

Gamma-rays from large-scale outflows in starburst galaxies

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Starburst Galaxies



A starburst galaxy is a galaxy undergoing an exceptionally high rate of star formation.
 Sometimes form when galaxies collide

Enhanced star formation results in:

- * Large number of massive stars with strong winds.
- * High infrared luminosity (from heated gas).
- * High rate of supernova explosions.
- High density of cosmic rays (CRs) accelerated by the SNRs.
- * Radio and gamma-ray emission produced by these CRs.

e.g. Paglione et al. 1996, Romero & Torres 2003, Torres 2004, etc...



Superwinds

The combined effect of SN and stellar winds creates a high-temperature bubble in the starburst. This hot gas escapes powering a galactic superwind.



Chevallier & Clegg 1985

The large number of corecollapse supernovae in starbursts results in the merge of SNRs.

Shocks formed in these collisions thermalize the energy released by the explosions creating a cavity filled with hot $(T \sim 10^8 \text{ K})$ gas that is unbound by the gravitational potential.

The hot gas expands adiabatically and escapes the system sweeping up cooler and denser gas from the disk.

Superwinds



Numerical simulations show that the wind expands preferentially along the minor axis of any disk-like galaxy forming a bipolar flow, sweeping up cooler, denser, ambient disk or halo gas, and stripping ambient gas from the walls of the cavity.

Superwinds

- Superwinds are complex multiphase outflows of cool, warm, and hot gas, dust, and magnetized relativistic plasma.
- Evidence for superwinds include: detection of metals in the IGM, measurements of outflows in Halpha and molecular lines, X-ray diffuse emission above the galaxies, radio halos, hard X-ray radiation with T~10⁷⁻⁸ K, etc.

M82 - Composite: *Chandra* + *HST* + *Spitzer*



Scaling relations

 $\dot{E} = \epsilon \dot{E}_{*}. \quad \epsilon: \text{ thermalization parameter}$ $\dot{M} = \dot{M}_{*} + \dot{M}_{\text{ISM}} = \beta \dot{M}_{*}. \quad \beta: \text{ mass load}$ $\dot{E}_{*} = 7 \times 10^{41} \text{ (SFR}/M_{\odot} \text{yr}^{-1}) \text{ erg s}^{-1}$ $\dot{M}_{*} = 0.26 \text{ (SFR}/M_{\odot} \text{yr}^{-1}) M_{\odot} \text{yr}^{-1}$ $\dot{\tau}_{\text{SN}} = 0.02 \text{ (SFR}/M_{\odot} \text{yr}^{-1}) \text{ yr}^{-1}$

SFR
$$\approx 17 \frac{L_{IR}}{10^{11} \text{erg s}^{-1}} M_{\odot} \text{ yr}^{-1}$$

As a consequence of the previous relations the outflow final velocity depends on the thermalization efficiency and the mass load factor:

$$v_{\infty} = \sqrt{\frac{2\epsilon \dot{E}_*}{\beta \dot{M}_*}} \approx 3000 \sqrt{\frac{\epsilon}{\beta}} \text{ km s}^{-1}$$

 β can be determined from high-resolution CO observations.

The nearby starburst NGC 253

- * Edge-on galaxy.
- * Very nearby (~3 Mpc).
- High star forming rate (~3 solar masses / yr in the inner starburst).
- * Superwind detected in X-rays, H-alpha, CO.
- * Superwind is asymmetric (NW bubble stronger).
- * Non-thermal radiation form the central source.
- * Radio halo (also non-thermal).

Radio polarization + X-rays

VLA + XMM-Newton



Heesen et al. 2009



LETTER

450 | NATURE | VOL 499 | 25 JULY 2013

doi:10.1038/nature12351

Suppression of star formation in the galaxy NGC 253 by a starburst-driven molecular wind

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Table 1. Physical properties of NGC 253 and its superwind.

Starburst parameters	Value
d: Distance [Mpc]	$2.6^1 - 3.9$
$L_{\rm IR}$: Infrared luminosity $[L_{\odot}]$	$1.7 imes 10^{10}$
SFR: Star forming rate $[M_{\odot} \text{ yr}^{-1}]$	3
β : Mass loading factor	12
\dot{M} : Mass outflow $[M_{\odot} \text{ yr}^{-1}]$	9
$L_{\gamma}[E > 200 \text{ MeV}]$: γ -ray luminosity [erg s ⁻¹]	$4.3 imes10^{39}$
$T_{\rm c}$: Temperature of the central region [K] ²	$2 imes 10^7$
SW region parameters	Value
R: Radius of the SW bubble [kpc]	5
$L_{\rm x}$: X-ray luminosity [erg s ⁻¹]	$5 imes 10^{38}$
$T_{\rm h}$: Temperature of the gas in the bubble [K]	$3 imes 10^6$
n : Particle density $[\rm cm^{-3}]$	2×10^{-3}
B: Magnetic field $[\mu G]$	5
$v_{\rm A}$: Alfvén velocity [km s ⁻¹]	240
$v_{\rm s}$: Sound speed [km s ⁻¹]	164

