

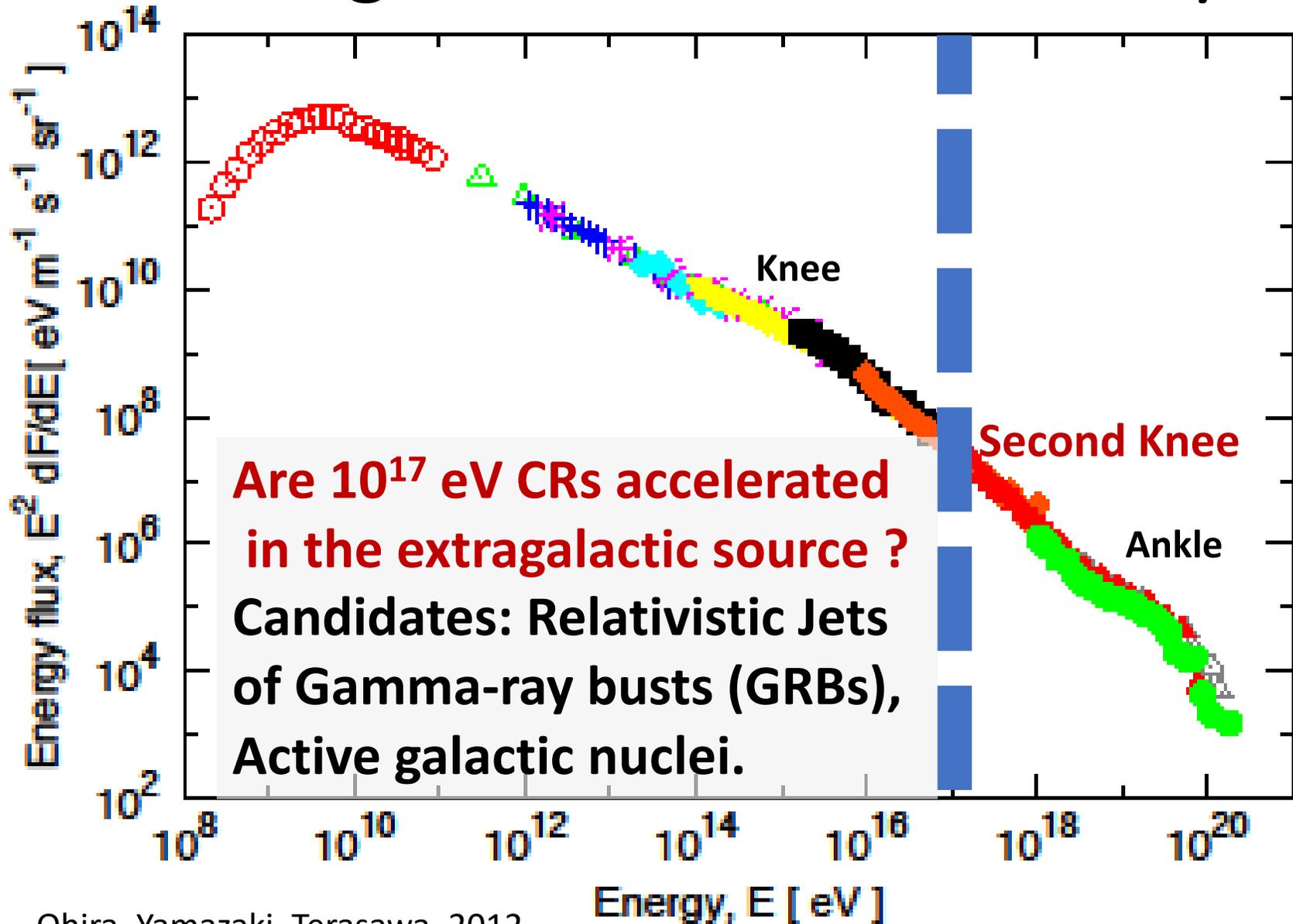
The Weibel-mediated shocks Propagating into Inhomogeneous Media

