000 km)	Mean shock speed (1000 km s^{-1})
s9.0-2D	1.41	15.24	14.19
s9.0-3D	1.042	12.42	16.29
s10.0-2D	1.41	7.70	10.62
s10.0-3D	0.767	1.96	6.65
s11.0-2D	1.41	9.18	7.41
s11.0-3D	0.568	2.75	8.00
s12.0-2D	1.41	8.72	8.08
s12.0-3D	0.694	2.66	6.85
s13.0-2D	1.311	0.06	0.067
s13.0-3D	0.674	0.09	0.048



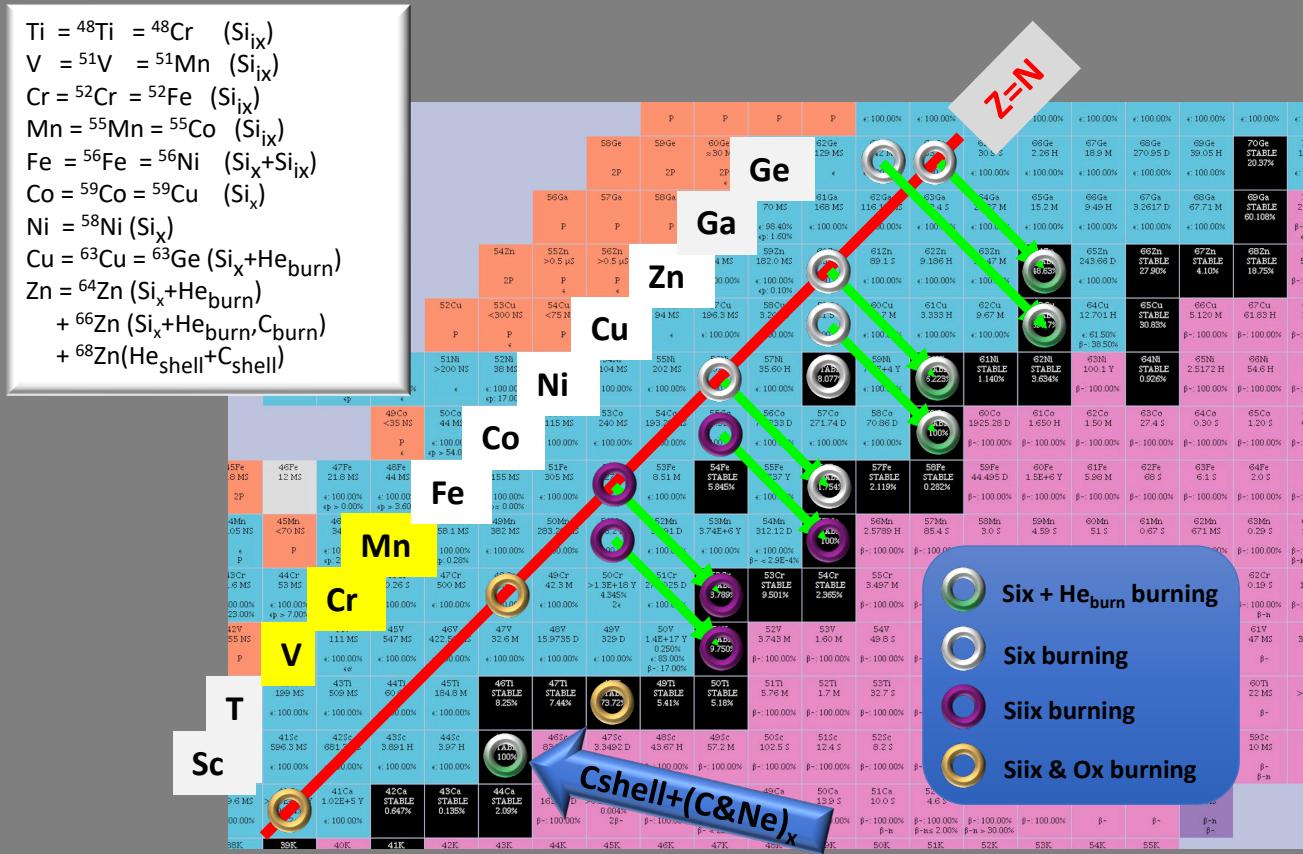


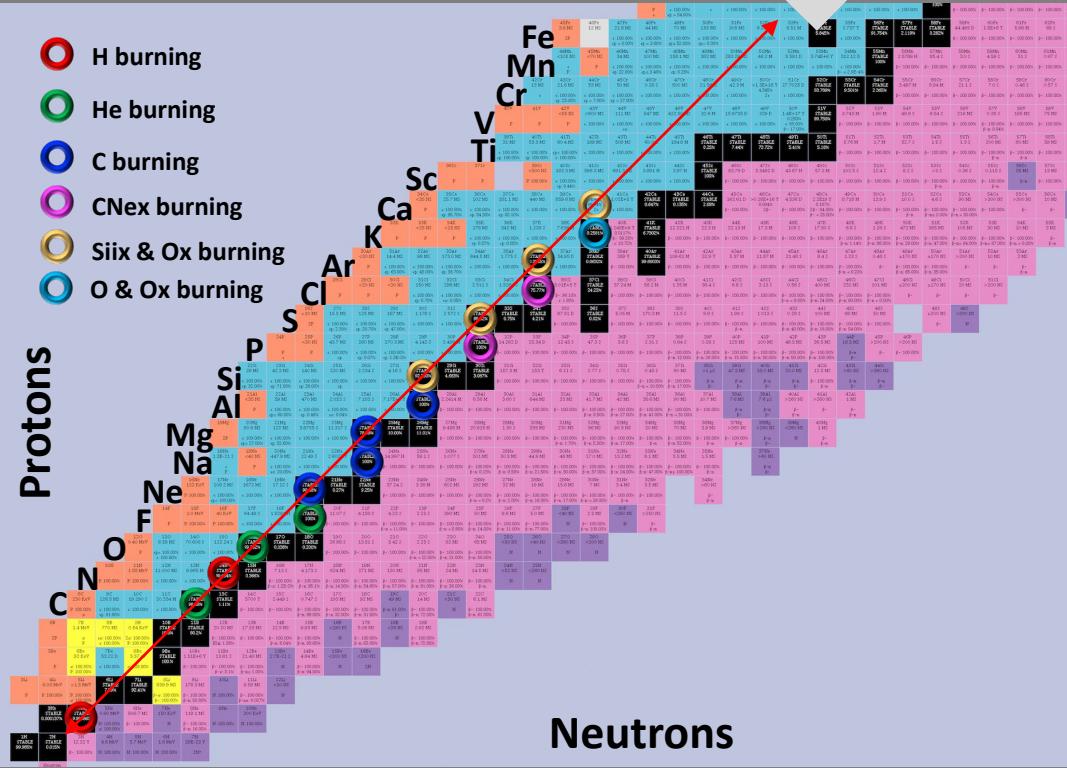


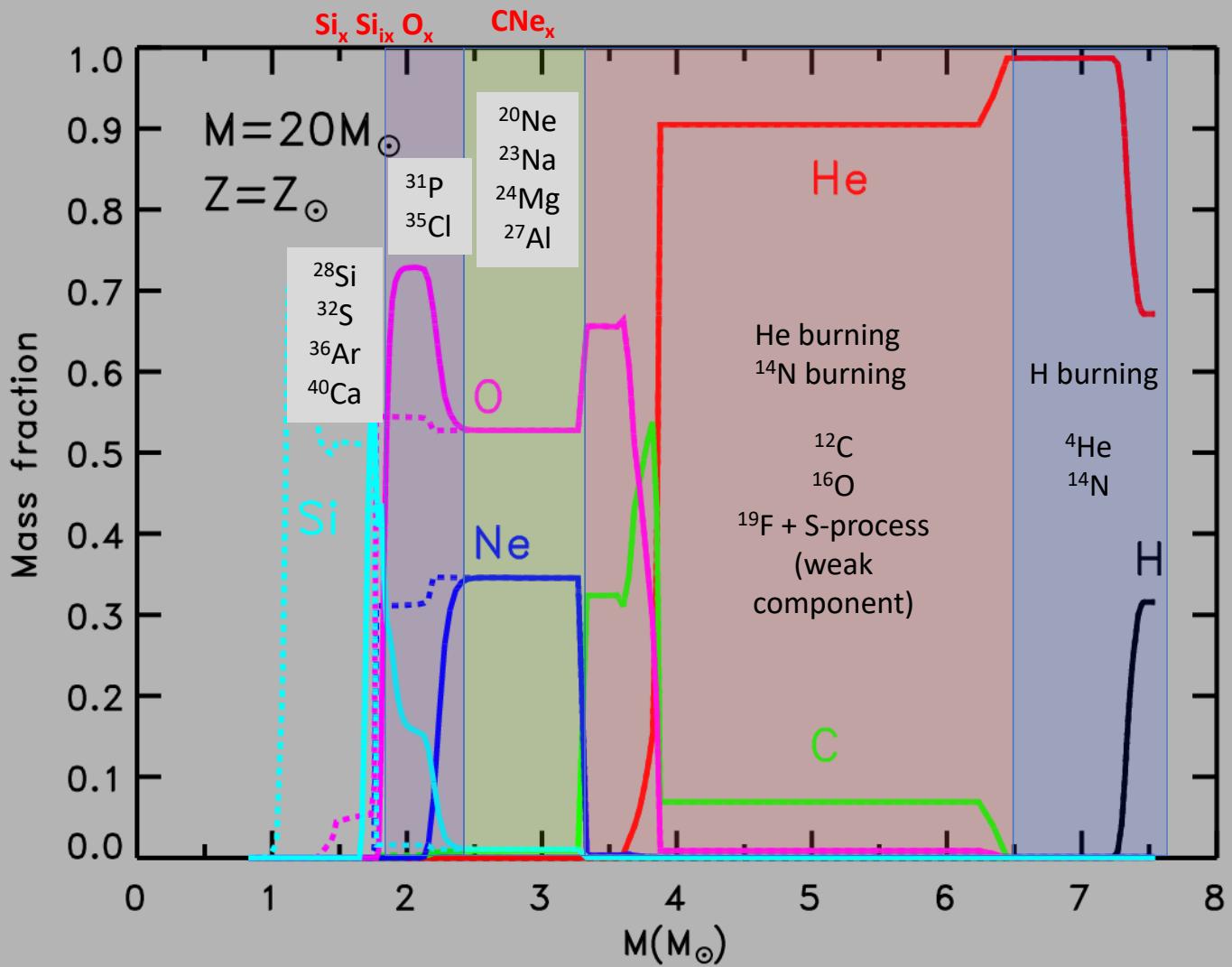
$Ti = ^{48}Ti = ^{48}Cr$ (Si_x)
 $V = ^{51}V = ^{51}Mn$ (Si_x)