Superwind region



DSA

In the NE superwind bubble:

Parameters	$\epsilon = 1$	$\epsilon = 0.75$
\dot{E} : Mechanical luminosity of the superwind [erg s ⁻¹]	1.5×10^{42}	1.1×10^{42}
$v_{\rm rev}$: Velocity of the reverse shock [km s ⁻¹]	866	750
$v_{\rm shell}$: Velocity of the expanding shell [km s ⁻¹]	494	407

$$\frac{dE}{dt} = \frac{3}{20} ec \left(\frac{D}{D_{\rm B}}\right)^{-1} \left(\frac{v_{\rm rev}}{c}\right)^2 B$$
$$t_{\rm acc} \approx 2.1 \left(\frac{D}{D_{\rm B}}\right) \left(\frac{v_{\rm rev}}{1000 \,\,\mathrm{km \, s^{-1}}}\right)^{-2} \left(\frac{B}{\mu \rm G}\right)^{-1} \left(\frac{E}{\rm GeV}\right) \,\,\mathrm{yr}.$$

Losses









For $t_{\rm acc} = \tau$, in the Bohm limit $D = D_{\rm B}$ and for a thermalization $\epsilon = 1$:

$$E_{\text{max}}^p = 1.7 \times 10^{16} \text{ eV} \text{ protons}$$

 $E_{\text{max}}^{\text{Fe}} = 4.4 \times 10^{17} \text{ eV}$ iron nuclei.

The CR luminosity is:

$$L_{\rm CR} = 4\pi\xi R_{\rm shock}^2 \rho v_{\rm shock}^3 \sim \xi \dot{M} v_{\rm shock}^2$$

With an efficiency of 10%:

$$L_{\rm CR} \sim 3.2 \times 10^{41} \ {\rm erg \ s^{-1}}$$





Fig. 6. Proton distributions for a = 1 and a = 100 with $a = L_p/L_e$ in the case of DSA in the reverse shock, assuming thermalization $\epsilon = 1$.



Fig. 7. Electron distributions for a = 1 and a = 100 with $a = L_p/L_e$ in the case of DSA in the reverse shock, assuming thermalization $\epsilon = 1$.

Fig. 8. Spectral energy distribution in the case of DSA with ratio a = 1, magnetic field $B = 5 \ \mu$ G, thermalization $\epsilon = 1$, and shock efficiency $\xi = 0.012$.



Fig. 9. Spectral energy distribution in the case of DSA with ratio a = 100, magnetic field $B = 5 \ \mu$ G, thermalization $\epsilon = 1$, and shock efficiency $\xi = 6 \times 10^{-3}$.



Shocked cloud



Fractal and spherical clouds: different behavior.







Cooper et al. 2009

Losses and SEDs



Table 1. Parameters of the models. The magnetization $\beta = 0.9$ and the wind velocity $v_w = 1000$ km s⁻¹ are the same in both cases.

Model	$R_{\rm c}$ pc	$n_{ m w} \ [m cm^{-3}]$	$n_{ m c} \ [m cm^{-2}]$	$v_{ m sw}$ [km s ⁻¹]	$v_{ m sc}$ [km s ⁻¹]
M1	5	0.01	100	1320	4.2
M2	100	0.01	10	1292	13.2

Table 2. Dynamical timescales

Model	$t_{ m crush}$ [Myr]	$t_{ m KH}$ [Myr]	$t_{ m RT}$ [Myr]	$t_{ m \Lambda_{sc}} \ [m Myr]$	$t_{\Lambda_{ m sw}} \ [{ m Myr}]$
M1 M2	0.37 2.39	0.49 3.09	$0.37 \\ 2.39$	$\begin{array}{c} 1.15\times 10^{-3} \\ 1.46\times 10^{-5} \end{array}$	64.78 60.53





a=100



Conclusions

- * NGC 253 and other starbursts are sources of CRs below the Ankle in the CR spectrum.
- Gamma rays up to 10¹⁵ eV can be produced in the halo. This emission can be as important as emission from the disk, but reaching higher energies.
- * Discrete gamma-ray sources should exist in the halo because of the interaction of the superwind with fragments of the disks.
- * Some X-ray sources in the halo might be also associated with such a phenomenon as well.



Thanks!





FIG. 7.—A configuration image at 3.6 cm. Contours are at logarithmic intervals of 2^{1/2}, beginning at 0.09 mJy beam⁻¹.