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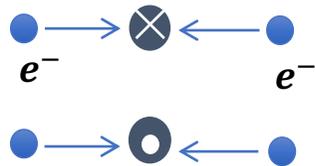
The origin of 10^{17} eV cosmic-rays (CRs)



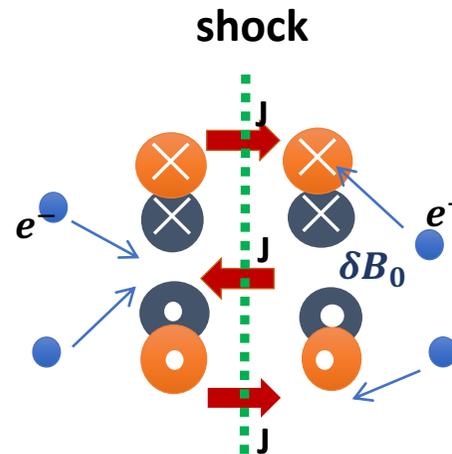
The relativistic jet interacts with interstellar medium (ISM) or circum-stellar medium (CSM).

→ **Relativistic collisionless shocks mediated by the Weibel instability**

Counter-streaming plasmas
= anisotropic-temperature plasmas



Perturbation of
magnetic field δB_0



The streaming plasmas are deflected by δB , so that currents are generated. The currents amplify δB . → Unstable.

The Weibel instability makes magnetic fields.

Particle-in-Cell simulations of Relativistic Collisionless Shocks:

CRs are accelerated up to 10^{17} eV if the b-fields persist in the far downstream region.

For $r_g \gg \lambda_{\delta B}$, $D(E) \sim c \lambda_{\delta B} \left(\frac{E}{eB \lambda_{\delta B}} \right)^2$

$$\frac{\gamma_{\max}}{\gamma_0} \sim (2 \epsilon_B \lambda_{\delta B} \omega_{pi} t)^{1/2}$$

At $t = t_{\text{age}} \sim R_{\text{dec}} / \Gamma c$

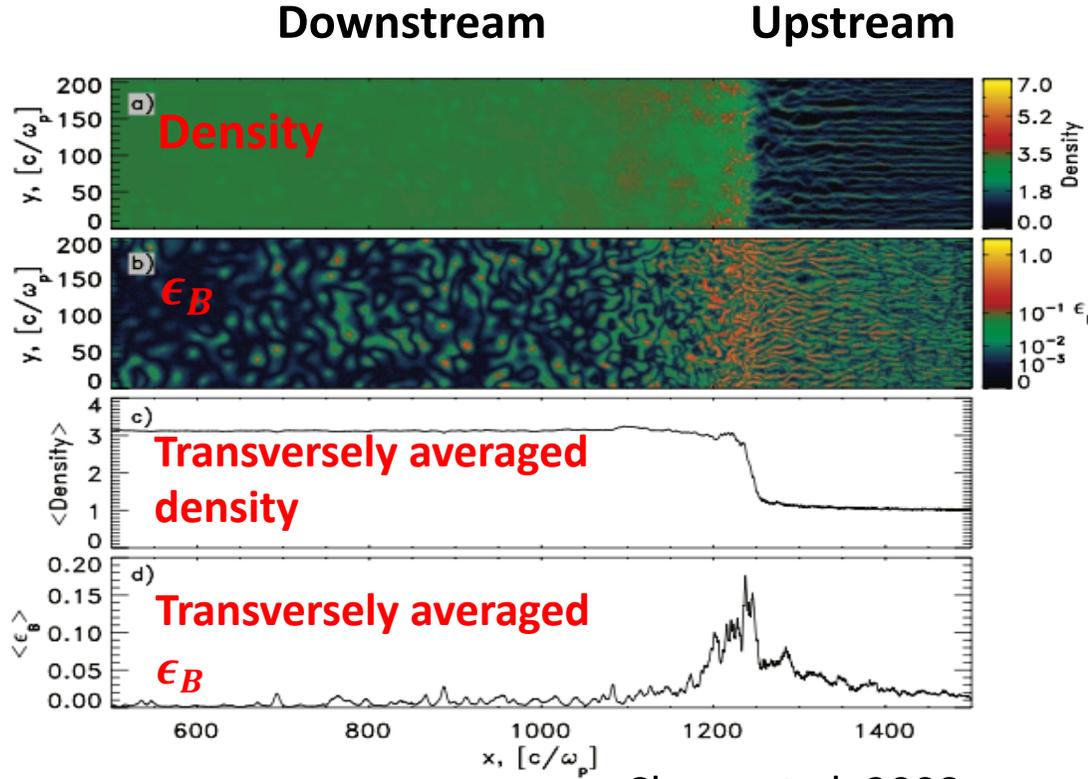
$$R_{\text{dec}} \sim 10^7 E_{0,54}^{1/3} n_{\text{ISM}}^{-1/3} \Gamma_{0,2}^{-2/3} \text{ cm}$$