 $Cr = ^{52}Cr = ^{52}Fe$ (Si_x)
 $Mn = ^{55}Mn = ^{55}Co$ (Si_x)
 $Fe = ^{56}Fe = ^{56}Ni$ ($Si_x + Si_x$)
 $Co = ^{59}Co = ^{59}Cu$ (Si_x)
 $Ni = ^{58}Ni$ (Si_x)
 $Cu = ^{63}Cu = ^{63}Ge$ ($Si_x + He_{burn}$)
 $Zn = ^{64}Zn$ ($Si_x + He_{burn}$)
 + ^{66}Zn ($Si_x + He_{burn}, C_{burn}$)
 + ^{68}Zn ($He_{shell} + C_{shell}$)



$Ti = {}^{48}Ti = {}^{48}Cr$ (Si_{ix})
 $V = {}^{51}V = {}^{51}Mn$ (Si_{ix})
 $Cr = {}^{52}Cr = {}^{52}Fe$ (Si_{ix})
 $Mn = {}^{55}Mn = {}^{55}Co$ (Si_{ix})
 $Fe = {}^{56}Fe = {}^{56}Ni$ ($Si_x + Si_{ix}$)
 $Co = {}^{59}Co = {}^{59}Cu$ (Si_x)
 $Ni = {}^{58}Ni$ (Si_x)
 $Cu = {}^{63}Cu = {}^{63}Ge$ ($Si_x + He_{burn}$)
 $Zn = {}^{64}Zn$ ($Si_x + He_{burn}$)
 + ${}^{66}Zn$ ($Si_x + He_{burn} + C_{burn}$)
 + ${}^{68}Zn$ ($He_{shell} + C_{shell}$)







recap

$$T > 5 \cdot 10^9 \text{ K}$$

NSE: all nuclei are at the equilibrium
Complete Explosive Silicon burning

$$Y_i = f(T, \rho, Y_e) \quad \begin{cases} Y_e > 0.49 \\ (\eta < 0.02) \end{cases}$$