$$\epsilon_{B,-1} \lambda_{c/\omega_{pi}} \sim 0.1 \text{ (from PIC simulation)}$$

$$\gamma_{\max,i} \approx 10^8 E_{0,54}^{3/4} n_{\text{ISM},0}^{-1/2} R_{\text{dec},17}$$

In ISM frame

Shalchi & Dosch 2009; Plotnikov et al. 2011, 2013; Sironi et al. 2013



Chang et al. 2008

However, the Weibel generated b-field rapidly decays. $\lambda_{\delta B} \ll r_g$ is inefficient for particle acceleration.

Amplified B-field in GRBs afterglows

Observations:

B-fields are strongly amplified in **large** emission regions.

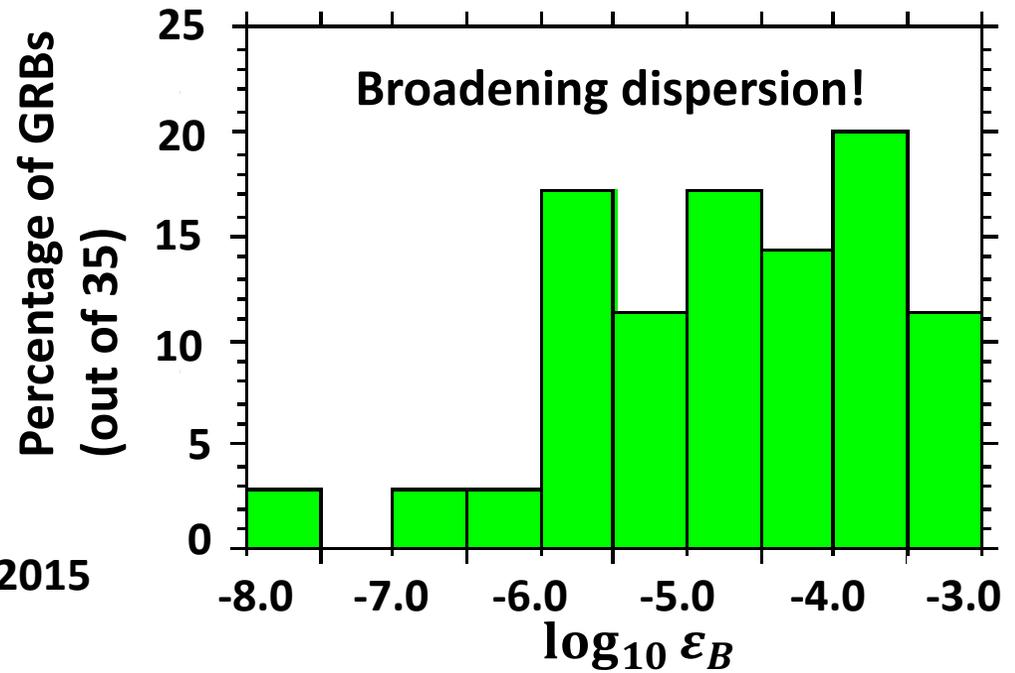
$$\epsilon_B = \frac{B^2/8\pi}{4\Gamma mnc^2} \sim 10^{-8} - 10^{-3} \quad \text{Average: } 10^{-5}$$

(Santana et al. 2014)

→ **The post-shock b-field is amplified to about 100 times the shock-compressed value.**

the process of the b-field amplification is an open question.

Kumar et al. 2015



Previous PIC simulation :

Nonlinear evolution of the Weibel instability for shocks in uniform plasmas → B-fields rapidly decay.

There should be density fluctuations, $\delta\rho$, in the ISM or CSM.

Ex.)

- **Large-scale inhomogeneity:**

The injection scale of $\delta\rho$ is 1-100 pc for the ISM turbulence.

Stellar winds or binary systems can generate $\delta\rho$ with a smaller scale.

Armstrong et al. 1995, Chugai & Danziger 1994, Smith et al. 2009, Yalinewich & Zwart 2019

- **Small-scale inhomogeneity:**

The CR precursor can generate a much smaller scale less than the precursor scale.

Drury & Falle 1986, Ohira 2013, 2014, 2016

Purpose of our study

In previous MHD simulation :

Large-scale magnetic field amplification by the turbulent dynamo

Inoue et al. 2011, Mizuno et al. 2014

However, MHD simulation cannot solve physics of the kinetic scale.

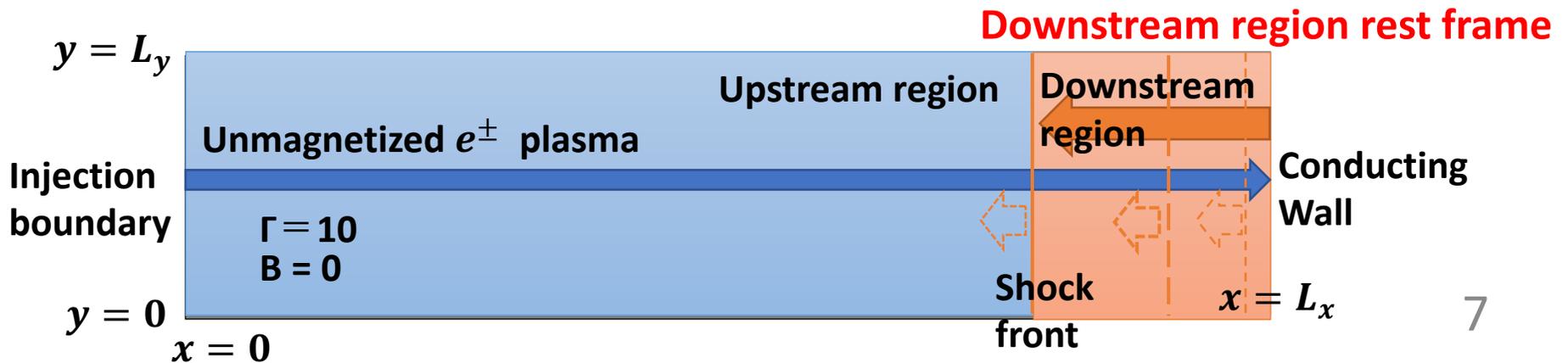
In this study:

We investigate a magnetic field amplification in the downstream region of the relativistic collisionless shock propagating into **the inhomogeneous ISM**, by using two-dimensional PIC simulations.