^{56}Ni The most abundant nucleous

Sc Ti Co Ni Zn Cu

$$5 \cdot 10^9 \text{ K} > T > 4 \cdot 10^9 \text{ K}$$

QSE: ^{28}Si is not at the equilibrium
Incomplete Explosive Silicon burning

$$Y_i = f(T, \rho, Y_e)$$

V Cr Mn Si S Ar Ca

$$4 \cdot 10^9 \text{ K} > T > 3.3 \cdot 10^9 \text{ K}$$

2QSE: matter clustered in two groups, one peaked at
 ^{28}Si (99%) and the second one at Fe
Explosive Oxygen burning

$$Y_i = f(T, \rho, Y_e)$$

Si S Ar K Ca

$$3.3 \cdot 10^9 \text{ K} > T > 1.9 \cdot 10^9 \text{ K}$$

Explosive Neon and Carbon + C shell burning

Ne Na Mg Al P Cl

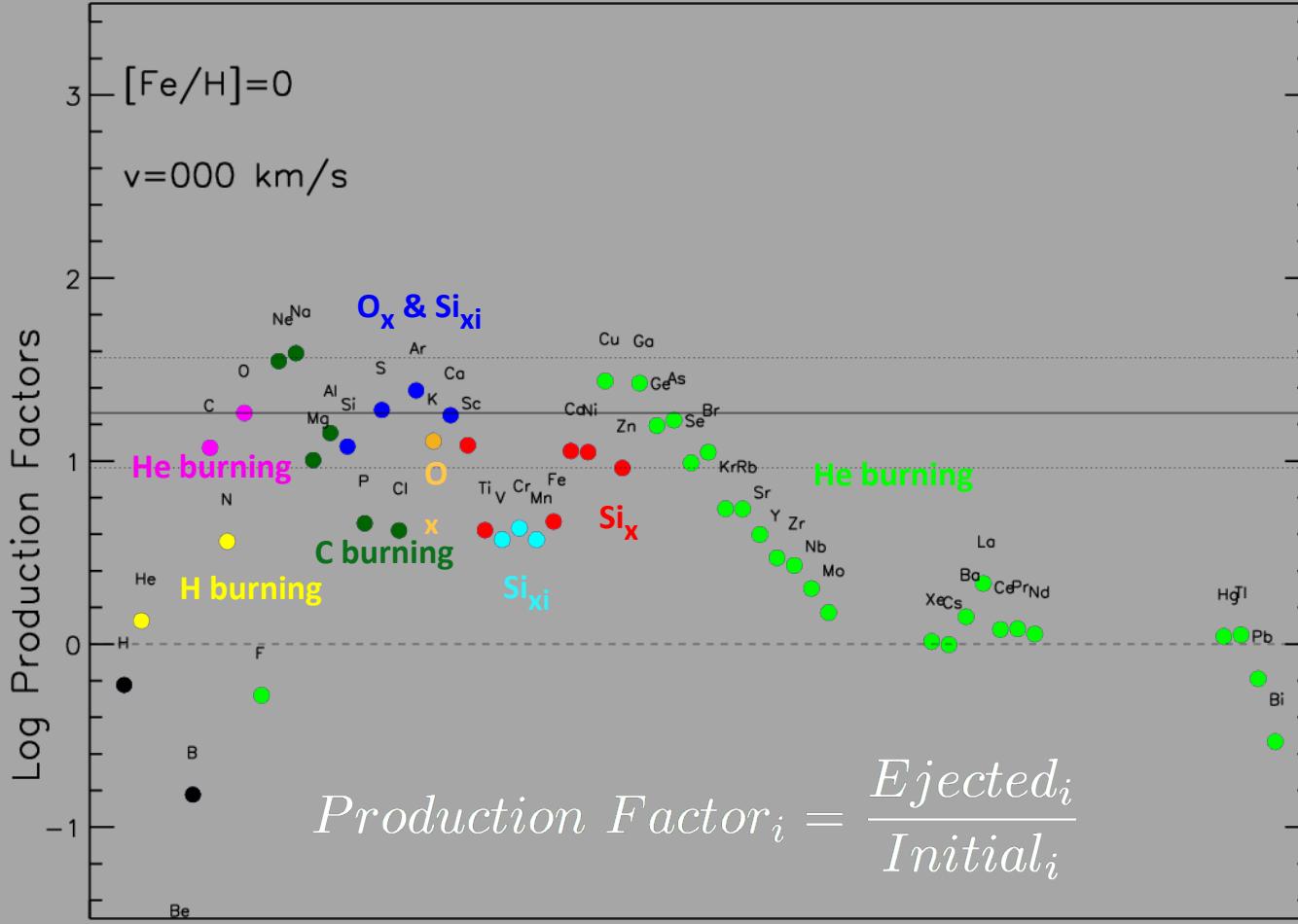
$$T < 1.9 \cdot 10^9 \text{ K}$$

NO Explosive burning

Element	produced	destroyed
He (⁴ He)	H _{c,s}	He _{c,s}
C (¹² C)	He _{c,s}	C _{c,s}
N (¹⁴ N)	H _{c,s}	He _{c,s}
O(¹⁶ O)	He _c	C _{c,s}
F (¹⁹ F)	He _s	
Ne (²⁰ Ne)	C _s	Ne _x