Simulation Set Up

- Two-dimensional electromagnetic PIC code (pCANS)
- X-Y plane with periodic boundary condition in the y-direction.
- Box Size: $L_x = 1.2 \times 10^4 c/\omega_p$, $L_y = 86 c/\omega_p$, Cell Size: $\Delta X = \Delta Y = 0.1 c/\omega_p$
- Unmagnetized e^\pm plasmas ($n_{e^+} = n_{e^-}$) with Lorentz factor $\Gamma = 10$, $v_{th} = 0.1c$.
- Initial spatial distribution of e^\pm :

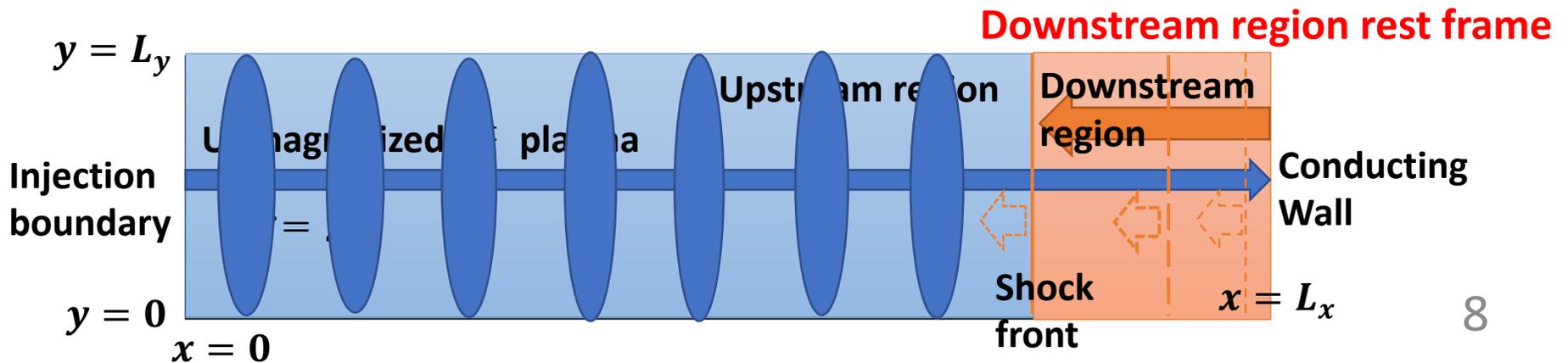
$$n(x, y) = n_0 \{1 + 0.5 \sin(2\pi kx/L_x)\}$$



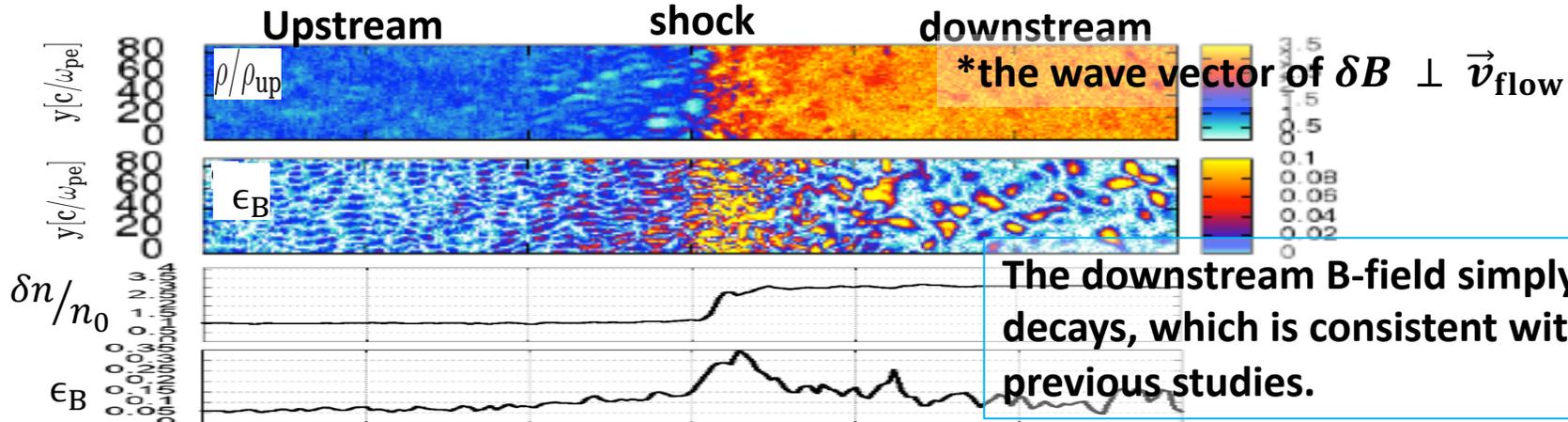
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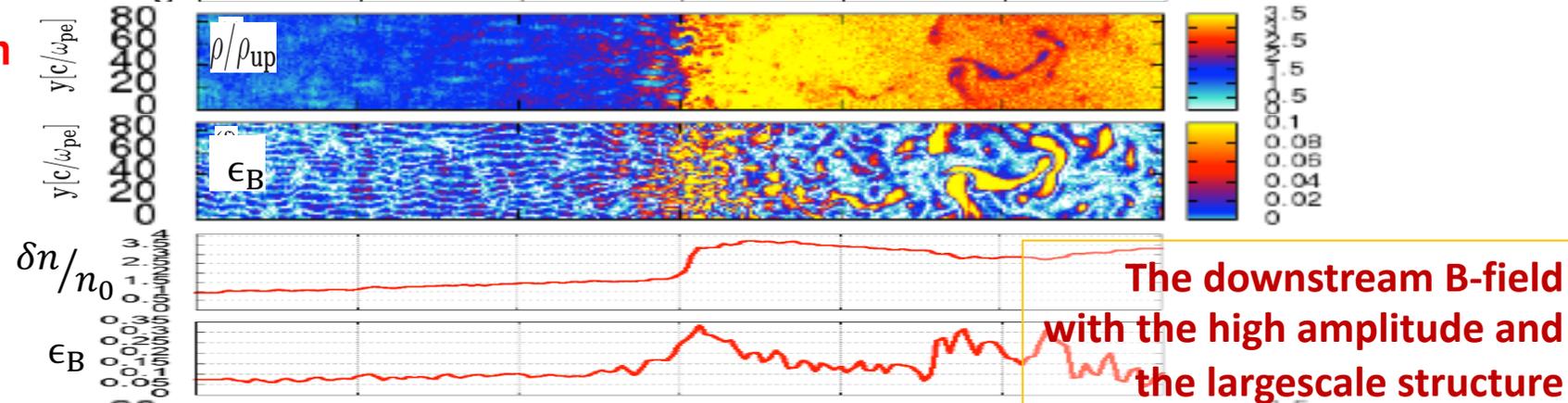
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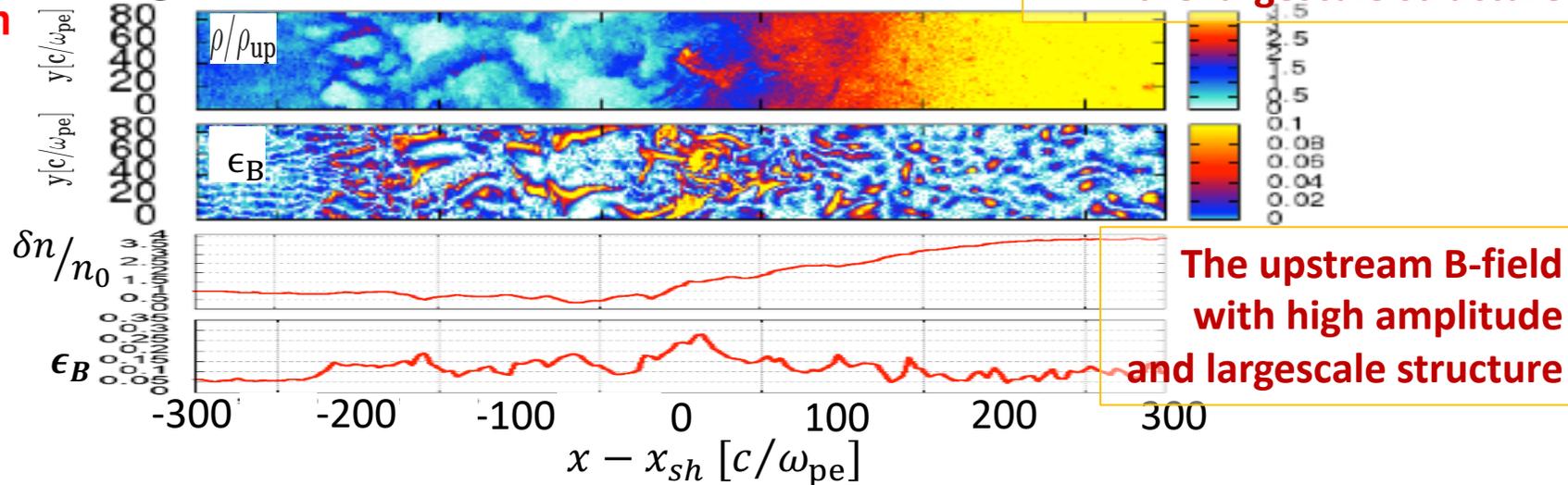
Uniform
(at $t\omega_{pe} = 5000$)