Na (²³ Na)	C _s	Ne _x
Mg (²⁴ Mg)	C _s	Ne _x
Al (²⁷ Al)	C _s	Ne _x
Si (²⁸ Si)	Si _{ix} O _x Ne _x	
P (³¹ P)	Ne _x C _s	Ne _x
S (³² S)	Si _{ix} O _x Ne _x	
Cl	³⁵ Cl O _x Ne _x - ³⁷ Cl C _s	³⁷ Cl Ne _x
Ar (³⁶ Ar)	Si _{ix} O _x	
K (³⁹ K)	He C _s O _x	O _x
Ca (⁴⁰ Ca)	Si _{ix} O _x	

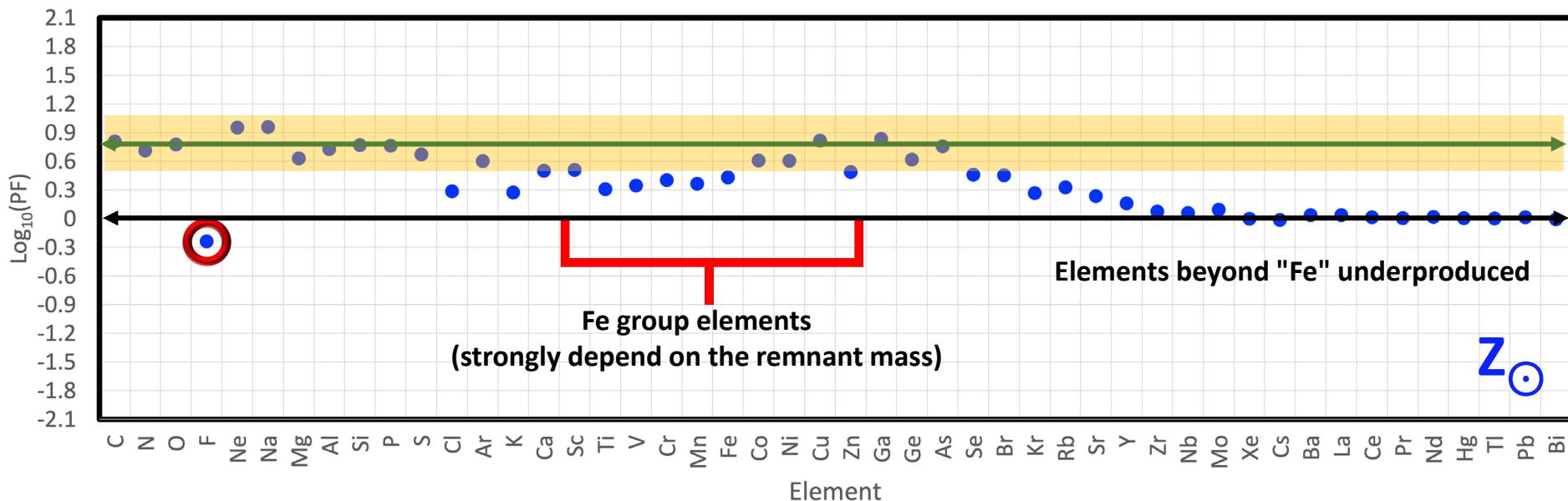
Element	produced	destroyed
V (⁵¹ V)	Si _{ix}	
Cr (⁵² Fe)	Si _{ix}	
Mn (⁵⁵ Fe ⁵⁵ Co)	Si _{ix}	
Sc (⁴⁵ Sc, ⁴⁵ Ca)	He C _s Si _x	Ne _x
Ti (⁴⁸ Cr)	O _x Si _x Si _{ix}	
Fe (⁵⁶ Ni ⁵⁶ Fe ⁵⁴ Fe)	Si _{ix} O _x	
Co (⁵⁹ Ni)	He _c Ne _x Si _x	
Ni (⁵⁸ Ni)	Si _x	
Cu(⁶³ Cu)	He _c Ne _x Si _x	
Zn(⁶⁴ Zn)	He _c Ne _x Si _x	

Solar metallicity



Production Factors produced by a generation of massive stars:

$$PF = \frac{X_{\text{ejecta}}^i}{X_{\text{initial}}^i} = \frac{Y^i/M_{\text{ejected}}}{X_{\text{initial}}^i}$$



Critical masses (model dependent)