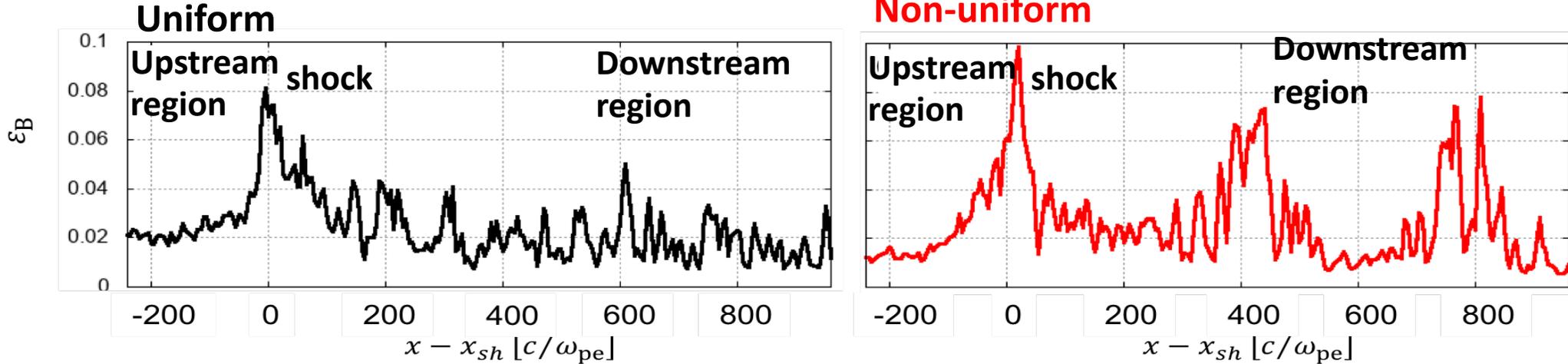
No-uniform
(shock in the
 $\frac{n_0 + \delta n}{n_0} = 1.5$
region)
at $t\omega_{pe} = 4800$



No-uniform
(shock in the
 $\frac{n_0 + \delta n}{n_0} = 0.5$
region)
at $t\omega_{pe} = 5200$



The spatial evolution of the transversely averaged ϵ_B in the larger downstream region ($t = 6100\omega_p^{-1}$)



Uniform:

The magnetic field simply decays.

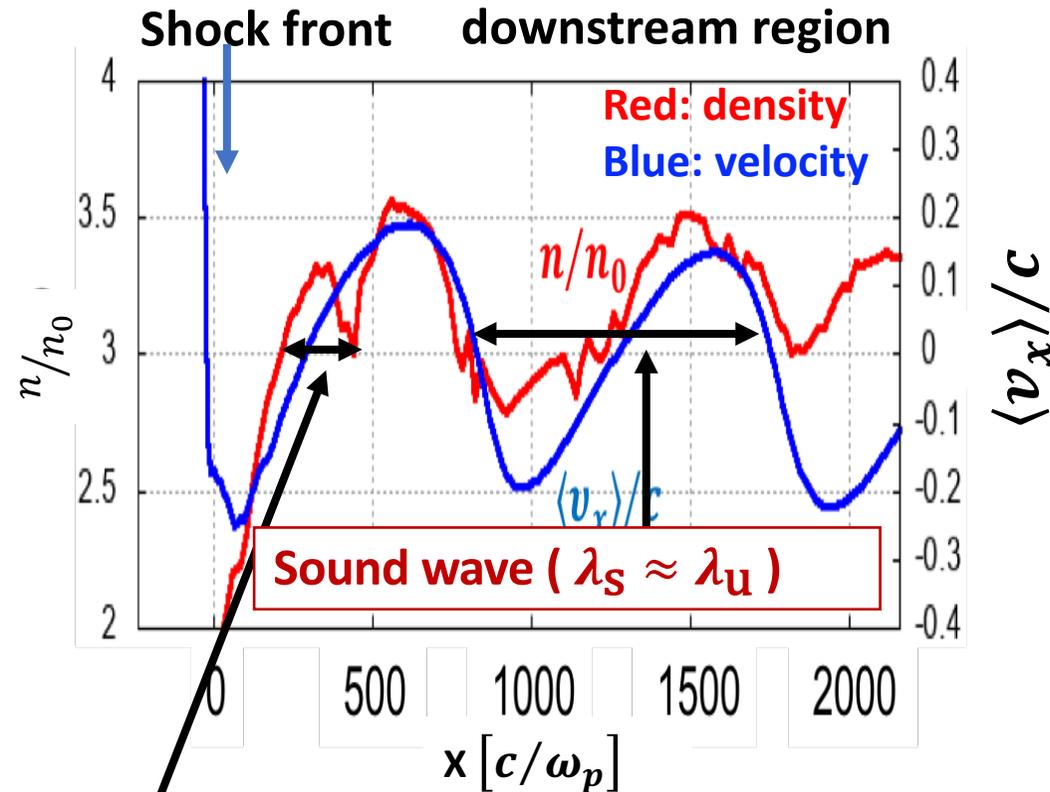
Non-uniform:

The magnetic field decays much slower than uniform case. Two peaks of ϵ_B were generated when the shock was passing through the low-density region.

The inhomogeneity of the upstream density is crucial in the generation of the downstream magnetic field.

$$L_x \times L_y = 1.2 \times 10^4 \times 86 \text{ c}/\omega_p$$

Density & velocity in the x-direction ($t = 6100 \omega_p^{-1}$)



Entropy wave ($\lambda_e \approx \frac{1}{3} \lambda_u$)

- λ_e : wavelength of the entropy wave
- λ_s : wavelength of the sound wave
- λ_u : wavelength of the upstream incoming wave

Sound waves and entropy waves are generated as expected by hydrodynamic analysis of the shock-wave interaction.

Each wavelength is consistent with result from the linear analysis for the shock front.

The sound wave might play crucial role on the particle acceleration!

Summary

- ✓ **There really are some density fluctuations in ISM or CSM**
We investigated the b-field amplification in the downstream region of collisionless shocks propagating into inhomogeneous media.
- ✓ **A larger-scale magnetic field is generated in the shock precursor region and hardly decay after it is advected downstream.**
- ✓ **The sound and entropy wave are generated in the downstream region. The sound wave could be crucial for the particle acceleration.**
- ✓ **Temperature anisotropy is produced by the sound wave and the diffusion of high-energy particles.**
It is expected that larger magnetic field is generated by the Weibel instability far downstream compared with the uniform case.